

# Numerical Simulation of the Precipitation Kinetics of Nitrides and Carbides in Microalloyed Steel

By

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*'Prediction is very difficult, especially about the future.'*

Nils Bohr (1885-1962), Nobel laureate in Physics (1922)

*'Everything should be made as simple as possible, but not simpler.'*

Albert Einstein (1879-1955), Nobel laureate in Physics (1921)

# Preface

The present work is submitted for the degree 'Doktor der technischen Wissenschaften' at the 'Faculty of Mechanical Engineering' of the 'Graz University of Technology'. The work was carried out in the frame of the project 'Modelling and simulation of the early stages of precipitation of second phase particles in technically relevant structural high-performance materials' which was part of module B in the Christian Doppler Laboratory 'Early Stages of Precipitation'. The project was realised in the time between November 2007 and October 2010 under the supervision of Prof. Ernst Kozeschnik. It was financed by the 'Christian Doppler Forschungsgesellschaft (CDG)'.

I declare that I have authored this thesis independently, that I have not used other than the declared sources/resources, and that I have explicitly marked all material which has been quoted either literally or by content from the used sources.

Rene Radis  
Graz, June 1, 2010

# List of Publications

The author of this dissertation was also author or co-author of several papers, published during his time as a PhD student. The following list shows the references of these papers, which cover investigations on microalloyed steels, nickel-based superalloys and maraging steels.

## Refereed Journal Publications:

1. R. Radis and E. Kozeschnik. Concurrent Precipitation of AlN and VN in Microalloyed Steel. *Steel Research International*, in press.
2. R. Radis and E. Kozeschnik. Kinetics of AlN Precipitation in Microalloyed Steel. *Modelling and Simulation in Materials Science and Engineering*, 18(055003):1-16, 2010.
3. R. Radis, G. A. Zickler, M. Stockinger, C. Sommitsch and E. Kozeschnik. Interaction of the Precipitation Kinetics of  $\delta$  and  $\gamma'$  Phases in Nickel-Base Superalloy ATI Allvac<sup>®</sup>718Plus<sup>™</sup>. *Materials Science Forum*, 638-642:2712-2717, 2010.
4. R. Radis and E. Kozeschnik. Precipitation Kinetics of Aluminium Nitride in Austenite in Microalloyed HSLA Steels. *Materials Science Forum*, 636-637:605-611, 2010.
5. R. Radis, M. Schaffer, M. Albu, G. Kothleitner, P. Pölt and E. Kozeschnik. Multi-Modal Size Distributions of  $\gamma'$  Precipitates During Continuous Cooling of UDIMET 720 Li. *Acta Materialia*, 57:5739-5747, 2009.



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7. R. Schnitzer, R. Radis, M. Nöhner, M. Schober, R. Hochfellner, S. Zinner, E. Povoden-Karadeniz, E. Kozeschnik and H. Leitner. Reverted Austenite in PH 13-8 Mo Maraging Steels. *Materials Chemistry and Physics*, 122:138-145, 2010.
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9. E. Kozeschnik, J. Svoboda, R. Radis and F. D. Fischer. Mean-Field Model for the Growth and Coarsening of Stoichiometric Precipitates at Grain Boundaries. *Modelling and Simulation in Materials Science and Engineering*, 18(015011):1-19, 2010.
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1. R. Radis, G. A. Zickler, M. Stockinger, Ch. Sommitsch and E. Kozeschnik. Numerical Analysis of the Simultaneous Precipitation of  $\delta$  and  $\gamma'$  Phases in the Ni-Base Superalloy ATI Allvac<sup>®</sup>.718Plus<sup>™</sup>, In *7<sup>th</sup> International Symposium on Superalloy 718 & Derivatives*, Pittsburgh, Pennsylvania, USA, October 2010, in press.
2. R. Radis and E. Kozeschnik. Kinetics of AlN Precipitation in the Ferrite Phase Field of Microalloyed (HSLA) Steels. In *139<sup>th</sup> Annual Meeting & Exhibition TMS, Vol. 2: Materials Characterization, Computation, Modeling and Energy*, pages 341-348, Seattle, Washington, USA, February 2010.

3. R. Radis and E. Kozeschnik. Simulation of Concurrent AlN and VN Precipitation During Thermal Treatment of Microalloyed Steels. In *3<sup>rd</sup> International Conference on Simulation and Modelling of Metallurgical Processes in Steelmaking STEEL-SIM*, pages 326-331, Leoben, Austria, September 2009.
4. R. Radis, G. A. Zickler, M. Stockinger, C. Sommitsch and E. Kozeschnik. Interaction of the Precipitation Kinetics of  $\delta$  and  $\gamma'$  Phases in Nickel-Base Superalloy ATI Allvac<sup>®</sup>718Plus<sup>™</sup>. In *6<sup>th</sup> International Conference on Processing & Manufacturing of advanced Materials THERMEC*, pages 2712-2717, Berlin, Germany, August 2009.
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**Poster:**

1. R. Radis and E. Kozeschnik. Kinetics of AlN Precipitation in the Ferrite Phase Field of Microalloyed (HSLA) Steels. *139<sup>th</sup> Annual Meeting & Exhibition TMS*, Seattle, Washington, USA, February 2010.
2. R. Radis and E. Kozeschnik. Simulation of Concurrent AlN and VN Precipitation During Thermal Treatment of Microalloyed Steels. *3<sup>rd</sup> International Conference on Simulation and Modelling of Metallurgical Processes in Steelmaking STEEL-SIM*, Leoben, Austria, September 2009.
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4. S. Schwarz, R. Radis, E. Kozeschnik, G. Rumlplmair. TEM Investigations of the Precipitation Kinetics of Mn(Cu)S and AlN in Microalloyed Steel. *MICROSCOPY Conference 2009*, Graz, Austria, September 2009.

**Supervised Master Thesis:**

1. Ch. Schlacher, *Untersuchung der Versprödung von sprengverfestigten Manganhartstählen*, Graz University of Technology, March 2010.

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# Abstract

The present thesis deals with the numerical simulation of the precipitation kinetics of nitrides and carbides during isothermal annealing of (HSLA) microalloyed steel. Therefore the thermodynamic data of the relevant phases are analysed with special emphasis on microalloyed steel and implemented in the thermodynamic database 'mc\_steel'. Using the software package MatCalc, computer simulations of the precipitation kinetics of microalloy nitrides and carbides are performed. A newly developed model is employed, which takes into account the predominant precipitation of particles at grain boundaries, together with the classical treatment for randomly distributed precipitates in the grain interior. It is demonstrated, that the precipitation of nitrides and carbides can be modeled in a consistent way if simultaneous precipitation at grain boundaries and at dislocations is taken into account, dependent on chemical composition, grain size and annealing temperature. It is also demonstrated that, for consistent simulations, it is necessary to account for several physical mechanisms that are often neglected in this type of simulations, among them the precipitate/matrix volumetric misfit, the temperature dependent Young's modulus, composition-, temperature- and size-dependent interfacial energies, as well as the ratio between bulk and grain boundary diffusion. Finally, a new simulation approach is presented for the calculation of the precipitation kinetics of microalloying phases during casting. It is shown that segregation effects during solidification play an important role and must be considered when doing precipitation kinetics calculations. All precipitation kinetics simulations are compared to independent experimental results from literature.

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# Chapter 1

## Introduction

The term 'microalloyed steel' came up prior to 1940 when niobium and/or vanadium were added to a class of higher strength low C steels. However, the absence of a market and the lack of understanding of the use of these microalloying additions discouraged any serious developments. Both, research and development of these steels started simultaneously in the late 1950s and the early 1960s [1]. 40 years later, before the turn of the century, microalloyed steels have found extensive applications and it is estimated that they represent 12% of world steel production [2]. The addition of essentially less than 0.1% of the alloying element increases the yield strength of these steels by two or three times compared to plain C-Mn steel. Therefore microalloyed steels are used for the production of e.g. crankshafts, connecting rods, drive couplings, pipeline tubes, pipe fittings or pistonshafts [3]. Especially in the automotive sector, a strong increase in the application of these materials is observed. While the weight fraction of iron and steel in an average family car has decreased from 74% in 1978 to 67% in 1997, that of high- and medium strength steels has increased from 3.7% to 9.1% within the same time period [4].

In addition to a significant increase of yield strength, compared to plain C-Mn steels, these materials provide toughness, hot ductility and hot workability. Moreover, due to the decrease of the C content, good weldability is obtained. These desired mechanical properties are reached by the precipitation of second phase precipitates, e.g. nitrides, carbides or complex carbonitrides. Small amounts of vanadium (V), niobium (Nb), ti-

tanium (Ti) and aluminium (Al) are used together with nitrogen (N) and carbon (C) to form these second phase particles. During the production process of microalloyed steel, the precipitation will take place at different stages [5]. Type I precipitates are very stable ones, which form in the liquid phase and during or after solidification. These coarse particles have no significant influence on microstructural properties. Type II particles precipitate in austenite after solution treatment or during hot deformation. This group of precipitates is made responsible for grain refinement of microalloyed steel. During or after the austenite to ferrite transformation, particles of Type III develop, nucleating on the  $\gamma/\alpha$  interface and in ferrite. Usually, a fine precipitate dispersion is established in this stage, which contributes significantly to the strength by dispersion strengthening. For achieving desired mechanical properties, it is necessary to control the chemical composition of the alloys as well as the precipitation process during the production procedure, e.g. continuous casting, hot strip rolling, compact strip production, thin slab casting and rolling or coiling of hot strips of microalloyed steels. Therefore, a sound understanding of the basic mechanism of the precipitation process is essential, including the description of the precipitation kinetics in terms of nucleation, growth and coarsening of particles. Studying these nucleation and diffusion processes experimentally is expensive. With the development of specialised simulation software, e.g. the thermo-kinetic software MatCalc [6–8], it is possible, nowadays, to investigate these complex processes with computational methods on the researchers desktop. Using these tools enables the development of new alloys and/or optimization of the heat treatment parameters for the production processes of existing ones. This is crucial for being competitive on a global steel market.

# Chapter 2

## Objectives

The major goal of the present thesis is the numerical description of the precipitation kinetics of various nitrides and carbides in (HSLA) microalloyed steels. Therefore, different investigations are carried out.

Firstly, the solubility products of nitrides and carbides are studied in detail with special emphasis on microalloyed steel and compared to different thermodynamic assessments. Appropriate phase descriptions are determined and implemented in the thermodynamic database 'mc\_steel'.

Secondly, precipitation kinetics simulations are carried out, utilizing a newly developed model for predominant precipitation at grain boundaries together with the classical treatment of randomly distributed precipitates. Moreover, in addition to different nucleation sites, misfit strain energies between precipitate/matrix interfaces are taken into account. Therefore, the characteristics of both, the novel grain boundary model as well as the classical treatment of precipitation, are investigated in terms of parameter studies.

Thirdly, the precipitation kinetics of AlN and VN are calculated, considering simultaneous precipitation at grain boundaries and at dislocations. It is shown that, for a consistent description of AlN precipitation, it is necessary to account for several physical mechanisms that are often neglected in this type of simulations, among them the precipitate/matrix volumetric misfit, temperature dependent Young's modulus, composition-, temperature- and size-dependent interfacial energies, as well as the ratio between bulk

and grain boundary diffusion.

Finally, the precipitation processes of microalloying phases during casting are analysed, taking into account the microsegregation of alloying elements during solidification of microalloyed steel.

# Chapter 3

## Literature Review

### 3.1 Fundamentals of Microalloyed High Strength Low Alloy Steel

This section gives an overview of the fundamental characteristics of microalloyed steels. After a short historical review of the development, structure property relationships are discussed with special emphasis on the two main strengthening mechanism in microalloyed steels, fine grain hardening and precipitation strengthening.

#### 3.1.1 The Development of Microalloyed Steel

The positive effect of adding small amounts of microalloying elements was discovered prior to 1940. However, the absence of a sophisticated market and the lack of understanding of the use of this observation discouraged any serious developments. After an introduction in the late 1950s, both, research and development started simultaneously in the early 1960s [1]. Different stages of the development of microalloyed steels, including external influences, are shown in figure 3.1. The graph shows, that the development of these materials was driven by three main environmental influences. First of all, there was an increasing demand on steels with low C content due to the development and expansion of modern welding methods. Traditionally, tensile strength was increased by increasing the C and/or Mn content. But these materials obviously showed poor

weldability in terms of weld cracking. With increasing yield strain by the addition of microalloying elements, much lower C contents could be realized with improved weldability [1].

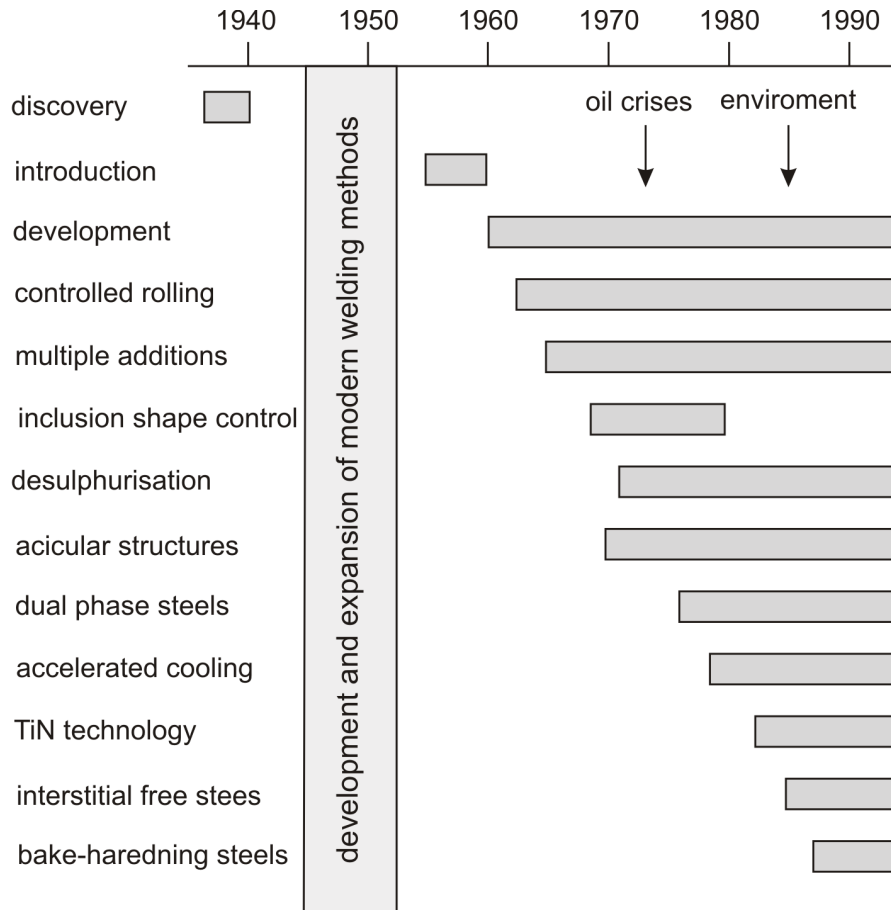


Figure 3.1: Historical review of the development of microalloyed steel with respect to major external influences [1].

Another major influence was the oil crises in 1973, which led to the development of lightweight vehicles with the aim of reducing the fuel consumption. Together with the substitution of lighter weight components by using aluminium and polymeric materials, microalloyed steels contributed significantly to vehicle weight reduction. And, last but not least, there was the ambition to reduce the environmental contamination and material conservation. Higher strength steels were substituted for other lower strength steels, the reduced material mass resulted in a substantially reduced materials requirement, with an attendant saving in energy required for iron extraction and steel refining [1].

### 3.1.2 Microstructure

Compared to classical high strength quenched and tempered martensitic steels, the microstructure of microalloyed (HSLA) steels consists of precipitation strengthened ferrite and pearlite containing simple stoichiometric nitrides and/or carbides as well as complex carbonitrides [3]. These desired microstructures are generally achieved during the production process, e.g. thermomechanical rolling. Thus, no further heat treatments are required to increase strength. The two major strengthening mechanism in microalloyed steels are fine grain hardening and dispersion strengthening [9]. While dispersion strengthening increases strength by hindering the movement of dislocations, fine grain hardening is made responsible for increasing strength, toughness, ductility and formability by controlling the grain size due to grain boundary pinning by second phase precipitates during the production process. Both mechanism are essential for achieving superior mechanical properties in microalloyed steels and are discussed subsequently.

### 3.1.3 Fine Grain Hardening

Grain refinement is extensively used as strengthening mechanism in the development of microalloyed steels. The quantitative relationship between grain size and yield strength was derived independently by Hall [10] and Petch [11] in the early 1950s. Today this relation is well established and known as the Hall-Petch relation,

$$\sigma_y = \sigma_0 + k_y d^{-1/2}, \quad (3.1)$$

where  $\sigma_y$  is the lower yield strength and  $d$  is the grain size. The quantities  $k_y$  and  $\sigma_0$  are material constants.  $k_y$  is dependent on various solutes in iron. While C, N, Ni and Si increase  $k_y$ , no influence of Cr is observed [12]. The quantity  $\sigma_0$  considers the starting stress for dislocation movement and the strengthening contribution of different solutes. Figure 3.2 shows the yield strength dependent on the final ferrite grain size for an air-cooled C-Mn steel containing 1.0% Mn, 0.25% Si and 0.01% N. In addition to the basic strength  $\sigma_0$ , which includes solid solution strengthening contributions of various



elements, a strong dependency on the final grain size is observed. While coarse grained materials show yield strengths in the range of 200 MPa, materials with approximately 5  $\mu\text{m}$  ferrite grain size double the yield strength to about 400 MPa.

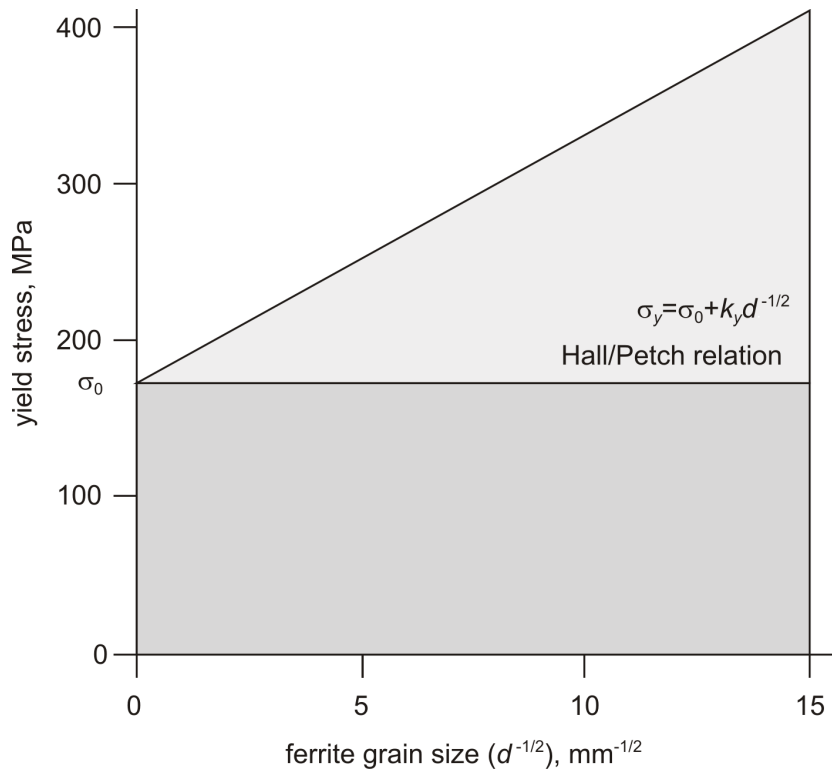


Figure 3.2: Dependency of yield stress on final ferrite grain size for an air-cooled C-Mn steel containing 1.0% Mn, 0.25% Si and 0.01% N [9].

In addition to the increase of yield strength, fine grain hardening improves mechanical properties like toughness, hot ductility and formability. Figure 3.3 shows the embrittlement vectors for various microstructural strengthening effects, defined by the change in the impact transition temperature per unit increase in yield strength. The vectors, depicted in this figure relate to ferrite-pearlite structures such as found in the matrix of low C structural steels [1]. The image clearly demonstrates that in addition to the positive strengthening effects (see Figure 3.2), grain refinement is the most effective way for reducing the impact transformation temperature and hence improving mechanical properties in terms of toughness or formability.

In microalloyed steel, this grain refinement effect is realized by pinning of grain bound-

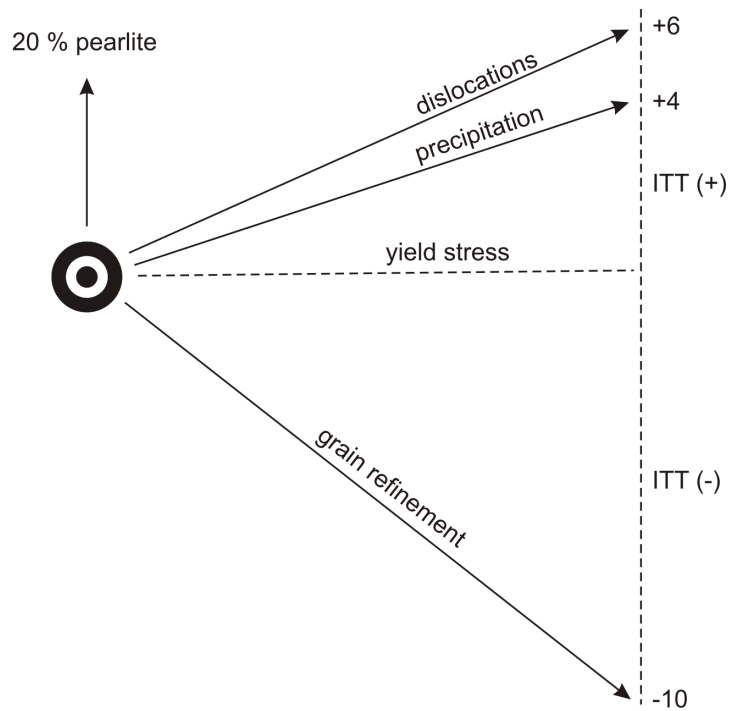


Figure 3.3: The embrittlement vectors, defined by the change in the impact transition temperature per unit increase in yield strength for various microstructural strengthening mechanism [1].

aries. This mechanism is well known for many years and essential for controlling grain growth, especially in the austenite phase field. Several theoretical descriptions of grain growth in particle containing materials have been developed, by Zener [13], Gladman [14], Heallman and Hillert [15], Elst et al. [16] and Rios [17]. The mathematical descriptions of these models as well as a comparison of the results with experimental data can be found in the work of Gao and Baker [18].

### 3.1.4 Dispersion Strengthening

Apart from fine grain hardening, dispersion strengthening is another very important strengthening mechanism in microalloyed steels. Shear stress induced dislocation movement in a dispersion hardened system, such as bcc ferrite, will cease when it meets one or several hard undeformable precipitates along its length. Two different dislocation particle interactions are made responsible for hindering the movement of dislocations, see

Figure 3.4. Depending on the precipitate size, smaller or larger than  $R_0$ , the dislocation will overcome the precipitate by cutting or bypassing it.  $R_0$  indicates the critical radius, where the increase of yield strength is maximised. Due to the hardness of the microalloy nitrides and carbides, the transition from a cutting mechanism to the bypassing mechanism will occur at rather small precipitate size [1]. Thus, the latter mechanism, called Orowan [19] type, is the major type of dispersion strengthening in these materials.

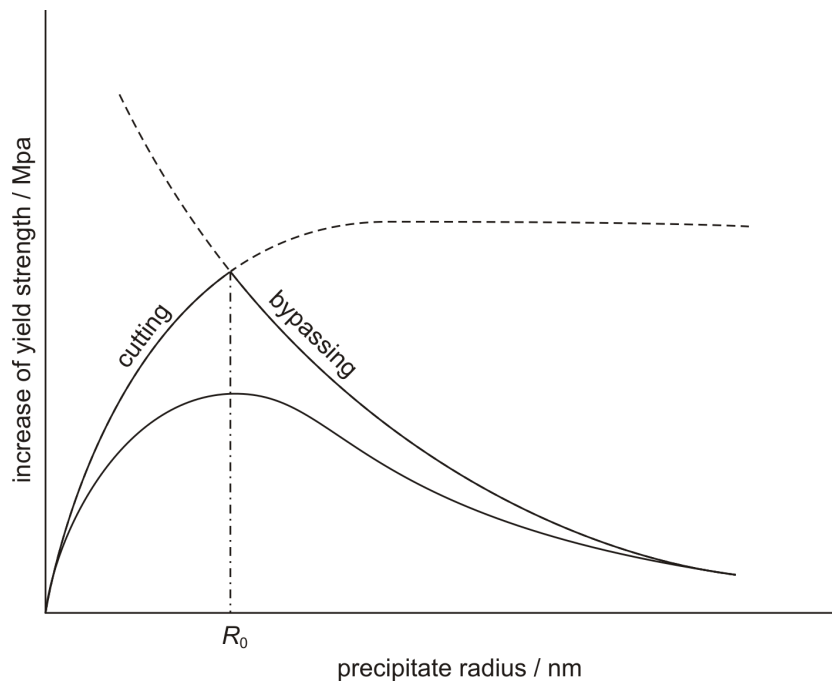


Figure 3.4: Relationship between particle radius and yield strength, considering cutting or bypassing of precipitates by dislocations [20].

This strengthening effect of hard carbides and/or nitrides can be described with theories based on dislocation bypassing. In the original work, Orowan developed his theory assuming a regular array of precipitates, which leads to an overestimation of the required stress for dislocation bypassing of the precipitates. Ashby [21] extended the work of Orowan taking into account that particles are arranged randomly as well as considering surface to surface interparticle distances rather than center to center spacing. According to that, the Ashby-Orowan [21] relationship is given with

$$\Delta\sigma_y = MCGb\frac{1}{\lambda} \ln\left(\frac{r_a}{r_i}\right), \quad (3.2)$$

where  $M$  is the Taylor factor,  $C$  is a constant,  $G$  is the shear modulus,  $b$  is the Burgers-vector,  $\lambda$  is the mean particle distance and  $r_i$  and  $r_a$  are the inner- and outer cut-off radii, respectively. Figure 3.5 shows the calculated increase of yield strength as a function of the volume fraction and diameter of precipitates according to equation 3.2. The graph clearly demonstrates the increase of yield strength with increasing volume fraction of small precipitates.

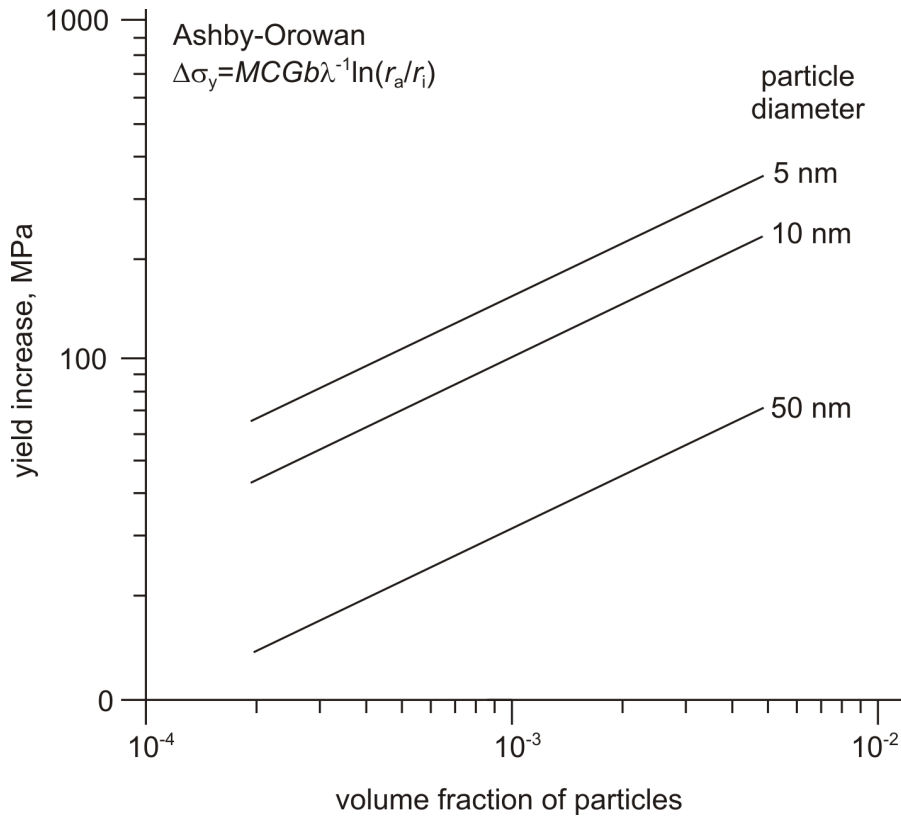


Figure 3.5: Calculated increase of yield strength by the addition of small volume fractions of particles of different sizes according to the Ashby-Orowan [21] relationships [1].

## 3.2 Precipitates in Microalloyed Steels

The precipitation of second phase particles, e.g. nitrides, carbides or complex carbonitrides, plays an important role in the production process of high-strength low-alloy (HSLA) steel. This section deals with the effect of various alloying elements on microalloyed steel with special emphasis on the precipitation of second phase particles. Furthermore, the influences of these second phase particles on the mechanical properties of microalloyed steel are presented.

### 3.2.1 Microalloying Elements and their Precipitates

The major microalloying elements in steel are Al, V, Nb and Ti. The content of these elements in microalloyed steel is rather low and generally below 0.1 wt%. These microalloying elements are added solely or in combination and precipitate in different matrix phases as simple carbides, nitrides or complex carbonitrides. Generally, microalloying additions have higher affinity to N compared to C, thus, all these elements prefer the precipitation in form of nitrides [1]. Since V, Nb and Ti have all fcc structure these elements have mutual solubility in each other, leading to the precipitation of complex carbonitrides.

#### 3.2.1.1 Aluminium

Al is often used for deoxidation during casting of Al killed steels. Therefore, the use of Al as a microalloying element can be seen as byproduct which is free of charge. However, important mechanical and technological properties, such as hot ductility [22–24], impact toughness [24], deep drawability [25] or weldability [26] can be improved by microalloying additions of Al. On the other hand, precipitation of AlN can cause embrittlement and induce cracking phenomena, such as, e.g. rock candy fracture [27–33]. Some authors [34] report that AlN without any other trace element has no influence on hot ductility of high purity iron. A summary of the effects of AlN in steel can be found in the extensive review of Wilson and Gladman [26].

The most significant effect of AlN in steel is in grain size control [24, 35–37], which

directly influences these properties. Especially in the austenite phase field, AlN precipitates predominately at austenite grain boundaries [38–43]. Therefore, in austenite, grain growth can be effectively controlled by AlN precipitation. For instance, Militzer et al. [44] show that, in the presence of AlN precipitation, there is very little grain growth (40  $\mu\text{m}$  after 10 min.) during isothermal annealing at temperatures below 1100 °C, whereas there is fast grain growth at 1150 °C (200  $\mu\text{m}$  after 10 min.). They suggested that complete dissolution of AlN above 1100 °C is responsible for the large grain size in their steel. And also Gao and Baker [18] claim that the presence of AlN precipitates is responsible for grain size control in Al-V-N steel. In addition to AlN precipitates at the austenite grain boundaries, AlN is also found at ferrite grain boundaries [41].

The equilibrium crystallographic structure of AlN is the hexagonal wurtzite structure with a lattice constant of  $a=0.311$  nm and  $c=0.4978$  nm. At the early stages of precipitation, also, simple cubic structure with  $a=0.412$  nm is observed [43,45–47]. A detailed investigation of the two crystallographic structures can be found in reference [48].

Depending on the heat treatment and processing conditions, the morphology of AlN precipitates cover a wide range of different precipitate shapes. AlN particles can be dendritic (up to several hundred microns in length), large plates, rod, needle shaped, rectangular/cuboidal or simple finley distributed within the matrix ( $< 5$  nm). A large number of references summarizing these different morphologies, can be found in reference [26].

It should be noted that Al differs from other microalloying elements, like V, Nb and Ti, since it does not form a carbide in steel. Due to the hexagonal wurtzite structure of AlN precipitates, there is no mutual solubility with the carbides/nitrides of V, Nb and Ti, which all have NaCl cubic crystal structure [49].

### 3.2.1.2 Vanadium

World-wide, 94 % of the total V is used for alloying of steel and cast iron [50]. Compared to Al, the positive influences of V are not related to grain size control but to dispersion strengthening effects due to precipitation of finley distributed particles in a ferritic matrix (see section 3.1.4). Li et al. [51] showed that the dispersion plus dis-

location strengthening in V-microalloyed steels is in the range of 80 - 250 MPa, and Lagneborg et al. [49] concluded in their extensive review, that in a 0.12 % V-steel, the precipitation hardening potential is typically in the range of 6 MPa per 0.001 % N.

Therefore, V microalloyed steels are widely used for forging steels, providing high strength as well as good fatigue resistance, which is essential for applications for structures, machine and vehicle parts [3]. Furthermore V strengthened cold rolled steels are used in the automotive sector. According to Scott et al. [52], the ideal microalloying element for cold rolled steels should have following properties:

- Large precipitates formed in slabs can be redissolved during reheating
- No significant increase in hardness during hot rolling and no decrease in hot ductility
- No precipitation during finishing and coiling
- Intense precipitation during continuous annealing
- Precipitation should occur predominantly in the ferrite phase field
- Precipitation should be intergranular
- The number density should be as high as possible

These points require that the precipitating phase is sufficiently soluble in austenite, but shows a maximum possible precipitation driving force to generate a high particle nucleation density in ferrite. Considering all above mentioned criteria, the microalloying element which best satisfies these constraints is V [52]. In addition to these excellent properties for the production of cold rolled sheets, V provides also excellent metallurgical properties for production of hot rolled sheet products. Therefore, the important aspects can be summarized as, the high solubility of V(C,N) in austenite, the low solute drag coefficient [53] of V in austenite and the fact that N is the preferred element in V(C,N) precipitation [54].

### 3.2.1.3 Niobium

Nb has become a preferred microalloying element in HSLA steels sheet, since it increases the strength without reducing toughness. Therefore, the predominate strengthening mechanism is grain refinement of the ferrite grains, whereas other microalloying additions apply the strengthening mechanism dispersion strengthening to a major extent (Ti) or completely (V). Furthermore, compared to Nb, higher amounts of Ti or V are needed to obtain the same increase in strength [55,56]. Consequently, Nb is widely used in civil construction, automobile and line pipe applications. Plate sheets are produced by applying thermomechanical rolling, a technique that simultaneously increases strength and toughness by grain refinement. The most important role of Nb as a microalloying element in thermomechanically rolled steel is the retardation of austenite recrystallization, which provides more nuclei for the  $\alpha/\gamma$  transformation and thus a finer ferrite grain size [57]. Two reasons can be attributed to this important effect [58]. Firstly, Nb in solid solution can retard the austenite recrystallization after deformation through the solute drag [53] effect as a consequence of its relatively large atom size. Secondly, and more important, Nb can also retard the recrystallization process by the precipitation of finely dispersed NbC particles. Depending on the hot rolling parameters, like deformation, temperature and interpass time, NbC precipitation is generally incomplete in the austenite phase field and does not reach its thermodynamical equilibrium state. Consequently, part of Nb retained in solid solution may form nanometer sized particles during or after the  $\gamma/\alpha$  transition, which increases strength by precipitation hardening.

### 3.2.1.4 Titanium

Ti is often used to stabilize the N in liquid steel and it is well established that Ti has a very low solubility in steel. Therefore, it forms nitrides during and/or early after solidification in Ti-V and Ti-Nb microalloyed steels, providing a significant contribution in grain growth inhibition [49], see also figure 3.7. Furthermore it is also well known that there is a significant contribution of precipitation strengthening by the precipitation of fine dispersed TiN to the matrix strength of ferrite in microalloyed steel [59]. And,



especially, in Ti containing microalloyed steels, precipitation of complex carbonitrides with inhomogeneous element distribution is observed. The particles, which precipitate in the high temperature range, act as nucleation sites for nitrides and/or carbides with higher solubility. Thus, core/shell constructions of precipitates are observed with a TiN rich core surrounded by other microalloying elements, e.g. V, Nb(C,N) [60–62]. The morphology of pure TiN or Ti-rich complex carbonitrides are plates and cuboids, see e.g. [63,64]. TiN is mainly used for the production of high strength steels sheets, e.g. in thin strip produced by compact strip production or hot strip rolling.

### 3.2.2 Effects of Microalloying on the Properties of Microalloyed Steels

The previous sections demonstrated the different characteristics of the microalloying additions. This section now deals with the influence of different precipitate phases on the mechanical properties of microalloyed steels in terms of fine grain hardening (section 3.1.3) and dispersion strengthening (section 3.1.4). Figure 3.6 shows a comparison of the effects of classical microalloying elements on the properties of hot strips, after Hulka [55]. The x-axis represents the amount of microalloying additions, whereas the y-axis shows the influences on the mechanical properties, strength and toughness. The letters P and F indicate the increase or decrease of the mechanical properties due to precipitation strengthening (P) and/or fine grain hardening (F). Both, Nb as well as Ti show big influences on the strength caused by both strengthening mechanism (P and F), whereas V is a solely precipitation strengthening element. Compared to Ti and V, there are much less additions of Nb required to achieve the same increase in strength. Furthermore, the graph shows that Nb is the only element enhancing considerable the toughness by fine grain hardening. However, Hulka [55] neglected in his comparison the positive influence of AlN on toughness.

Speich et al. [65] also compared the positive effects of microalloying additions on fine grain hardening. They suggested, that, next to Ti additions, Al is maybe the most effective element in avoiding grain coarsening at high temperatures, see figure 3.7. Con-

sequently the precipitation of AlN provides positive influences on toughness and embrittlement.

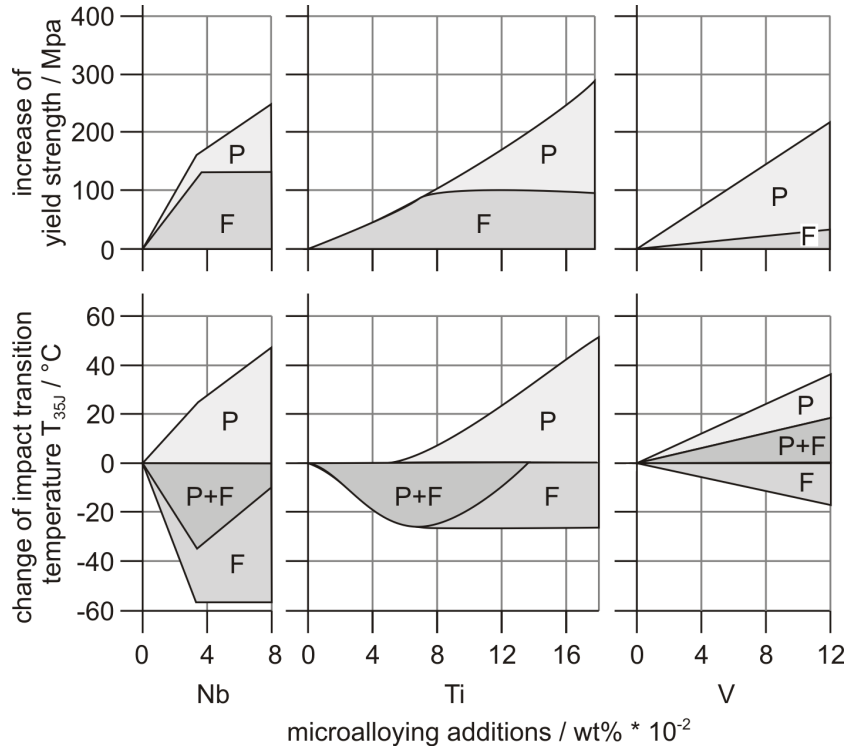


Figure 3.6: Effect of microalloying on the properties of hot strip [55].

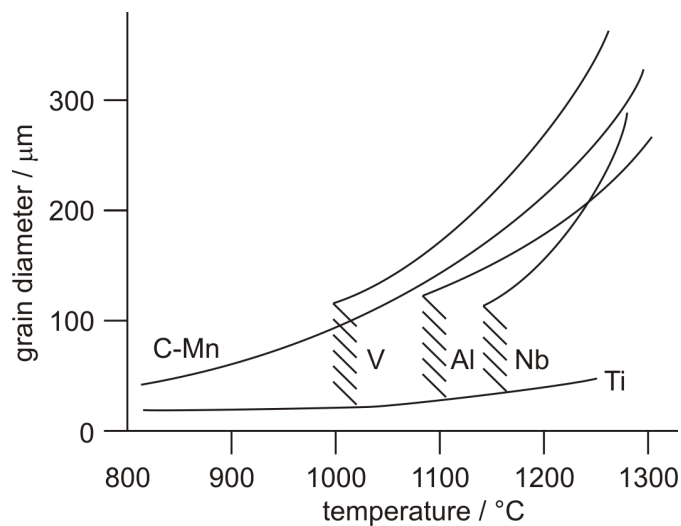


Figure 3.7: Austenite grain growth characteristics in steels containing various microalloying additions [65].

Another comparison of the microalloying elements was made by Morrison [2] with special emphasis on the achievement of different desired microstructures during various production processes, see table 3.1.

Table 3.1: Roles of the classical microalloying elements V, Nb and Ti, after Morrison [2].

microalloying element	V	Nb	Ti
precipitation strengthening after hot rolling	VN, VC	NbCN	TiC
precipitation strengthening after normalising	VC	-	-
influences recrystallisation during hot rolling	-	Nb, NbCN	-
refines grain size on normalising	VN	NbCN	TiC
refines grain size during high temperature austenitising	-	-	TiN
influences transformation characteristics after hot rolling	-	Nb	-

Morrison [2] also neglected the important role of Al in his summary. However, previous sections of this work as well as detailed investigations of other authors, e. g. Wilson and Gladman [26], show the high significance of this element, providing fine grain hardening and hence, enhanced toughness properties.

### 3.2.3 The Solubility of Nitrides and Carbides in Steel

In order to control the strengthening mechanisms, dispersion strengthening and fine grain hardening, it is necessary to understand the dissolution and precipitation processes. Fine ferrite grain sizes are generally achieved by controlling the austenite grain size in the high temperature range prior to the  $\gamma/\alpha$  transformation. Thus, the precipitation of particles at sub-solvus temperatures as well as particles remaining undissolved in the austenite phase field are required. On the other hand, the precipitation of fine dispersed particles (3-5 nm in diameter) in the ferrite phase field after transformation is essential to increase strength by dispersion hardening. Therefore, such nitrides and/or carbides should be dissolved in the austenite region. In order to achieve the desired metallurgical

states, a detailed knowledge of the solubilities of the microalloy carbides and nitrides is required, together with a knowledge of their precipitation behavior [1].

### 3.2.3.1 The Solubility Product

The reaction between a microalloying element  $M$  with an interstitial  $X$  to give a compound  $MX$  at a given temperature  $T$  is given with



where  $[M]$  and  $[X]$  represent the concentrations of the microalloying element and the interstitial, dissolved in the matrix, whereas  $(MX)$  stands for the compound. The equilibrium constant, also known as the solubility product can be written as

$$k_s = \frac{a_{[M]} \cdot a_{[X]}}{a_{(MX)}}, \quad (3.4)$$

where  $a_{[M]}$  and  $a_{[X]}$  represent the activity of the dissolved elements and  $a_{(MX)}$  is the activity of the compound. These activities can be expressed by the products of the elemental concentrations and the activity coefficients  $\gamma$  leading to

$$a_{[M]} = \gamma_{[M]}[M], \quad (3.5)$$

and hence

$$k_s = \frac{\gamma_{[M]}[M] \cdot \gamma_{[X]}[X]}{a_{(MX)}}. \quad (3.6)$$

In a rather empirical manner, it is assumed that the activity coefficients  $\gamma_{[M]}$  and  $\gamma_{[X]}$  are equal to unity and of course the activity of the product  $(MX)$  is unity [1]. Thus it can be written

$$k_s = [M][X]. \quad (3.7)$$

$[M]$  and  $[X]$  should be the molare fractions. Since the concentrations are very low, traditionally, these element contents are expressed in weight percentage, which has only minor influence on the value of  $k_s$ .

The temperature dependence can be expressed using an Arrhenius relationship

$$\ln k_s = - \left( \frac{Q}{RT} \right) + C, \quad (3.8)$$

where  $R$  represents the universal gas constant,  $Q$  the heat of dissolution,  $T$  the temperature and  $C$  a constant. Historically, metallurgist used the equation

$$\log k_s = A - \frac{B}{T}, \quad (3.9)$$

where  $A$  and  $B$  are system constants, and  $T$  stands for the temperature in  $K$ . Combining equations 3.7 and 3.9 leads to the universal definition of the solubility product

$$\log k_s = \log[M[wt\%]] \cdot [X[wt\%]] = A - \frac{B}{T}. \quad (3.10)$$

Additional information on the deduction of the solubility product can be found in reference [1].

### 3.2.3.2 Solubility Products of Various Stoichiometric Phases

Different microalloying phases show a wide range of different solubility products. Figure 3.8 shows a comparison of the solubility data of different microalloy nitrides and carbides in austenite [49], since this is the temperature range of prime importance, when doing precipitation kinetics calculations.

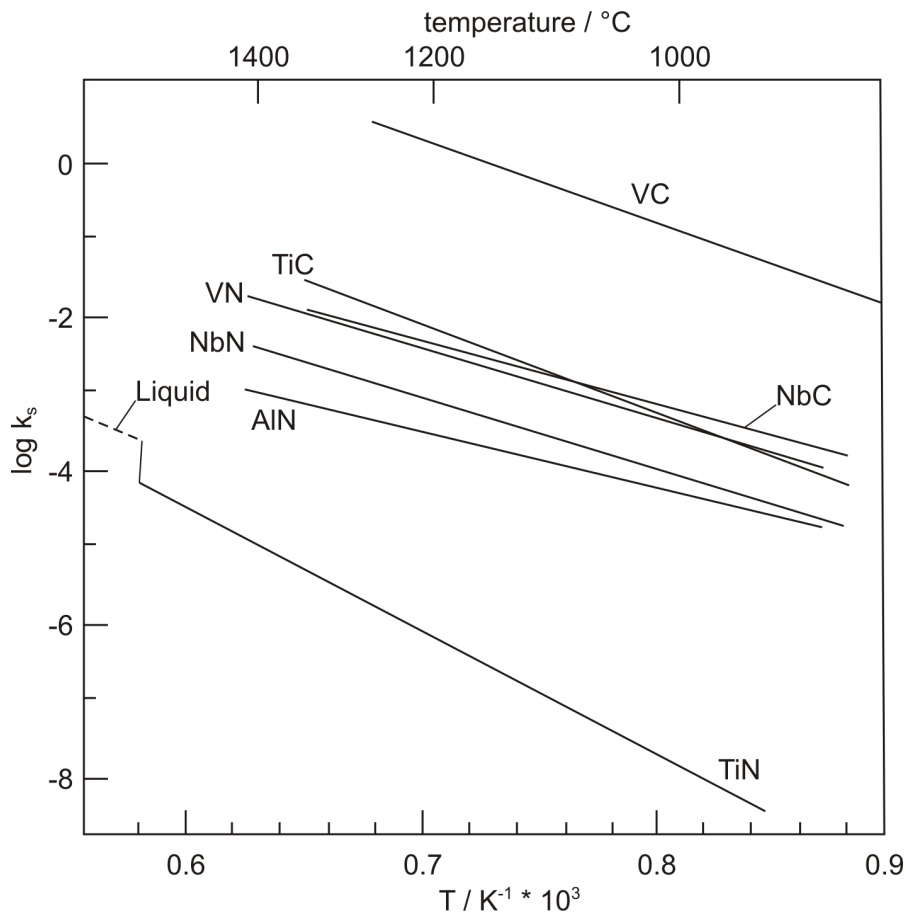


Figure 3.8: Comparison of the solubility products of the microalloy carbides and nitrides, in austenite [49].

Despite the variations in cited solubility product equations summarized in section 4.2 Gladman [1] proposed several important conclusions.

- In the austenite phase field, all microalloyed nitrides are more stable than their corresponding carbide. Large differences in solubility between the microalloy nitride and carbide are observed for the elements V and Ti, whilst the difference between the solubilities of NbN and NbC is much less. The solubility of AIN is slightly less compared to the solubility of NbN. However, AIN has never been reported to occur in steel, even in steel with considerable Al and C contents.
- Most of the microalloying carbides and nitrides show broadly similar solubility in austenite, with two exceptions. TiN is more stable by a factor of 1000, and VC is

far more soluble by a factor of 1000, than other microalloy carbides and nitrides.

- The solubility of TiN in liquid at temperature up to 1600 °C is similar to the solubility of most other microalloy nitrides and carbides in austenite at around 1200 °C.

The knowledge of these individual solubilities of microalloy compounds give a clear direction for the selection of microalloying elements for special applications. While the high solubility of VC suggests the use of V in normalised steels and high C steels, Ti and Al are the preferred grain refining elements in the austenite phase field as a direct consequence of their higher phase stability.

### 3.3 Theory of Precipitation

Precipitation is a special case of phase transformation where a new phase grows on the expense of an existing one. If we are talking from precipitates, we talk from particles, which are very small (few nanometers to few micrometers) compared to the parent phase (matrix phase). Typical precipitates are intermetallic phases or oxides, sulphides, nitrides and carbides [66]. The present work deals exclusively with the precipitation of nitrides and carbides, which are the essential precipitates in microalloyed steel.

In general, precipitation occurs when the amount of alloying elements exceeds the soluble equilibrium amount of elements in an alloying system. Apart from the dependence on chemical composition and temperature these precipitation processes are strongly time dependent, which is why we are talking of kinetics. The precipitation process can be divided into three different stages, compare for instance [66]:

- **Nucleation:** The initial stages in precipitate life. Stochastic process driven by microscopic thermal and compositional fluctuations.
- **Growth:** Controlled attachment of atoms to the precipitate. Deterministic process driven by chemical and/or mechanical driving forces.
- **Coarsening:** Dissolution of small precipitates in favor of larger ones. Driven by curvature-induced pressure, also called the Gibbs-Thomson effect.

Figure 3.9 shows the volume fraction  $f_T$ , the precipitate diameter  $d_T$  and the interparticle spacing  $\lambda_T$  versus time, indicating the three different stages, nucleation, growth and coarsening, during isothermal annealing at a given temperature  $T$ . After the incubation time  $\tau$ , precipitates are generated with a critical precipitate diameter  $d^*$ . After this nucleation process, the particles pass through the stage of growth. The volume fraction  $f_T$  increases along with the increase of the particle diameter  $d_T$ . In the special case of randomly distributed precipitates, which is depicted in figure 3.9, the increase of the diameter goes with  $t^{\frac{1}{2}}$ . When the equilibrium phase fraction  $f_T$  is reached, stage three begins. Big precipitates grow at the expense of smaller ones, while the volume fraction



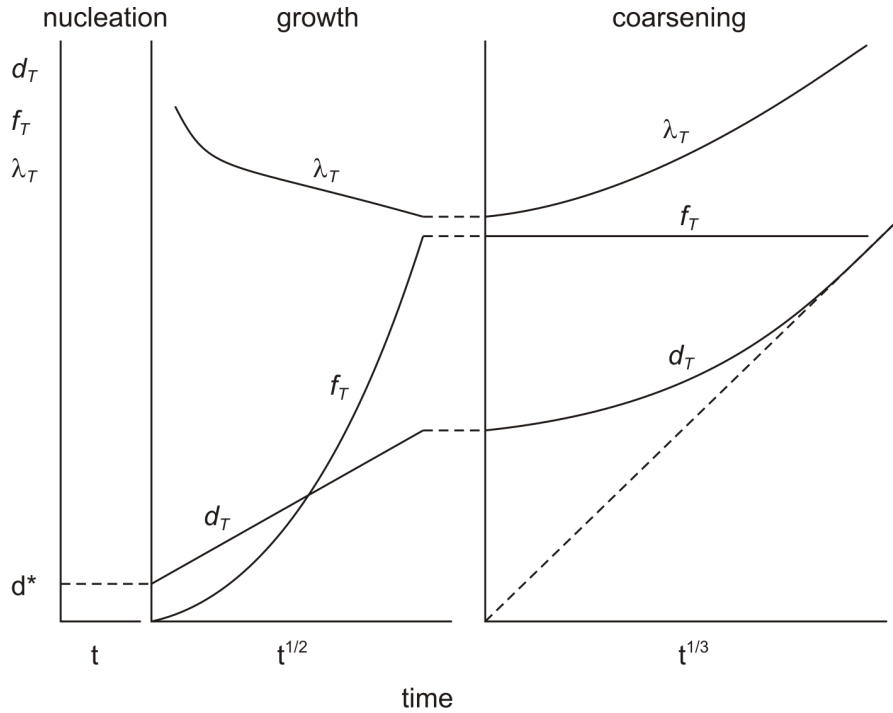


Figure 3.9: Development of precipitate diameter  $d_T$ , interparticle distance  $\lambda_T$  and phase fraction  $f_T$  versus time during isothermal annealing at a given temperature  $T$  [67].

$f_T$  remains nearly constant. This stage is known as the coarsening stage and the precipitate diameter obeys  $t^{1/3}$  if the particles are distributed randomly in the matrix phase. It is worthy to note here again that the parabolic growth law as well as the cubic coarsening law, describing the development of the mean radius during isothermal annealing, are only fulfilled in the special case of randomly distributed precipitates. This is the case for homogeneous precipitation in the bulk and in the broadest sense for heterogeneous precipitation at one or two dimensional lattice defects (e.g. dislocations), due to their nearly random arrangement. These laws are clearly not fulfilled when precipitation occurs predominantly at grain boundaries, which is demonstrated later on in the present work.

## 3.4 Precipitation Kinetics Models Embedded in the Software MatCalc

MatCalc [6–8] is a thermo-kinetic software package developed for the simulation of the precipitation kinetics of various metallurgical phases. The characteristics of this program can be summarized as a mean-field treatment of nucleation, growth and coarsening of multi-component, multi-phase and multi particle systems. This approach is based on the Onsager thermodynamic extremal principle [68,69], assuming that a thermodynamic system evolves along the particular kinetic path, where maximum entropy is produced. The theory is described by Svoboda et al. [6] and the treatment of the numerical implementation in the software package MatCalc was done by Kozeschnik et al. [7]. The basis for the calculation of the thermodynamics of the system are CALPHAD type databases [70], providing chemical potentials of the phases as well as diffusivities of the alloying elements on the basis of the chemical composition of the alloy and the present temperature. The following sections give a brief introduction on the basic models incorporated in the software package MatCalc.

### 3.4.1 Classical Nucleation Theory for Multi-Component Systems

In the software package MatCalc [6–8], the nucleation kinetics of precipitates is calculated from Classical Nucleation Theory (CNT) [71] extended for multi-component systems [6,66,72]. Accordingly, the transient nucleation rate  $J$  is given as

$$J = N_0 Z \beta^* \cdot \exp\left(\frac{-G^*}{k \cdot T}\right) \cdot \exp\left(\frac{-\tau}{t}\right). \quad (3.11)$$

$J$  describes the rate, at which nuclei are created per unit volume and time.  $N_0$  represents the total number of potential nucleation sites. The Zeldovich factor  $Z$  takes into account that the nucleus is destabilized by thermal excitation as compared to the inactivated state, and it is given with [71]

$$Z = \left[ \frac{-1}{2\pi kT} \frac{\partial^2 \Delta G}{\partial n^2} \Big|_{n^*} \right]^{\frac{1}{2}}. \quad (3.12)$$

$\Delta G$  represents the total free energy change on nucleus formation,  $n$  is the number of atoms in the nucleus,  $k$  is the Boltzmann constant and  $T$  is temperature. The derivative of  $\Delta G$  is taken at the critical nucleus size  $n^*$ .

The atomic attachment rate  $\beta^*$  takes into account the long-range diffusive transport of atoms, which is needed for nucleus formation if the chemical composition of the matrix is different from the chemical composition of the precipitate. A suitable multi-component equation for  $\beta^*$  has been derived by Svoboda et al. [6], and it is given with

$$\beta^* = \frac{4\pi\rho^{*2}}{a^4\Omega} \left[ \sum_{i=1}^n \frac{(c_{ki} - c_{0i})^2}{c_{0i}D_{0i}} \right]^{-1}, \quad (3.13)$$

where  $\rho^*$  is the critical nucleation radius,  $a$  is the atomic distance,  $\Omega$  is the molar volume,  $c_{ki}$  and  $c_{0i}$  are the element concentrations in precipitate with index  $k$  and matrix, respectively, and  $D_{0i}$  is the diffusivity of element  $i$  in the matrix. The incubation time  $\tau$  is given with [71]

$$\tau = \frac{1}{2\beta^*Z^2}, \quad (3.14)$$

and the critical energy for nucleus formation  $G^*$  is

$$G^* = \frac{16\pi}{3} \frac{\gamma^3}{(\Delta G_{vol} - \Delta G_s)^2}, \quad (3.15)$$

with the interfacial energy  $\gamma$ , the volume free energy change  $\Delta G_{vol}$  and the misfit strain energy change  $\Delta G_s$ . The interfacial energy  $\gamma$  is calculated using a generalized  $n$ -next nearest-neighbour broken-bond approach [73], assuming a planar and sharp interface, see also section 3.4.3. Corrections for the small size of the precipitate are taken into account according to reference [74].

It is important to note that  $G^*$  is a most essential quantity in nucleation theory, when compared to the other quantities occurring in equation 3.11. Since  $G^*$  appears in the exponent of the nucleation rate equation 3.11, small changes in  $\gamma$  and/or  $\Delta G_{vol} - \Delta G_s$  can lead to huge variations in  $J$ , which has been demonstrated recently for  $\gamma'$ -precipitation in Ni-base superalloy UDIMET 720 Li [75–77].

Finally, the misfit strain energy is given with

$$\Delta G_s = \frac{E}{1 - \nu} \epsilon^2, \quad (3.16)$$

where  $E$  is the Young's modulus and  $\nu$  the Poisson constant of the matrix.  $\epsilon$  is the linear misfit strain, approximately given as

$$\epsilon = \frac{1}{3} \frac{V_{mol}^P - V_{mol}^M}{V_{mol}^M}, \quad (3.17)$$

where  $V_{mol}^P$  and  $V_{mol}^M$  are the molar volumes of the precipitate and the matrix, respectively.

### 3.4.2 Growth and Coarsening - The SFFK Model

Once a precipitate is nucleated, its further growth is evaluated based on the evolution equations for the radius and composition of the precipitate derived by Svoboda et al. [6] in a mean-field approach utilizing the thermodynamic extremal principle [68, 69], also known as the SFFK model. The next two sections give a brief overview of this approach for randomly distributed precipitates (section 3.4.2.1) and precipitates arranged at grain boundaries (section 3.4.2.2).

#### 3.4.2.1 Randomly Distributed Precipitates

The SFFK model [6] enables the treatment of multi-component, multi-phase and multi-particle systems. Figure 3.10 shows a schematic precipitate distribution for random

precipitation of particles with different size and chemical composition, surrounded by spherical diffusion fields (shaded areas).

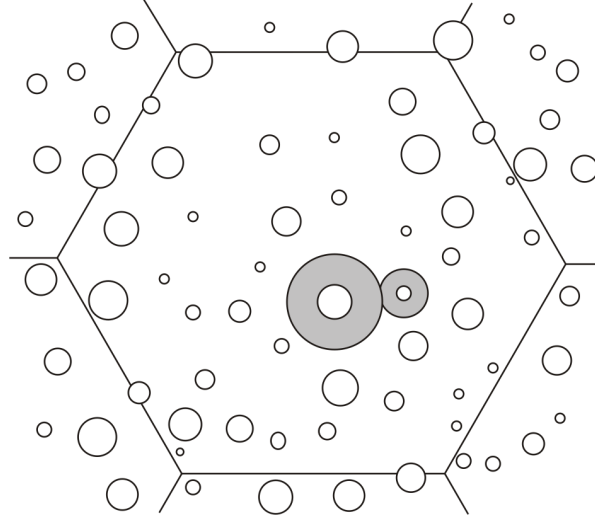


Figure 3.10: Schematic precipitate distributions and diffusion fields (shaded areas) for random precipitation in 2D [78].

In a system containing  $n$  chemical elements and  $m$  randomly distributed precipitates with spherical diffusion fields, the total free energy of the system can be written as [6,8]

$$G = \sum_{i=1}^n N_{0i} \mu_{0i} + \sum_{k=1}^m \frac{4\pi \rho_k^3}{3} \left( \lambda_k + \sum_{i=1}^n c_{ki} \mu_{ki} \right) + \sum_{k=1}^m 4\pi \rho_k^2 \gamma_k, \quad (3.18)$$

where  $\mu_{0i}$  ( $i = 1, \dots, n$ ) is the chemical potential of component  $i$  in the matrix and  $\mu_{ki}$  ( $k = 1, \dots, m, i = 1, \dots, n$ ) is the chemical potential of component  $i$  in the precipitate  $k$ . The chemical potentials are expressed as functions of the concentrations  $c_{ki}$  ( $k = 0, \dots, m, i = 1, \dots, n$ ).  $\gamma$  is the interface energy density and  $\lambda_k$  takes into account the contribution of the elastic energy and plastic work due to the volume change of precipitates. Therefore, equation 3.18 contains only mean values and no concentration profiles are considered. Thus, the SFFK model is a classical approximation procedure and can be seen as a 'mean-field approach'.

In the formulation of the total free energy (equation 3.18), each of the three terms can be related to a particular energy contribution [66]. Accordingly,

- The first term describes the Gibbs energy of the matrix by summing up the product of all moles of components  $N_{0i}$  with the corresponding chemical potentials  $\mu_{0i}$ .
- The second term is related to the sum of the bulk free energy of all  $m$  precipitates, where the specific quantities for the mechanical free energy contribution  $\lambda_k$  and the chemical free energy  $\sum_{i=1}^n c_{ki}\mu_{ki}$  are multiplied with the precipitate volume  $4/3\pi\rho_k^3$ . The summation over  $k$  in the second term is performed over all  $m$  precipitates.
- The third term takes into account the energy contribution of the precipitate-matrix interfaces. The specific interfacial energy  $\gamma_k$  is multiplied by the precipitate surface area  $4\pi\rho_k^2$  and summation is performed over all  $m$  precipitates.

According to equation 3.18, the formulation of the total free energy of the system in the present state, is defined by the independent state parameters  $\rho_k$  and  $c_{ki}$ . The other parameters can be determined, using the mass conservation law for each component  $i$  [66],

$$N_i = N_{0i} + \sum_{k=1}^m \frac{4\pi\rho_k^3}{3} c_{ki} \quad (3.19)$$

and the global mass conservation

$$\sum_{i=1}^n N_i = N. \quad (3.20)$$

If a thermodynamic system evolves to a more stable state, the difference in free energy between initial and the final stage is dissipated. The classical products of this reaction are entropy and transformation heat. Three different dissipation effects are considered in the SFFK model [66]:

- $Q_1$  - Dissipation due to interface movement
- $Q_2$  - Dissipation due to diffusion inside the precipitate
- $Q_3$  - Dissipation due to diffusion inside the matrix

**Dissipation due to interface movement** can be related to the resistance against the driving pressure for interface migration and can be written as

$$Q_1 = \sum_{k=1}^m \frac{4\pi\rho_k^2}{M_k} \dot{\rho}_k^2, \quad (3.21)$$

where  $M_k$  stands for the interface mobility.

**Dissipation due to diffusion inside the precipitate** is given with

$$Q_2 = \sum_{k=1}^m \sum_{i=1}^n \int_0^{\rho_k} \frac{RT}{c_{ki}D_{ki}} 4\pi r^2 j_{ki}^2 dr, \quad (3.22)$$

where  $R$  is the universal gas constant.  $j_{ki}$  and  $D_{ki}$  are the flux and the diffusion coefficient of component  $i$  in the precipitate  $k$ .

**Dissipation due to diffusion inside the matrix** can be expressed as

$$Q_3 = \sum_{k=1}^m \sum_{i=1}^n \int_{\rho_k}^Z \frac{RT}{c_{0i}D_{0i}} 4\pi r^2 J_{ki}^2 dr, \quad (3.23)$$

assuming that the distance between the individual precipitates is sufficiently large and their diffusion profiles do not overlap. Thereby,  $Z$  represents a characteristic length between two precipitates. The flux  $J_{ki}$  can be determined using the mass conservation law across the interface.

Finally, the total rate of dissipation is calculated by summing up the expressions of the individual contributions

$$Q = Q_1 + Q_2 + Q_3. \quad (3.24)$$

Applying the thermodynamic extremal principle after Onsager [6], the growth rate  $\dot{\rho}_k$  and the rate of change of chemical composition  $\dot{c}_{ki}$  of the precipitate with index  $k$  are obtained from solution of the linear system of equations

$$A_{ij}y_j = B_i, \quad (3.25)$$

where the variable  $y_j$  represents the rates  $\dot{\rho}_k$  and  $\dot{c}_{ki}$ , as well as the Lagrange multipliers from the stoichiometry constraints (see reference [6]). The index  $i$  comprises the variables for radius and chemical composition of the precipitate as well as the stoichiometry boundary conditions implied from the crystal structure of the precipitate phase. The system of equations 3.25 is solved for each precipitate  $k$  separately. The full expressions for the coefficients  $A_{ij}$  and  $B_i$ , as used in the present work, are given in reference [8]. The numerical time integration of  $\dot{\rho}_k$  and  $\dot{c}_{ki}$  is performed in the software MatCalc, based on the classical numerical Kampmann - Wagner approach [79], see section 3.4.4.

### 3.4.2.2 Precipitates Arranged at Grain Boundaries

If precipitation occurs predominantly at grain boundaries, the assumption of randomly distributed precipitates with spherical diffusion fields is not valid anymore. There are two reasons for the predominant precipitation of particles at grain boundaries. Firstly, grain boundaries are preferred nucleation sites due to less misfit strain energies, compared to nucleation sites with larger precipitate-matrix lattice mismatch. Secondly, grain boundaries represent very efficient short circuit diffusion paths with diffusivities several orders of magnitude larger than in the bulk [80]. Thus, different diffusion coefficients  $D_{gi}$  and  $D_{bi}$  must be considered, representing the diffusivities of component  $i$  at grain boundaries and in the bulk, respectively.

Figure 3.11 shows a schematic arrangement of precipitates aligned at grain boundaries. Two processes are assumed to occur parallel in the system [78]:

- Growth of precipitates on the account of components dissolved in the matrix
- Coarsening of precipitates by means of redistribution of the fixed total volume of the grain boundary precipitates amongst precipitates of different sizes

According to this image, the model takes into account two different diffusion geometries,



indicated by arrows. On the one hand, the slow diffusion from the grain interior towards the grain boundary is considered. On the other hand, the model takes into account the many orders of magnitude faster diffusion along the grain boundaries [78].

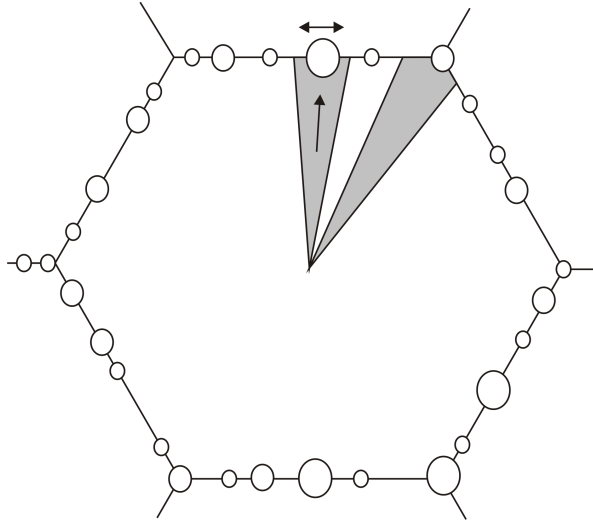


Figure 3.11: Schematic precipitate distributions for heterogeneous precipitation at grain boundaries in 2D, indicating the diffusional zone in the bulk and in the grain boundary (arrows) [78].

According to the treatment of randomly distributed precipitates with spherical diffusion fields, one can find the formulation of the total Gibbs free energy of the system as well as the corresponding dissipation equations. In the case of predominant precipitation at grain boundaries, two different dissipation mechanisms are considered in the model [78]:

- $Q^g$  - Dissipation due to diffusion in the grain boundary
- $Q^b$  - Dissipation due to diffusion in the grain

By applying the thermodynamic extremal principle after Onsager, the growth rate  $\dot{\rho}_k^g$  and the rate of coarsening  $\dot{\rho}_k^c$  of each precipitate with index  $k$  can be determined, similar to the approach for randomly distributed precipitates. The corresponding equations can be found in reference [78].

### 3.4.3 Calculation of the Interfacial Energy

Apart from Gibbs energies and diffusion rates, the interfacial energy  $\gamma$  is another important, and maybe also the most uncertain, ingredient when doing precipitation kinetics simulations. Since the critical energy for nucleus formation  $G^*$  (equation 3.15) contains the cube of the interfacial energy and  $G^*$  appears in the exponent of the nucleation rate equation  $J$  (equation 3.11), small changes in  $\gamma$  lead to huge variations in  $J$ . Thus, precipitation kinetics simulations are very sensitive to this input quantity, which is often used in literature as 'fit - parameter' to adjust the simulations to experimental data.

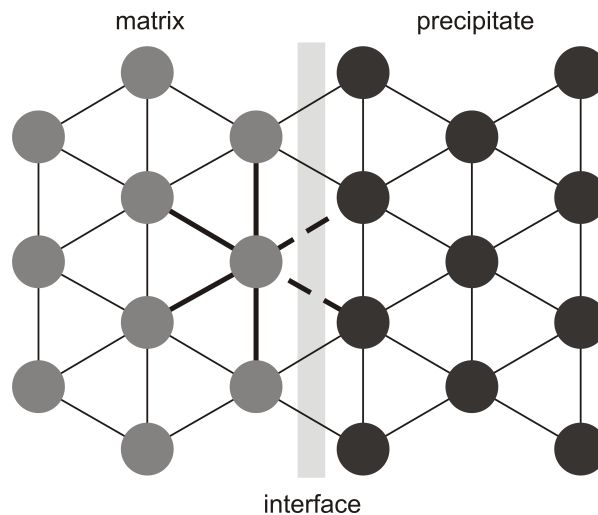


Figure 3.12: Two-dimensional interface with nearest-neighbor broken bonds [66].

Recently, Sonderegger and Kozeschnik [73] extended the classical 'Nearest-Neighbor-Broken-Bond' (NNBB) approach for the calculation of the interfacial energy  $\gamma$  to an  $n$ -next nearest-neighbor generalization of the NNBB concept, which is incorporated in the software package MatCalc [6–8]. Based on thermodynamic input data, this approach evaluates the interfacial energy, assuming a planar and sharp interface between precipitate and matrix, see figure 3.12. According to this image, the classical treatment of the 'NNBB' model considers only the first-nearest-neighbors, whereas the new extended approach takes also into account the second, third, ...  $n$ -nearest neighbors, which provides higher accuracy of the prediction.

The interfacial energy for a planar sharp interface  $\gamma_0$  is calculated as [73]

$$\gamma_0 = \frac{n_s \cdot z_{S,eff} \Delta E_{sol}}{N \cdot z_{L,eff}} \quad (3.26)$$

where  $\Delta E_{sol}$  is the enthalpy of solution,  $n_s$  the number of surface atoms per unit area within the surface plane and  $N$  the Avogadro number.  $z_{s,eff}$  represents the number of broken bonds across the interface per atom and  $z_{L,eff}$  accounts for the number of neighbors of an atom. The subscript 'eff' stands for effective structural parameters, accounting also the second, third, ...  $n$ -nearest neighbors.

However, corrections regarding the spherical shape of the nucleus and entropic contributions from finite interface thickness must be considered in the interfacial energy. Thus, the curvature dependency of the interfacial energy is treated according to [74]

$$\gamma' = \alpha \cdot \gamma_0 \quad (3.27)$$

and

$$\alpha = 1 - \frac{6}{11} \cdot \frac{r_1}{\rho} + 0.08921 \cdot \frac{r_1^2}{\rho^2} + 0.045 \cdot \ln \left( \frac{10}{3} \cdot \frac{\rho}{r_1} \right) \cdot \frac{r_1^2}{\rho^2}, \quad (3.28)$$

where  $\rho$  is the precipitate radius and  $r_1$  is the nearest-neighbor interatomic distance. The corresponding curve is shown in figure 3.13. For large precipitate radii, the correction function converges towards unity, which corresponds to a planar interface. With

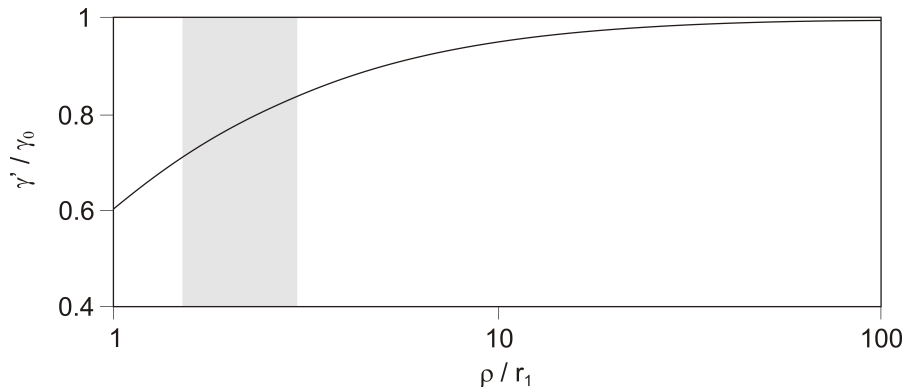


Figure 3.13: Size correction function (equation 3.28) [74].

decreasing radii, the size-correction function predicts a decrease in interfacial energy. The shaded area indicates critical nucleation radii with a size of 1.5 to 3 times  $r_1$ , which are typical cluster radii in the nucleation stage. The correction factor  $\alpha$  varies between 0.7 and 0.85 for these critical nucleation radii.

Particularly, close to the solution temperature, the effect of entropic contributions from finite interface thickness can become prominent. This behaviour of the interfacial energy has been observed in a previous work by Pudar et al. [81] and also treated theoretically in reference [82]. Hence, at temperatures close to the solution temperature of the precipitates, an additional correction factor  $\beta$  must be taken into account. Finally, the calculation of the interfacial energy can be summarized as,

$$\gamma = \gamma_0 \cdot \alpha \cdot \beta. \quad (3.29)$$

### 3.4.4 Numerical Treatment of the Evolution Equations

The fact that both, classical nucleation theory (CNT) (section 3.4.1) as well as the growth laws for the evolution of the precipitate radius (section 3.4.2) are differential equations, makes it necessary to integrate these expressions over time.

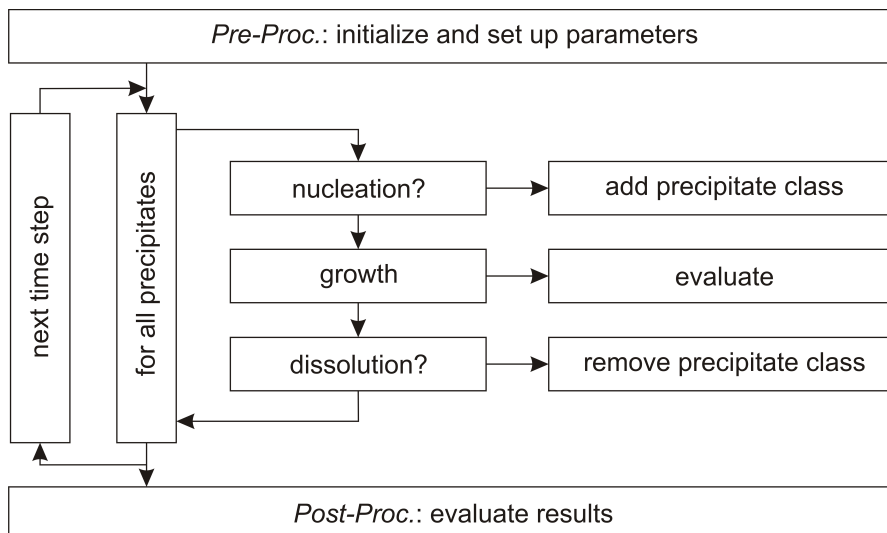


Figure 3.14: Flow chart for numerical integration of the evolution equations for precipitate nucleation and growth [7].

This is realized in MatCalc applying a numerical treatment of the evolution equations similar to the classical Kampmann - Wagner approach [79], according to the flow chart in figure 3.14. In the solution procedure, the time is discretised and the precipitate distribution is subdivided into different precipitate classes. All precipitate evolution equations are calculated for every precipitate class within each time step, determining all precipitate parameters. A detailed explanation of this algorithm can be found in reference [66].

# Chapter 4

## Results and Discussion

In this chapter, the models, previously introduced in section 3.4, are applied for the investigation of the precipitation of nitrides and carbides in microalloyed steel, using the thermo kinetic software MatCalc (version 5.30).

Starting with an analysis of the thermodynamic information in three different databases demonstrates, that there is room for improvement of the thermodynamic data. Therefore, the solubility products of several microalloy nitrides and carbides are studied in detail with special emphasis on microalloyed steel. Appropriate phase descriptions are determined and implemented in the thermodynamic database 'mc\_steel' (version 1.86) [83]. Additional essential ingredients for modeling and simulation of precipitation kinetics are analysed, among them the precipitate/matrix volumetric misfit, the temperature dependent Young's modulus, as well as the ratio between bulk and grain boundary diffusion, and temperature dependent formulations are presented. Performing detailed parameter studies, predominant precipitation at grain boundaries as well as the precipitation of randomly distributed particles are analysed, with special emphasis on the sensitivity of the calculations on these input quantities.

With the knowledge of the model characteristics of both, the novel grain boundary model (GBDG) as well as the classical treatment of randomly distributed precipitates with spherical diffusion fields (RSDG), together with the established input parameters, the precipitation kinetics of AlN and VN are calculated, considering simultaneous precipitation at grain boundaries and at dislocations. The calculations are carried out for

15 different chemical compositions at various different annealing temperatures, in ferrite as well as austenite, and compared to independent experimental data from literature. Finally, the precipitation processes of TiN, TiC and NbC during casting are analysed, taking into account the microsegregation of alloying elements during solidification of microalloyed steel. A new approach is presented, which allows calculating the precipitation parameters locally at every point of the dendrite.

## 4.1 Thermodynamic Information in Three Different Databases

When doing precipitation kinetics simulation, the main input quantity is the thermodynamic information. These thermodynamic data are generated from experimental investigations and summarised in thermodynamic databases. Hence, the thermodynamic assessments of different sources can show some discrepancies. The following section deals with a preliminary investigation of three different databases. In addition to the MatCalc database 'mc\_steel' (version 1.18) [83], the thermodynamic information in the commercial databases 'TCFE3' [84] and 'Fe-data6' [85] are analysed and compared to each other. Therefore, equilibrium calculations of microalloying phases in three alloys are carried out [86]. The chemical compositions of the alloys are given in table 4.1.

Table 4.1: Chemical composition of three typical microalloyed steels.

steel no.	C	Si	Mn	Al	Ti	V	N	Nb
1	0.22	0.243	1.2	0.025	0.002	0.001	0.0061	0.36
2	0.16	0.300	1.5	0.033	0.015	0.065	0.0097	0.03
3	0.40	0.590	1.4	0.030	0.021	0.150	0.0170	-

The simulations are performed using the description of the 'complex solution phase' (MX phase) of the three different databases. This complex solution phase is modeled according to the CALPHAD approach [70] as a part of the fcc matrix phase. For the simulation of microalloying phases with fcc structure, composition sets are defined

with microalloying elements as major constituents. However, minor amounts of other interstitial as well as substitutional elements are considered in the phase description, this is why we are talking of 'complex solution phases'.

Figures 4.1 to 4.3 show the calculated equilibrium phase fractions of the microalloy phases versus temperature, using the three different databases [86].

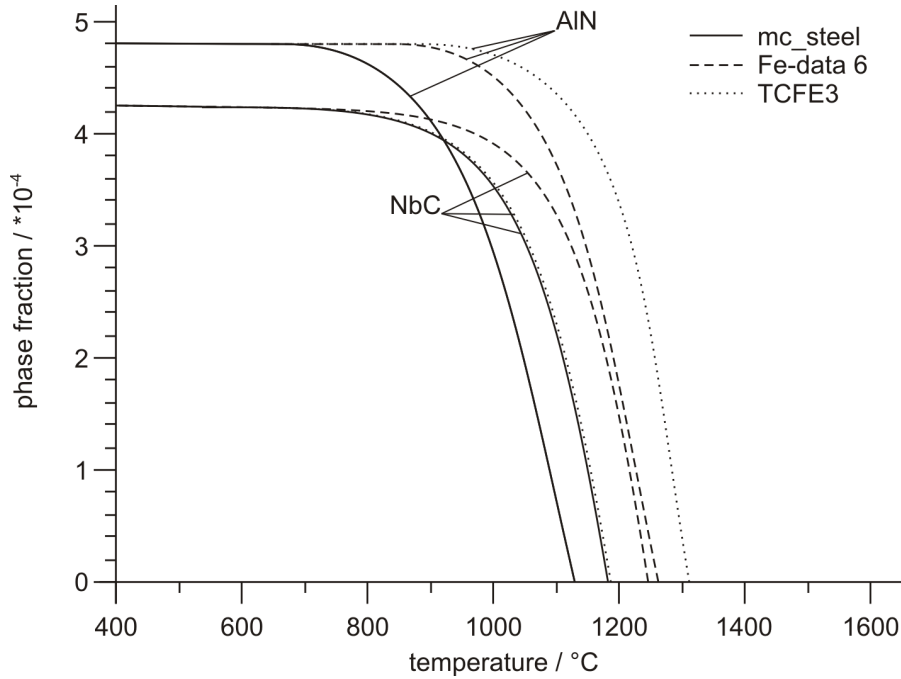


Figure 4.1: Comparison of the calculated phase fractions versus temperature of the phases NbC and AlN in steel 1, using three different databases.

Depending on the investigated alloy, precipitation of VC, NbC, Ti(C,N) or AlN is observed. The solution temperature for Ti(C,N) is calculated very similar, using the different databases. However, essential discrepancies in the calculated phase fraction of Ti(C,N) are observed with decreasing temperature. The database 'Fe-data6' provides generally smaller values compared to the other two databases. Additionally, it should be noted, that the chemical composition of Ti(C,N) differs also in a wide range, depending on the database used for the calculation.

Apart from Ti(C,N), most similar values concerning the solution temperature are calculated for NbC. The calculations with the databases 'mc.steel' and 'Fe-data6' deliver quite similar results and the database 'TCFE3' calculates a value, which is only 70°C



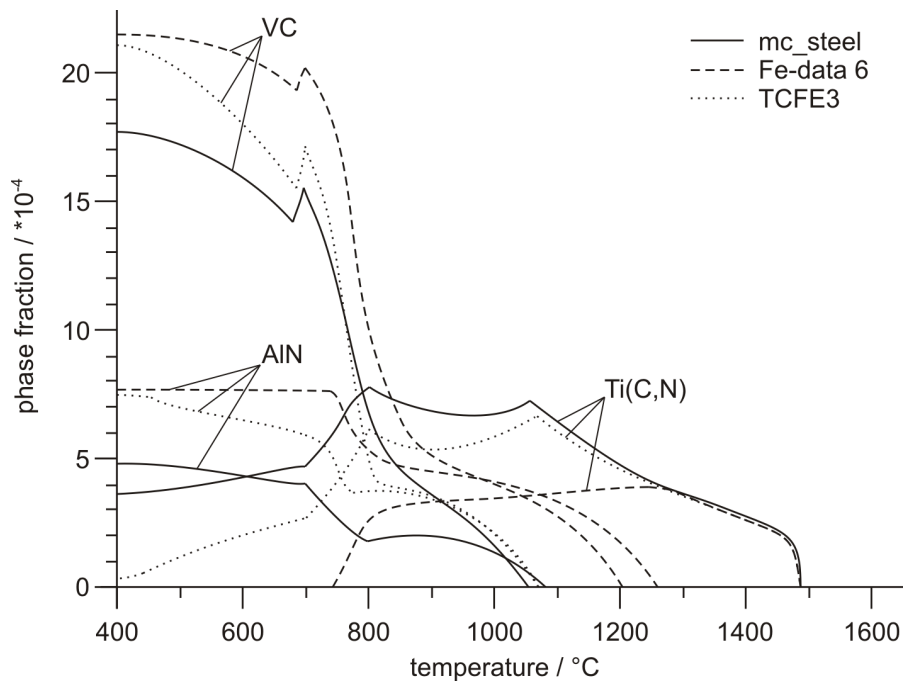


Figure 4.2: Comparison of the calculated phase fractions versus temperature of the phases VC, Ti(C,N) and AlN in steel 2, using three different databases.

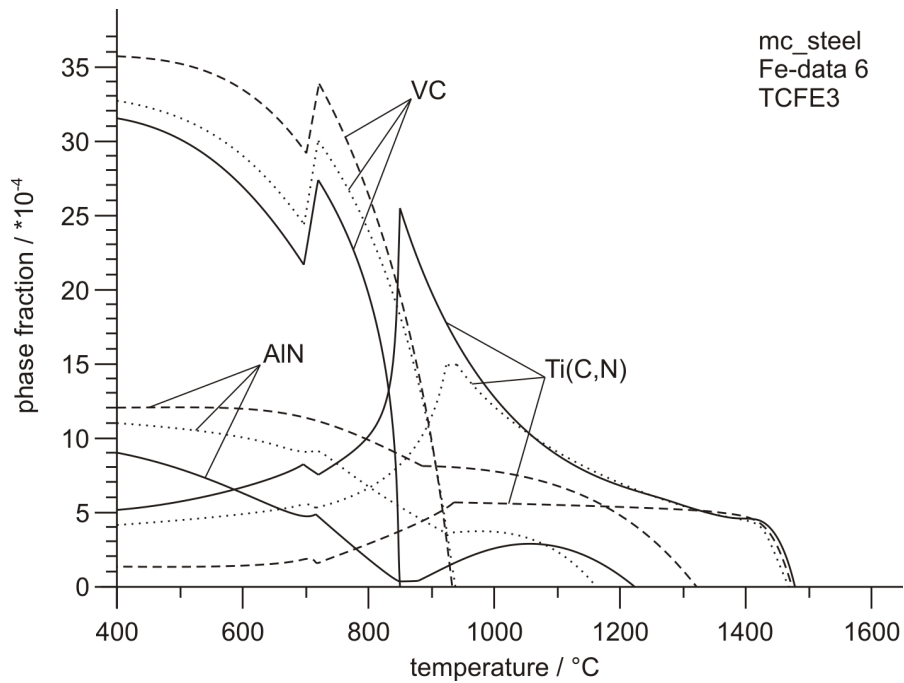


Figure 4.3: Comparison of the calculated phase fractions versus temperature of the phases VC, Ti(C,N) and AlN in steel 3, using three different databases.

higher, see figure 4.1.

The discrepancies of the predicted solution temperature are much more pronounced in the case of VC. Considering steel 2, the databases 'mc\_steel' and 'TCFE3' predict a solution temperature of 1050 °C, while the database 'Fe-data6' calculates 1200 °C.

Most dramatic differences are observed by the calculation of the equilibrium phase fraction of AlN. In the case of steel 1, the database 'TCFE3' calculates a solution temperature of 1315 °C, 'Fe-data6' predicts 1260°C and the calculation with the database 'mc\_steel' delivers 1125 °C. Thus, the difference in the predicted solution temperatures of AlN, using the databases 'TCFE3' and 'mc\_steel', is about 190 °C.

Substantial discrepancies are also observed regarding the predicted equilibrium phase fraction of the microalloy phases when using the different databases, especially for the calculation at low temperatures. These preliminary investigations clearly demonstrate the sensitivity of the calculations on the thermodynamic input data and thus the importance of the accuracy of these data. Therefore, the solubility of different stoichiometric phases is investigated in detail subsequently.

## 4.2 Solubility Products in Austenite

For microalloyed steels, it is common to express the solubility of nitrides and carbides using so-called solubility products, see section 3.2.3. With these data, thermodynamic parameters can be adjusted by comparing equilibrium calculations of solution temperatures of a nitride or carbide with the values obtained from the solubility product. The present section deals with the investigation of the solubility products of stoichiometric microalloy nitrides and carbides in the austenite phase field. Furthermore, the solution temperatures are calculated based on the thermodynamic databases 'TCFE3' [84], 'Fe-data6' [85] and 'mc\_steel' (version 1.86) [83] and compared to the solubility products given in literature. For the calculation of the solution temperatures, the ratio between the amount of the interstitial element (e.g. N) and the substitutional element (e.g. Al) is assumed to be 1/10.

### 4.2.1 Aluminium Nitride

A large number of papers are published on the solubility of AlN in austenite. Authors carried out experimental investigations [24, 29, 87–92] as well as thermodynamic calculations [93, 94] and determined various expressions for the description of the solubility of AlN in austenite. Table 4.2 and figure 4.4 summarize these data. In addition to the classical description of the solubility product (no. 1-14), Mayrhofer [89] introduced a formulation accounting for high Al contents (solubility product 15).

Unfortunately, the data obtained by the different authors show a very large scatter. This large scatter in the available solubility products for AlN precipitation gave motivation for detailed investigations already two decades ago. Wilson and Gladman [26] studied the solubility products published by different authors. Based on their assessment, they identified solubility product 5 [87] listed in table 4.2 as a 'most likely' one. They found support by the work of Irvine et al. [36], who report that their results are in close agreement with the results of Erasmus [24] and König et al. [88]. Particularly, the work of König et al. [88] seems to be much substantiated, because they investigated the solubility product of AlN using several different chemical compositions. In the present

assessment, we base our analysis on the experiments reported in references [41, 88, 89], where AlN precipitates have been clearly detected after different solution treatments, indicated by dots in figure 4.4, and have not been identified in treatments at higher temperature. The chemical compositions and different heat treatment temperatures are considered as the lower limit for the solubility product of AlN in austenite.

Table 4.2: Available solubility products of AlN.

no.	$\log k_s$	ref.
1	$4.382 - 11085/T$	[94]
2	$0.528 - 5938/T$	[90]
3	$1,79 - 7184/T$	[91]
4	$4,5989 - 11568/T$	[95]
5	$1,03 - 6770/T$	[87]
6	$1,95 - 7400/T$	[92]
7	$0,725 - 6180/T$	[24]
8	$0,309 - 6015/T$	[24]
9	$1,48 - 7500/T$	[35]
10	$1,8 - 7750/T$	[88]
11	$3,577 - 10020/T$	[29]
12	$2,923 - 9200/T$	[29]
13	$3,079 - 9295/T$	[29]
14	$6,4 - 14356/T$	[93]
15	$0,18 - 5675/T + 2,4 \cdot Al[wt\%]$	[89]

Together with the data of Saunders [96] and Zhang [97] for Al in the Fe matrix phases, following expression for the Gibbs energy of the AlN phase has been found most suitable for the simulations in the present work and is implemented in the database 'mc\_steel'.

$$G(FCC\_A1, AL : N; 0) = 298.15 + GHSERAL\# + GHSERNN\# + 80 \cdot T; \\ 6000.00 N$$

$$G(FCC\_A1, AL, FE : VA; 0) = 298.15 - 76066.1 + 18.6758 \cdot T; 6000.00 N$$

$$\begin{aligned}
 G(\text{FCC\_A1}, \text{AL}, \text{FE} : \text{VA}; 1) & 298.15 + 21167.4 + 1.3398 \cdot T; 6000.00 \text{ N} \\
 G(\text{BCC\_A2}, \text{AL} : \text{N}; 0) & 298.15 + 23000 + 10 \cdot T + \text{GHSERAL}\# \\
 & + 3 \cdot \text{GHSERNN}\#; 6000.00 \text{ N} \\
 G(\text{BCC\_A2}, \text{AL}, \text{FE} : \text{VA}; 0) & 298.15 - 122960 + 31.9888T; 6000.0 \text{ N} \\
 G(\text{BCC\_A2}, \text{AL}, \text{FE} : \text{VA}; 1) & 298.15 + 3089.2; 6000.0 \text{ N} \\
 G(\text{ALN}, \text{AL} : \text{N}; 0) & 298.15 - 262982 + 63 \cdot T + \text{GHSERAL}\# \\
 & + \text{GHSERNN}\#; 6000.00 \text{ N}
 \end{aligned}$$

The corresponding curve is shown as solid bold line in figure 4.4.

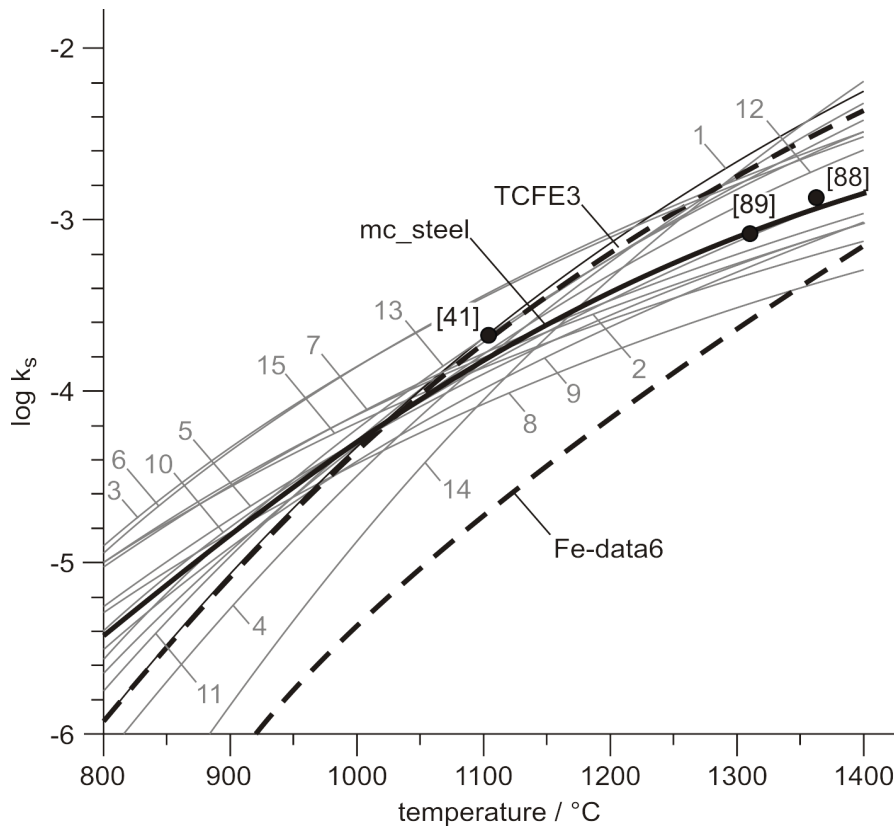


Figure 4.4: Comparison of solubility products and thermodynamic data of AlN.

### 4.2.2 Vanadium Nitride

The solubility of VN in the austenite phase field has been investigated using different experimental techniques [36,98,99] as well as thermodynamic calculations [49,100]. The corresponding relations are summarized in table 4.3 and figure 4.5. Nevertheless, it has been reported that, in V-microalloyed steels, the effect of Mn on the activity of V is especially important [36,49]. Mn is known to increase the activity coefficient of V and, at the same time, decrease the activity coefficient of C [49]. Therefore, the solubility product of VN in austenite increases slightly with increasing content of Mn [1]. Thus, Irvine et al. [36] suggested to take into account the Mn content, according to solubility product 1 in table 4.3, especially for microalloyed steels with significant Mn contents. All three databases predict the same solution behavior for VN in austenite. However, especially in the low temperature range of the austenite, they show some discrepancies compared to the experimental data [36,98,99] being in reasonable agreement.

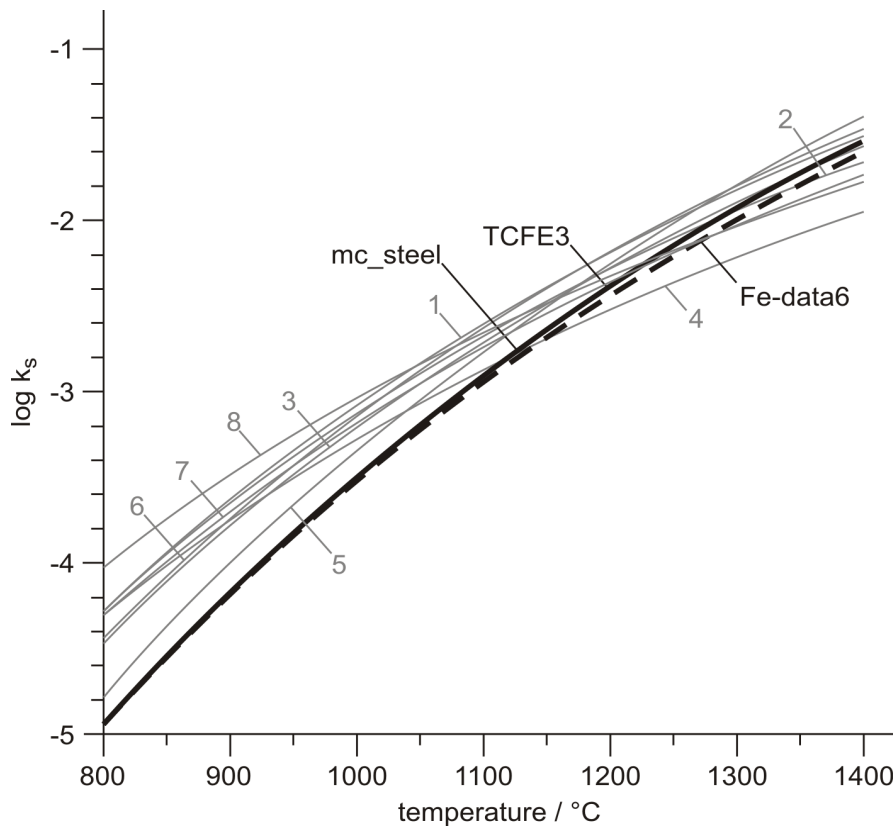


Figure 4.5: Comparison of solubility products and thermodynamic data of VN.

Table 4.3: Available solubility products of VN.

no.	$\log k_s$	ref.
1	$3.4 - 8330/T + 0.12 \cdot Mn[wt\%]$	[36]
2	$3.02 - 7840/T$	[98]
3	$3.63 - 8700/T$	[99]
4	$2.27 - 7070/T$	[100]
5	$-10.34 - 7600/T + 1.8 \cdot \ln T + 7.2 \cdot 10^{-5}T$	[49]
6	$3.84 - 8890/T$	[49]
7	$2.86 - 7700/T$	[27]
8	$2.21 - 6680/T$	[101]

The thermodynamic database 'mc.steel' incorporates the assessments of Ohtani and Hillert [102,103], which reads

$$\begin{aligned}
&G(FCC\_A1, V : N; 0) \quad 298.15 - 215000 + 101 \cdot T - 2.22 \cdot T \cdot LN(T) \\
&\quad + 0.00073 \cdot T^2 + GHSE RVV\# + GHSE RNN\#; \quad 6000.00 \quad N \\
&G(FCC\_A1, FE, V : VA; 0) \quad 298.15 - 15291 - 4.138T; \quad 6000.00 \quad N \\
&G(FCC\_A1, FE, V : N; 0) \quad 298.15 - 60000; \quad 6000.00 \quad N \\
&G(FCC\_A1, FE, V : N, VA; 0) \quad 298.15 - 120000; \quad 6000.00 \quad N \\
&G(FCC\_A1, V : N, VA; 0) \quad 298.15 - 131000 + 47.1 \cdot T; \quad 6000.00 \quad N \\
&G(BCC\_A2, V : N; 0) \quad 298.15 + 843000 + GHSE RVV\# + 3 \cdot GHSE RNN\#; \\
&G(BCC\_A2, V : N, VA; 0) \quad 298.15 - 1680000 + 325 \cdot T; \quad 6000.00 \quad N \\
&G(BCC\_A2, FE, V : VA; 0) \quad 298.15 - 23674 + 0.465 \cdot T; \quad 6000.00 \quad N \\
&G(BCC\_A2, FE, V : VA; 1) \quad 298.15 + 8283; \quad 6000.00 \quad N \\
&G(VN, V : N; 0) \quad 298.15 - 215000 + 101 \cdot T - 2.22 \cdot T \cdot LN(T) \\
&\quad + 0.00073 \cdot T^2 + GHSE RVV\# + GHSE RNN\#; \quad 6000.00 \quad N \\
&G(VN, V : VA; 0) \quad 298.15 + 7500 + 1.7 \cdot T + GHSE RVV\#; \quad 4000.00 \quad N \\
&G(VN, V : N, VA; 0) \quad 298.15 - 131000 + 47.1 \cdot T; \quad 6000.00 \quad N.
\end{aligned}$$

Only solubility product 5, introduced by Lagneborg et al. [49], is in acceptable agreement with the three thermodynamic databases. However, this solubility product is not representative, because it is also deduced from the assessment of Ohtani and Hillert [102].

### 4.2.3 Vanadium Carbide

The solubility of VC in the austenite phase field is significantly higher than those of the other microalloy carbides and nitrides. However, the description of VC is a little more complicated due to stoichiometric effects. There is some evidence that the chemical formula for vanadium carbide is close to  $V_4C_3$  [1,104]. Therefore, the solubility product may be defined as

$$k_s = \frac{[V]^{4/3}[C]}{a_{(V_4C_3)}^{1/3}}, \quad (4.1)$$

rather than equation 3.7. However, not all authors accepted this definition. For instance Narita [99] based his solubility product on the classical formulation according to 3.7, whereas Bungardt et al. [105], Wriedt et al. [104] and Aronsson [106] use equation 4.1 assuming the formation of  $V_4C_3$ . The different expressions for the solubility product of VC are listed in table 4.4 and pictured in figure 4.6.

Table 4.4: Available solubility products of VC.

no.	$\log k_s$	ref.
1	$6.72 - 9500/T$	[99]
2	$4.45 - 6560/T$	[104]
3	$7.06 - 10800/T$	[105]
4	$5.35 - 8000/T$	[106]
5	$3.685 - 5531/T$	[60]
6	$1.32 - 2980/T + (1160/T - 0.2) \cdot C[wt\%]$	[107]

Especially, the work of Wriedt et al. [104] (solubility product 2) seems to be substantiated, since they investigated several alloys with different V/C ratios. Acevedo et al. [60]



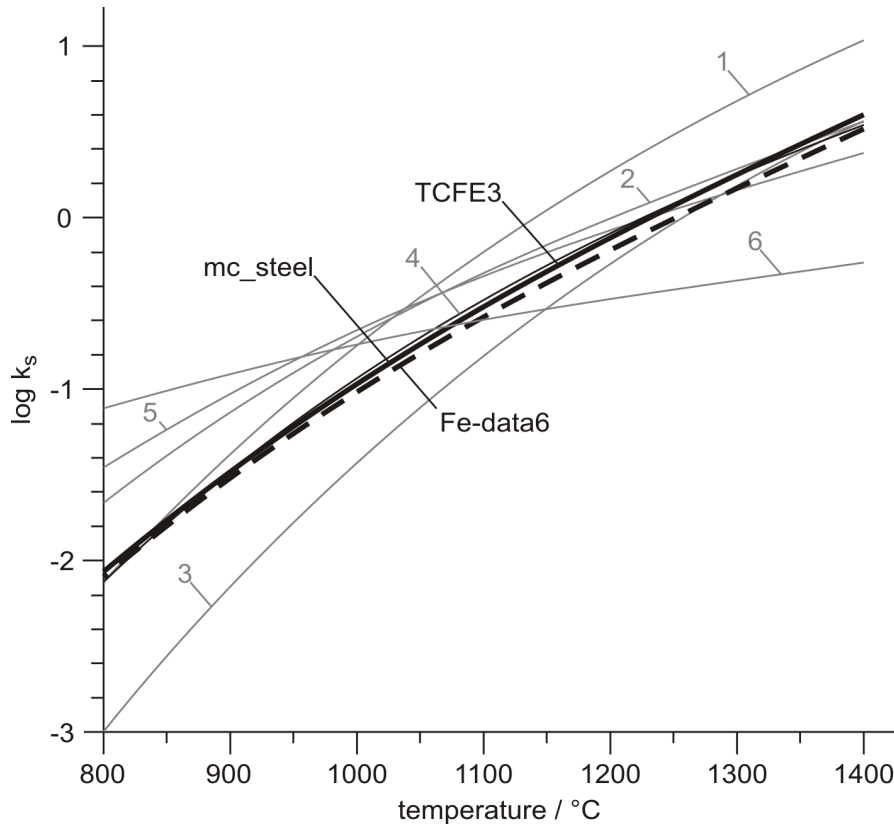


Figure 4.6: Comparison of solubility products and thermodynamic data of VC.

identified the formation of  $V_6C_5$  carbides by doing TEM investigations and, therefore, they introduced solubility product 5. Ohtani et al. [107] related their expression of the solubility of VC to the C content, see solubility product 6.

Similar to the observations for VN, all three databases predict the same solution behavior for VC in austenite, which is in close agreement to solubility product 4 [106].

The database 'mc\_steel' contains the thermodynamic assessment of Huang [108], which describes the solubility of VC with

$$G(FCC\_A1, V : C; 0) \quad 298.15 \quad - \quad 117302 \quad + \quad 262.57 \cdot T \quad - \quad 41.756 \cdot T \cdot LN(T)$$

$$\quad - \quad 0.00557101 \cdot T^2 \quad + \quad 590546 \cdot T^{-1}; \quad 6000.00 \quad N$$

$$G(FCC\_A1, V : C, VA; 0) \quad 298.15 \quad - \quad 74811 \quad + \quad 10.201 \cdot T; \quad 6000.00 \quad N$$

$$G(FCC\_A1, V : C, VA; 1) \quad 298.15 \quad - \quad 30394; \quad 6000.00 \quad N$$

$$G(FCC\_A1, FE, V : VA; 0) \quad 298.15 \quad - \quad 15291 \quad - \quad 4.138 \cdot T; \quad 6000.00 \quad N$$

$$\begin{aligned}
G(FCC\_A1, FE, V : C; 0) & 298.15 - 7645.5 - 2.069 \cdot T; 6000.00 N \\
G(FCC\_A1, FE, V : C; 1) & 298.15 - 7645.5 - 2.069 \cdot T; 6000.00 N \\
G(FCC\_A1, FE, V : C, VA; 0) & 298.15 - 40000; 6000.00 N \\
G(BCC\_A2, V : C; 0) & 298.15 + 108449 + GHSERVV\# + 3 \cdot GHSE RCC\#; \\
& 6000.00 N \\
G(BCC\_A2, V : C, VA; 0) & 298.15 - 297868; 6000.00 N \\
G(BCC\_A2, FE, V : VA; 0) & 298.15 - 23674 + 0.465 \cdot T; 6000.00 N \\
G(BCC\_A2, FE, V : VA; 1) & 298.15 + 8283; 6000.00 N \\
G(BCC\_A2, FE, V : C; 0) & 298.15 - 23674 + 0.465 \cdot T; 6000.00 N \\
G(BCC\_A2, FE, V : C; 1) & 298.15 + 8283; 6000.00 N \\
G(VC, V : C; 0) & 298.15 - 117302 + 262.57 \cdot T - 41.756 \cdot T \cdot LN(T) \\
& - 0.00557101 \cdot T^2 + 590546 \cdot T^{-1}; 6000.00 N \\
G(VC, V : VA; 0) & 298.15 + 7500 + 1.7 \cdot T + GHSERVV\#; 4000.00 N \\
G(VC, V : C, VA; 0) & 298.15 - 74811 + 10.201 \cdot T; 6000.00 N \\
G(VC, V : C, VA; 1) & 298.15 - 30394; 6000.00 N.
\end{aligned}$$

#### 4.2.4 Niobium Nitride

Several authors have reported about the solubility product of NbN in austenite. The equations are summarized in table 4.5 and figure 4.7. Whereas the solubility products 1-4 are determined experimentally [99, 109–111] (e.g. chemical separation of particles), others are determined by thermodynamic calculations, e.g. solubility product 7 [112]. However, solubility product 7 is an approximate expression for the chemical compound  $NbN_{7/8}$ . All other authors formulated the solubility of NbN based on the stoichiometric compound NbN. Gladman [1] recommends to use solubility product 1, determined by Narita [99], which is in very close agreement with the data of Smith [109], Mori et al. [111] and Klinkenberg et al. [57].

The thermodynamic database 'TCFE3' reproduces these recommended values best. The databases 'Fe\_data6' and 'mc\_steel' calculate approximately the same results, which show

Table 4.5: Available solubility products of NbN.

no.	$\log k_s$	ref.
1	$2.8 - 8500/T$	[99]
2	$4.04 - 10230/T$	[109]
3	$3.7 - 10800/T$	[110]
4	$3.79 - 10150/T$	[111]
5	$1.71 - 7000/T$	[49]
6	$6.63 - 12500/T$	[49]
7	$2.86 - 7927/T$	[112]
8	$4.2 - 10000/T$	[113]
9	$3.82 - 9940/T$	[57]
10	$5.09 - 11880/T$	[101]

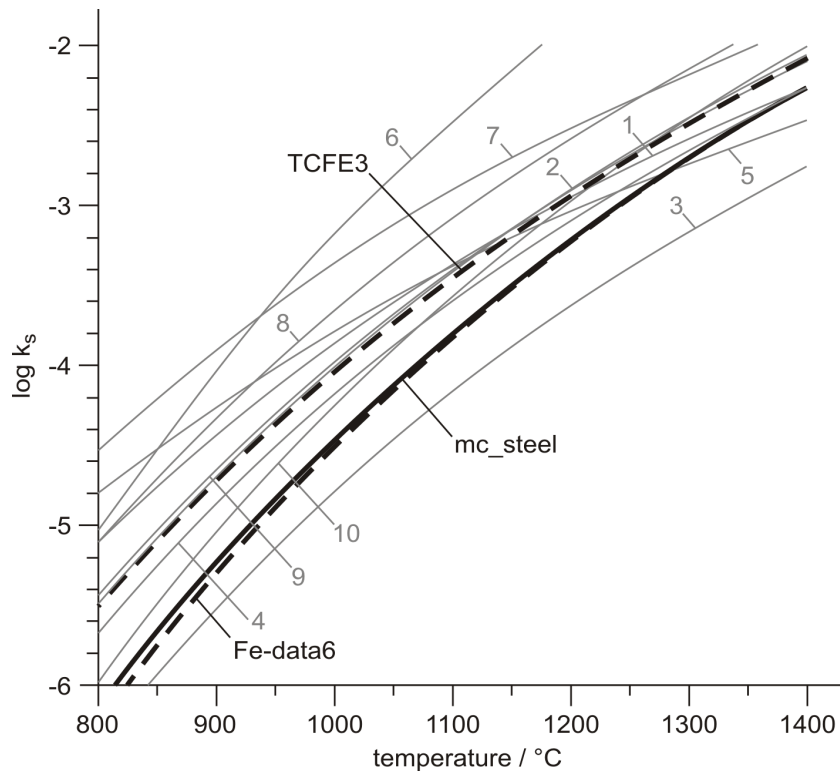


Figure 4.7: Comparison of solubility products and thermodynamic data of NbN.

some discrepancies to the experimentally obtained and recommended solubility products. However, the basis for the calculation of NbN in the database 'mc\_steel' is the thermodynamic assessment after Lee [114], which is given with

$$\begin{aligned}
 &G(FCC\_A1, NB : N; 0) \text{ 298.15} - 227779 + 120.567 \cdot T - 4 \cdot T \cdot LN(T) \\
 &\quad + GHSEARNB\# + GHSEARNN\#; \text{ 6000.00 } N \\
 &G(FCC\_A1, NB : N, VA; 0) \text{ 298.15} - 65218; \text{ 6000.00 } N \\
 &G(FCC\_A1, FE, NB : VA; 0) \text{ 298.15} - 2000; \text{ 6000.00 } N \\
 &G(BCC\_A2, NB : N; 0) \text{ 298.15} - 207283 + 110.085 \cdot T + GHSEARNB\# \\
 &\quad + 3 \cdot GHSEARNN\#; \text{ 6000.00 } N \\
 &G(BCC\_A2, NB : N, VA; 0) \text{ 298.15} - 357277 + 129.474 \cdot T; \text{ 6000.00 } N \\
 &G(BCC\_A2, FE, NB : VA; 0) \text{ 298.15} - 4400 + 6.333 \cdot T; \text{ 6000.00 } N \\
 &G(NBN, NB : N; 0) \text{ 298.15} - 227779 + 120.567 \cdot T - 4 \cdot T \cdot LN(T) \\
 &\quad + GHSEARNB\# + GHSEARNN\#; \text{ 6000.00 } N \\
 &G(NBN, NB : VA; 0) \text{ 298.15} + 13500 + 1.7 \cdot T + GHSEARNB\#; \text{ 5000.00 } N \\
 &G(NBN, NB : N, VA; 0) \text{ 298.15} - 65218; \text{ 6000.00 } N.
 \end{aligned}$$

#### 4.2.5 Niobium Carbide

Most papers have been published on the solubility of NbC in austenite. The determination methods cover many experimental techniques as well as thermodynamic calculations or statistical treatments of existing solubility products. Lots of them are summarized in the work of Palmiere et al. [115]. The different solubility products for NbC in austenite are given in table 4.6 and shown in figure 4.8.

Different authors report about different chemical formulas, similar to the case of VC. Whereas some authors assume the stoichiometric compound NbC [99,107,115–119,122–124], others [111,112,120,121] base the formulation of their solubility product on NbC<sub>7/8</sub>. Palmiere et al. [115] performed detailed atom probe investigations and demonstrated that precipitates ranging in size from 0.5 to 10 nm have compositions ranging from NbC<sub>8/10</sub> to stoichiometric NbC. However, they suggested to use a stoichiometric de-

Table 4.6: Available solubility products of NbC.

no.	$\log k_s$	ref.
1	$2.06 - 6700/T$	[115]
2	$-0.63 - 2500/T$	[116]
3	$3.04 - 7290/T$	[117]
4	$3.7 - 9100/T$	[118]
5	$3.42 - 7900/T$	[99]
6	$4.37 - 9290/T$	[119]
7	$3.18 - 7700/T$	[111]
8	$3.11 - 7520/T$	[120]
9	$2.96 - 7510/T$	[120]
10	$3.4 - 7200/T$	[120]
11	$2.81 - 7019.5/T$	[112]
12	$3.4 - 7920/T$	[121]
13	$2.26 - 6670/T$	[36]
14	$3.89 - 8030/T + (1150/T - 0.05) \cdot C[wt\%]$	[122]
15	$1.18 - 4880/T + (1590/T - 0.1) \cdot C[wt\%]$	[107]
16	$1.74 - 5600/T + (1380/T - 0.027) \cdot C[wt\%]$	[123]
17	$3.31 - 7970/T + (1371/T - 0.9) \cdot Mn[wt\%] -$ $(75/T - 0.0504) \cdot Mn^2[wt\%]$	[124]
18	$3.31 - 7970/T - (735/T - 0.348) \cdot Si[wt\%]$	[124]
19	$3.31 - 7970/T + (1113/T - 0.691) \cdot Cr[wt\%] -$ $(38/T - 0.0228) \cdot Cr^2[wt\%]$	[124]
20	$3.31 - 7970/T + (148/T - 0.0904) \cdot Ni[wt\%] +$ $(8.5/T - 0.0068) \cdot Ni^2[wt\%]$	[124]

scription according to solubility product 1. Compared to other solubility products in figure 4.8, this description predicts markedly lower solubility levels of Nb in austenite. Therefore, the authors argued that the measured data using the atom probe yielded a

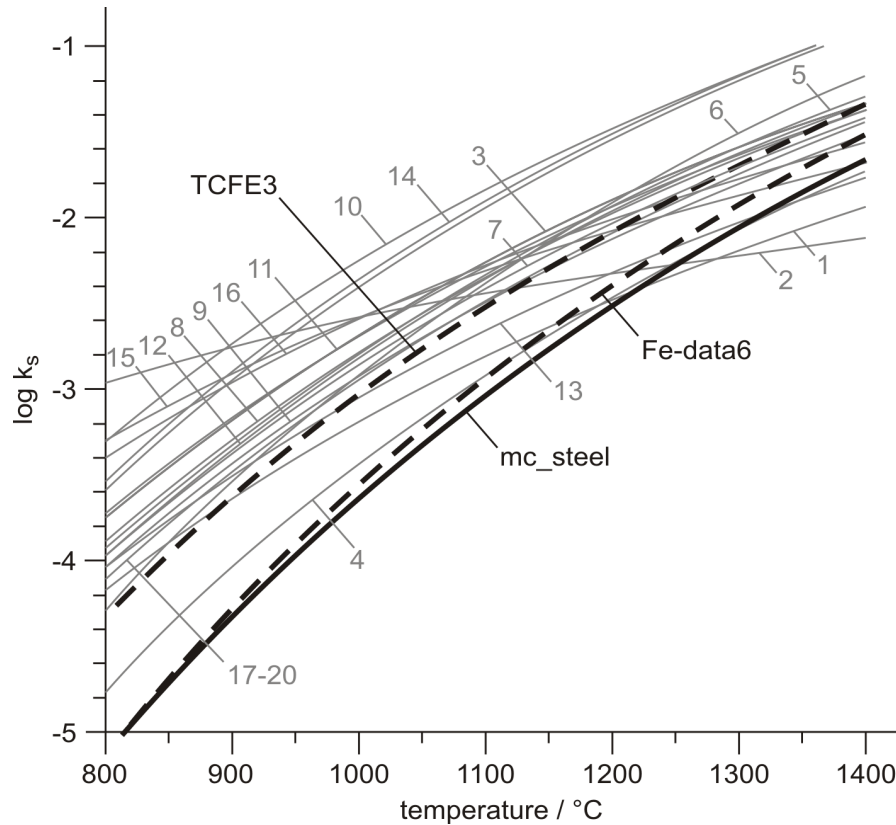


Figure 4.8: Comparison of solubility products and thermodynamic data of NbC.

more reliable description of the solution behavior compared to other techniques employed for previous investigations. Nordberg and Aronsson [120] did statistical investigations of existing solubility products using the method of least squares. Table 4.6 shows three different solubility products provided by these authors. Solubility product 8 is based on the assumption of the formation of  $\text{NbC}_{7/8}$ , whereas solubility product 9 underlies the stoichiometric compound of NbC. In addition to the statistical analysis, they performed thermodynamic calculations resulting in solubility product 10, which is also based on the formation of  $\text{NbC}_{7/8}$ . Some authors extended the classical description of the solubility product (equation 3.7) by an additional term, considering ternary interaction parameters. Therefore, Balasubramanian et al. [122] as well as Ohtani et al. [107, 123] added a term taking into account the C-concentration and the strong interactions between Nb and C at higher C concentrations, see solubility products 14-16. Moreover, Koyama et al. [124] studied the influence of Mn, Si, Cr and Ni on the solubility of NbC in the

austenite phase field, see solubility products 17-20. Mn and Cr increase the solubility of NbC remarkably. Also, Ni increases the solubility of NbC in austenite at comparatively low temperatures, while the effect is weak at higher temperatures. On the other hand, additions of Si decrease the solubility of NbC. Koyama et al. [124] attributed these effects to the effect of each element on the activity of Nb as well as that of C.

The three thermodynamic databases predict lower solubilities for NbC in austenite compared to most of the solubility products. This is in agreement with the atom probe investigation of Palmiere et al. [115]. Furthermore, the slope of the curves of the thermodynamic assessments differ significantly from those of most solubility products, which is related to deviations in entropic contributions. The thermodynamic database 'mc\_steel' contains the description of Povoden et al. [125] given with

$$\begin{aligned}
G(FCC\_A1, NB : C; 0) & 298.15 - 162000 + 283.912 \cdot T - 46.34274 \cdot T \cdot LN(T) \\
& - 0.0029287 \cdot T^2 - 1.02788144 \cdot 10^9 \cdot T^{-3} + 563374 \cdot T^{-1}; 6000.00 N \\
G(FCC\_A1, FE, NB : VA; 0) & 298.15 - 2000; 6000.00 N \\
G(FCC\_A1, FE, NB : C, VA; 0) & 298.15 - 40000; 6000.00 N \\
G(FCC\_A1, NB : C, VA; 0) & 298.15 - 83000 + 20 \cdot T; 6000.00 N \\
G(FCC\_A1, NB : C, VA; 2) & 298.15 - 61342 + 25.207 \cdot T; 6000.00 N \\
G(BCC\_A2, NB : C; 0) & 298.15 + 520879 - 70.807 \cdot T \\
& + GHSEARNB\# + 3 \cdot GHSEARCC\#; 6000.00 N \\
G(BCC\_A2, NB : C, VA; 0) & 298.15 - 588143; 6000.00 N \\
G(NBC, NB : C; 0) & 298.15 - 162000 + 283.912 \cdot T - 46.34274 \cdot T \cdot LN(T) \\
& - 0.0029287 \cdot T^2 - 1.02788144 \cdot 10^9 \cdot T^{-3} + 563374 \cdot T^{-1}; 6000.00 N \\
G(NBC, NB : VA; 0) & 298.15 + 13500 + 1.7 \cdot T + GHSEARNB\#; 5000.00 N \\
G(NBC, NB : C, VA; 0) & 298.15 - 83000 + 20 \cdot T; 6000.00 N \\
G(NBC, NB : C, VA; 2) & 298.15 - 61342 + 25.207 \cdot T; 6000.00 N,
\end{aligned}$$

which is deduced from the assessment of Huang et al. [126, 127].

### 4.2.6 Titanium Nitride

Studies of the solubility of TiN have shown that, from all microalloy carbides and nitrides, TiN is by far the most stable one [1]. Different expressions of solubility products have been obtained by experimental investigations [99, 128–131] as well as thermodynamic analysis [102, 131, 132]. The results are presented in table 4.7 and illustrated in figure 4.9.

Table 4.7: Available solubility products of TiN.

no.	$\log k_s$	ref.
1	$3.82 - 15020/T$	[99]
2	$4.35 - 14890/T$	[131]
3	$0.32 - 8000/T$	[128]
4	$8.48 - 12900/T + 1.457 \cdot \ln T + 0.00028 \cdot T$	[102]
5	$5.19 - 15490/T$	[129]
6	$5 - 14400/T$	[132]
7	$5.15 - 13925/T$	[129]
8	$4.94 - 14400/T$	[130]
9	$5.4 - 15791/T$	[27]

The data of Wada and Pehlke [130] are in very good agreement with those of Roberts [132] (solubility products 6 and 8). Furthermore, the values of Kunze [129] are very similar to those used by Turkdogan [27] (solubility products 5 and 9). The lowest solubility of TiN in austenite is predicted by Narita [99] (solubility product 1), whereas the highest values are proposed by using an estimation after Roberts [132], reported in the work of Kunze [129] (solubility product 7). Furthermore, figure 4.9 shows that the solubility data of the different authors underlie a wide spread. The scatter can perhaps be attributed to the low solubility of TiN in austenite, bearing in mind that TiN can precipitate in considerable amounts in liquid steel. However, apart from the data of Matsuda and Okumura [128], the slope of the curves are similar and parallel to each other, but the intercept varies significantly. All three databases predict the same



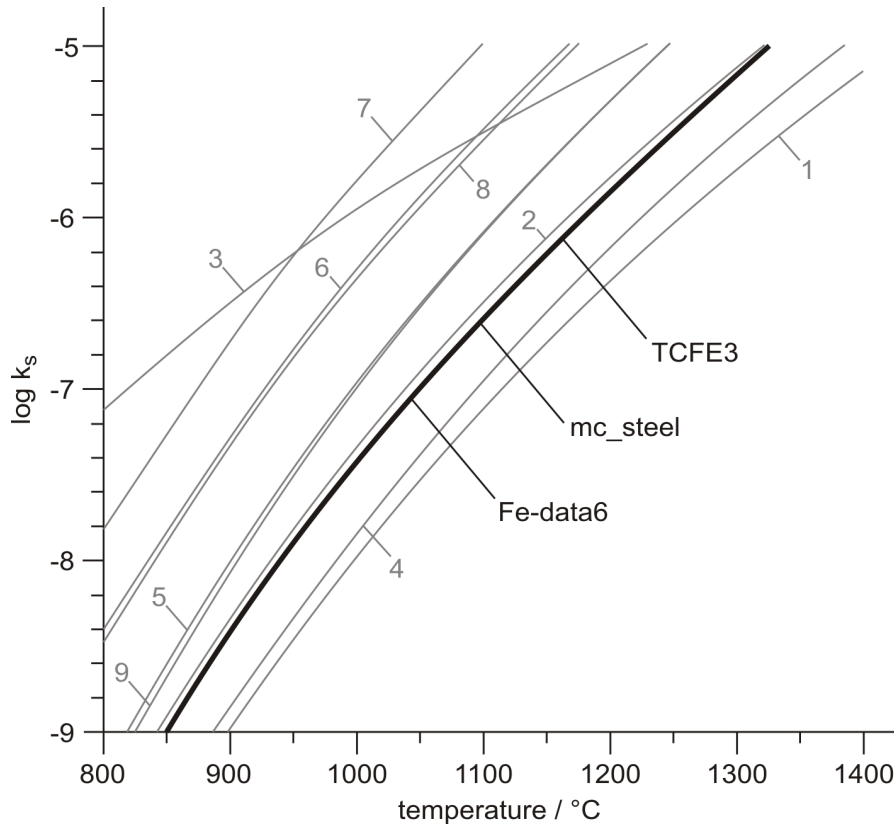


Figure 4.9: Comparison of solubility products and thermodynamic data of TiN.

solubility temperatures for TiN in austenite. The thermodynamic database 'mc\_steel' contains thermodynamic information of Lee [114] and Frisk [133], which is given with

$$\begin{aligned}
 G(\text{FCC\_A1}, \text{TI} : \text{N}; 0) & 298.15 \text{ GHSERTIN\#}; 6000.00 \text{ N} \\
 G(\text{FCC\_A1}, \text{TI} : \text{N}, \text{VA}; 0) & 298.15 - 42704; 6000.00 \text{ N} \\
 G(\text{FCC\_A1}, \text{TI} : \text{N}, \text{VA}; 1) & 298.15 - 13989; 6000.00 \text{ N} \\
 G(\text{FCC\_A1}, \text{FE}, \text{TI} : \text{VA}; 0) & 298.15 - 51136 + 10.4 \cdot T; 6000.00 \text{ N} \\
 G(\text{FCC\_A1}, \text{FE}, \text{TI} : \text{VA}; 1) & 298.15 - 1950 - 6 \cdot T; 6000.00 \text{ N} \\
 G(\text{FCC\_A1}, \text{FE}, \text{TI} : \text{VA}; 2) & 298.15 + 14875; 6000.00 \text{ N} \\
 G(\text{BCC\_A2}, \text{TI} : \text{N}; 0) & 298.15 + 2604202 + 118.04 \cdot T + \text{GHSERTIN\#} \\
 & + 2 \cdot \text{GHSENN\#}; 6000.00 \text{ N} \\
 G(\text{BCC\_A2}, \text{TI} : \text{N}, \text{VA}; 0) & 298.15 - 3215338; 6000.00 \text{ N} \\
 G(\text{BCC\_A2}, \text{FE}, \text{TI} : \text{VA}; 0) & 298.15 - 59098 + 11.5 \cdot T; 6000.00 \text{ N}
 \end{aligned}$$

$$G(BCC\_A2, FE, TI : VA; 1) 298.15 - 1796 + 1 \cdot T; 6000.00 N$$

$$G(BCC\_A2, FE, TI : VA; 2) 298.15 + 5602 + 3.5 \cdot T; 6000.00 N$$

$$G(TIN, TI : N; 0) 298.15 GHSERTIN\#; 6000.00 N$$

$$G(TIN, TI : VA; 0) 298.15 + 6000 - 0.1 \cdot T + GHSERTI\#; 3000.00 N$$

$$G(TIN, TI : N, VA; 0) 298.15 - 42704; 6000.00 N$$

$$G(TIN, TI : N, VA; 1) 298.15 - 13989; 6000.00 N.$$

The calculations with these input data are in very good agreement with the work of Inoue et al. [131] (solubility product 2), who did experimental investigations as well as thermodynamic calculations.

#### 4.2.7 Titanium Carbide

The solubility product equations for TiC in austenite are listed in table 4.8 and are shown graphically in figure 4.10.

Table 4.8: Available solubility products of TiC.

no.	$\log k_s$	ref.
1	$5.33 - 10475/T$	[99]
2	$2.75 - 7000/T$	[36]
3	$5.12 - 10300/T$	[1]
4	$4.03 - 8720/T$	[113]
5	$4.1 - 9070/T + (1205/T + 0.24) \cdot C[wt\%]$	[134]
6	$4.2 - 8970/T + (1600/T + 0.2) \cdot C[wt\%]$	[107]
7	$3.23 - 7430/T + (1300/T + 0.03) \cdot C[wt\%]$	[135]

Experimental investigations have been performed by Narita [99], Irvine et al. [36] and Balasubramanian et al. [113, 134] (solubility products 1-5), while thermodynamic calculations are done by Ohtani et al. [107, 135] (solubility products 6-7). Compared to the data for TiN (see figure 4.9), the results of the different authors are in good agreement. However, Gladman [1] suggested that some difficulties associated with TiC solubility

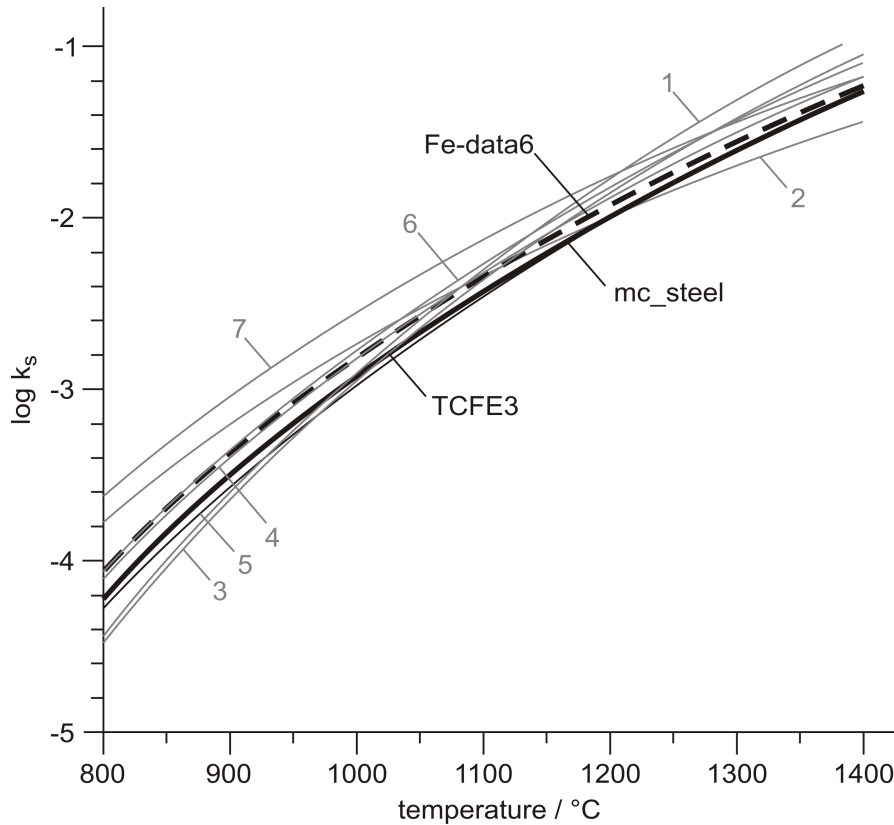


Figure 4.10: Comparison of solubility products and thermodynamic data of TiC.

data is related to the presence of N in the steels used. In section 4.2.6, it is already mentioned that TiN is extremely stable and needs to be considered when interpreting experimental results. Similar to the expressions for NbC (see table 4.6), Balasubramanian et al. [134] as well as Ohtani et al. [135] added a term to the solubility product accounting for the C-concentration by taking into account the strong interactions between Ti and C at higher C concentrations.

The databases 'mc\_steel' and 'TCFE3' predict the same solution temperatures and are very close to the data calculated with the database 'Fe-data6'. The database 'mc\_steel' contains thermodynamic information of Lee [114] and Frisk [133], which reads

$$G(FCC\_A1, TI : C; 0) \text{ 298.15} + GTIC\#; \text{ 6000.00 } N$$

$$G(FCC\_A1, TI : C, VA; 0) \text{ 298.15} - 40219.8 - 18.25 \cdot T; \text{ 6000.00 } N$$

$$G(FCC\_A1, TI : C, VA; 1) \text{ 298.15} - 108521 + 12.735 \cdot T; \text{ 6000.00 } N$$

$$G(FCC\_A1, FE, TI : VA; 0) \text{ 298.15} - 51136 + 10.4 \cdot T; \text{ 6000.00 } N$$

$$\begin{aligned}
G(FCC\_A1, FE, TI : VA; 1) & 298.15 - 1950 - 6 \cdot T; 6000.00 N \\
G(FCC\_A1, FE, TI : VA; 2) & 298.15 + 14875; 6000.00 N \\
G(FCC\_A1, FE, TI : C, VA; 0) & 298.15 - 180000; 6000.00 N \\
G(FCC\_A1, FE, TI : C, VA; 1) & 298.15 + 318559 + 39.74 \cdot T; 6000.00 N \\
G(FCC\_A1, FE, TI : C, VA; 2) & 298.15 - 318559 - 39.74 \cdot T; 6000.00 N \\
G(BCC\_A2, TI : C; 0) & 298.15 + 2 \cdot GHSERCC\# + GTIC\# + 600000; \\
& 4000.00 N \\
G(BCC\_A2, TI : C, VA; 0) & 298.15 - 957000 + 75 \cdot T; 6000.00 N \\
G(BCC\_A2, FE, TI : VA; 0) & 298.15 - 59098 + 11.5 \cdot T; 6000.00 N \\
G(BCC\_A2, FE, TI : VA; 1) & 298.15 - 1796 + 1 \cdot T; 6000.00 N \\
G(BCC\_A2, FE, TI : VA; 2) & 298.15 + 5602 + 3.5 \cdot T; 6000.00 N \\
G(TIC, TI : C; 0) & 298.15 - 168261 + 293.73187 \cdot T - 48.0195 \cdot T \cdot LN(T) \\
& - 0.00272 \cdot T^2 + 819000 \cdot T^{-1} - 2.03 \cdot 10^9 \cdot T^2 - 3; 6000.00 N \\
G(TIC, TI : VA; 0) & 298.15 + 6000 - 0.1 \cdot T + GHSERTI\#; 3000.00 N \\
G(TIC, TI : C, VA; 0) & 298.15 - 52702 - 4.6 \cdot T; 6000.00 N \\
G(TIC, TI : C, VA; 1) & 298.15 - 121367 + 31.5 \cdot T; 6000.00 N.
\end{aligned}$$

The results of all three databases are in good agreement with the solubility products listed in table 4.8.

#### 4.2.8 Solubility Data Calculated with the Database 'mc\_steel'

Figure 4.11 summarises the solubility of microalloy carbides and nitrides calculated with MatCalc (version 5.30) using the thermodynamic database 'mc\_steel (version 1.86)'. The results are similar to the data proposed by Lagneborg et al. [49], see figure 3.8 in section 3.2.3.2. Accordingly, all microalloying nitrides are more stable than their corresponding carbides. The solubility products of the different phases in austenite range from 1 to -9. Furthermore, the order of the solubility of the nitrides and carbides corresponds to the data proposed by Lagneborg et al. [49]. Thus, TiN is an extremely stable microalloying

phase, whereas VC is characterized by high solubility. The solubility of the other phases is somewhere in between these two extrema.

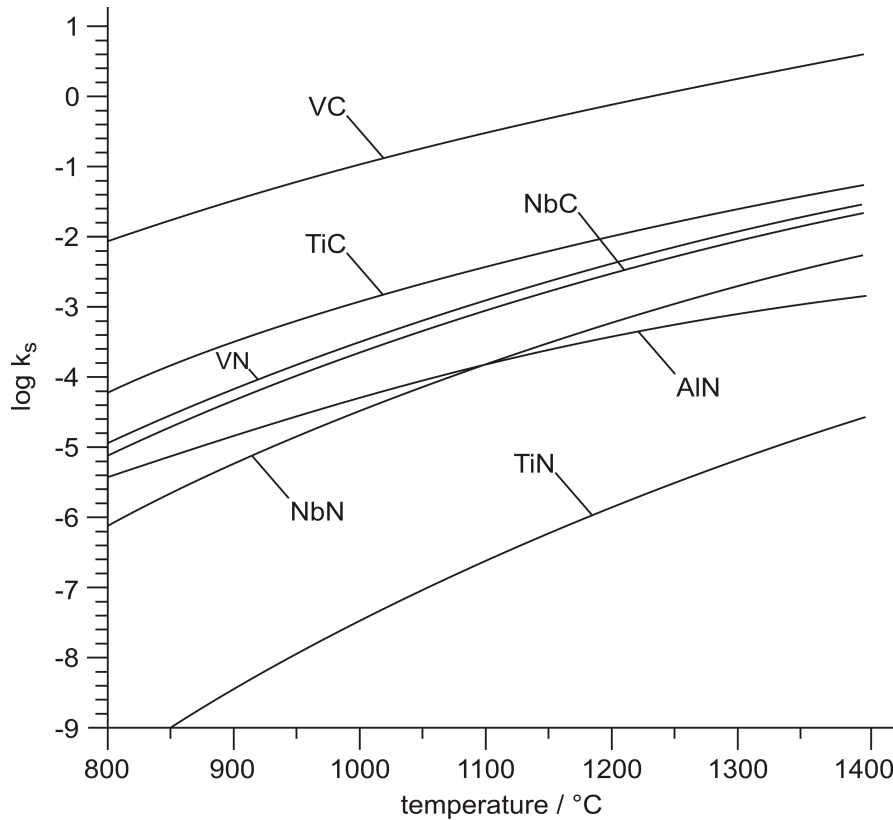


Figure 4.11: Comparison of the solubility of various carbides and nitrides, calculated with the thermodynamic database 'mc\_steel (version 1.86)'.

Due to the different solubility characteristics of the nitrides and carbides, depicted in figure 4.11, Gladman [1] proposed several important conclusions, which are summarized in section 3.2.3.2.

#### 4.2.9 Discussion

The solubility data of various nitrides and carbides have been determined by many workers using different theoretical or experimental (direct or indirect) methods. However, the evaluation of the solubility product by different authors show scatter to a greater or lesser extent, depending on the investigated phase, see section 4.2.1 to section 4.2.7.

One reason for these discrepancies can be related to the variety of methods used in ob-

taining the given solubility, as each approach has its own assumptions and limitations. Typical methods for determining solubility products of nitrides or carbides are briefly summarised in the following [115]:

- thermodynamic calculations
- chemical separation and isolation of precipitates
- equilibrating a series of steels with different concentrations of the microalloying element, after which the C or N contents are analyzed
- hardness measurements
- thermoelectric power measurements
- statistical treatment of existing solubility products

Thermodynamic calculations of solubility products often neglect any interaction between elements. Only few authors [107, 122, 134] consider these interaction parameters to account for the effect of alloying elements on the solubility of nitrides or carbides in austenite.

Experimental techniques, which indirectly account for these chemical interactions, have other limitations. One problem of the separation method is that very fine particles may not be included in the analysis. Particularly, the Beeghly method [136], which is used by numerous authors, seems to be critical for a reliable detection of fine precipitates [137]. Other methods, e.g. the equilibrating method, assume stoichiometric or nonstoichiometric compounds, not accounting for the exact chemical composition of the precipitate. Finally, indirect methods like hardness measurements or thermoelectric power measurements are questionable since they are based on the assumption that an increase in hardness or thermoelectric power is directly related to the precipitation of nitrides and/or carbides, not accounting other effects, e.g. fine grain hardening, solid solution hardening or dislocation strengthening [115].

Apart the limitations of the investigation methods, there exist other problems related to the chemical composition of the investigated alloys. Erasmus et al. [24] demonstrate that the AlN solution temperature in a 3,5 % Ni steel is about 100°C higher compared

to a plain C steel due to the lower solubility of N in austenite. Furthermore, Höner and Baliktay [29] determined three different solubility products for AlN in austenite for three different alloys, see table 4.2. And Mayrhofer [89] and Koyama et al. [124] related their solubility products to the presence of Al, Mn, Cr, Si and Ni. Furthermore the discrepancies can be related to the mutual solubility of different carbides or nitrides. The higher affinity of microalloying elements to N should be mentioned, which can be related to difficulties in the determination of various carbides, if N is present in the investigated steel. Also, the substitution of an element in the precipitate by other alloying elements, e.g. Al is replaced by Cr, is reported [45] as well as core shell structures of particles [60–62] in the presence of extremely stable compounds, e.g. TiN.

Another reason can be attributed to the presence of different crystallographic structures. Whereas the equilibrium crystallographic structure of AlN is the hexagonal wurtzite structure, in the early stages of precipitation cubic structure is observed [43,45–47]. Different crystallographic structures are also observed in the case of NbN. Mori et al. [138] measured NbN particles with cubic structure but also precipitates with hexagonal structures.

All these reasons account to a more or less extent for the discrepancies observed in the solubility data in section 4.2. However, the thermodynamic database 'mc\_steel' (version 1.86) [83] is in fair agreement with most of the reported solubility products, thus representing reasonable mean values.

### 4.3 Simple versus Complex Solution Phases

'Complex solution phases' (MX phases) with fcc structure are modeled as a part of the fcc matrix phase, according to the CALPHAD approach [70]. For the simulation of these phases, composition sets are defined with the microalloying elements as major constituents. However, these phase descriptions consider also minor amounts of other elements, this is why we call them 'complex solution phases'. The complex character of these phases can sometimes cause numerical difficulties and phase instabilities when performing precipitation kinetics calculations.

In some cases, if the chemical composition of the precipitates is rather stoichiometric, the description of the precipitation phases can be simplified to so-called 'simple solution phases'. Thus, simple phase descriptions are introduced in our database 'mc\_steel', which contain only the alloying elements of the desired phase, see section 4.2. The usage of these thermodynamic models lead to an enhancement of the stability of the calculations as well as to a reduction of computing time.

However, it is worthy to note once more that the usage of simple phases is a simplification for enhancing the stability of the calculations. Thus, before pure phases can be used, it must be clarified which phases are stable in a system by finding the miscibility gaps using the description of the complex solution phases (MX phase). Once the phases are defined, and the composition is rather stoichiometric, these complex phases can be substituted by the description of simple phases.

Figures 4.12 to 4.16 show the comparison of the calculations of three alloys (see table 4.1) using the simple phase description in 'mc\_steel' (version 1.18) as well as the description of the complex solution phase (MX phase) of this database. In addition to the calculated phase fractions of the microalloying phases (figures 4.12, 4.13 and 4.15), figures 4.14 and 4.16 show the calculated chemical composition of the complex solution phases Ti(C,N), (V,Nb)C and V(C,N) in alloys 2 and 3.

Since the thermodynamic model of AlN is stoichiometric in both cases and there are no additional alloying elements except Nb and C in steel 1 the calculated phase fraction of AlN and NbC are identically, see figure 4.12.



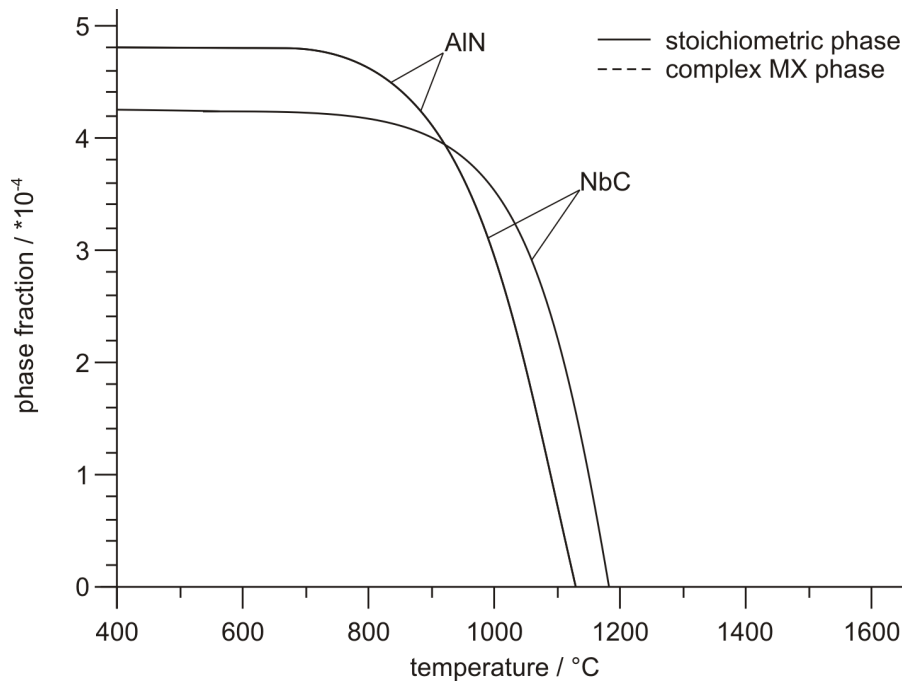


Figure 4.12: Equilibrium calculation of the phases NbC and AlN in steel 1, using the description of simple phases and complex solution phases in 'mc\_steel'.

The calculated solution temperatures of TiN in alloy 1 and 2 are in good agreement, see figures 4.13 and 4.15. However, there are some discrepancies in the phase fraction of TiN in the temperature range between 600 °C and 1200 °C, whereas the equilibrium phase fraction at low temperatures coincide well. This discrepancies can be explained by looking at the chemical compositions of the Ti(C,N) phases, see figures 4.14 and 4.16. In both cases, the Ti(C,N) precipitates are rather stoichiometric phases at high temperatures as well as at low temperatures. However, in the temperature range from 600°C to 1200°C the composition of the alloy 2 and 3 are more (Ti,V,Nb)(C,N) and (Ti,V)(C,N), respectively, rather than a simple TiN.

Using the stoichiometric VC phase description shows, that the calculated solution temperature as well as the phase fraction is in good agreement with the values obtained with the complex MX phase, in the case of steel 3 (figure 4.15). Dramatic discrepancies are only observed by the comparison of the two approaches concerning the VC in steel 2, see figure 4.13. This deviation can be attributed to the different number of stable phases, depending on the thermodynamic model. Whereas the complex phase descrip-

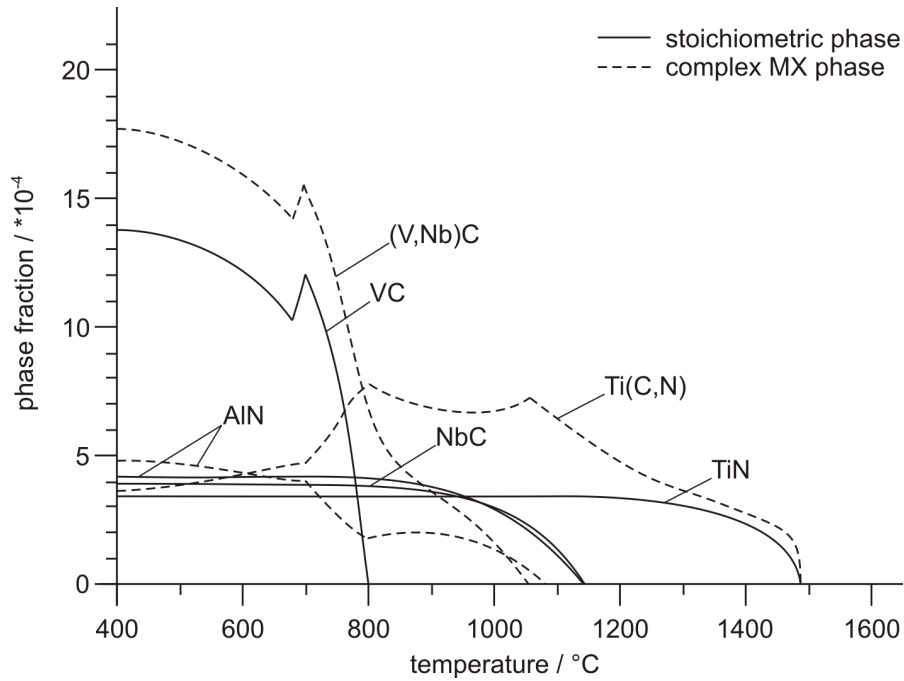


Figure 4.13: Equilibrium calculation of the phases VC, NbC, TiN and AlN in steel 2, using the description of simple phases and complex solution phases in 'mc\_steel'.

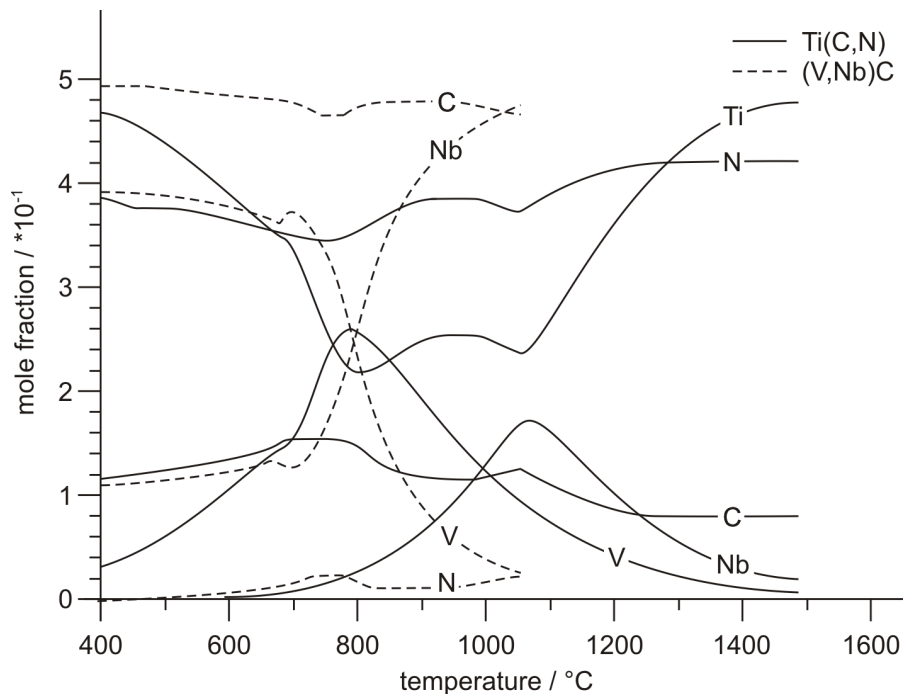


Figure 4.14: Equilibrium calculation of the chemical composition of the phases (V,Nb)C and Ti(C,N) in steel 2, using the description of complex solution phases in 'mc\_steel'.

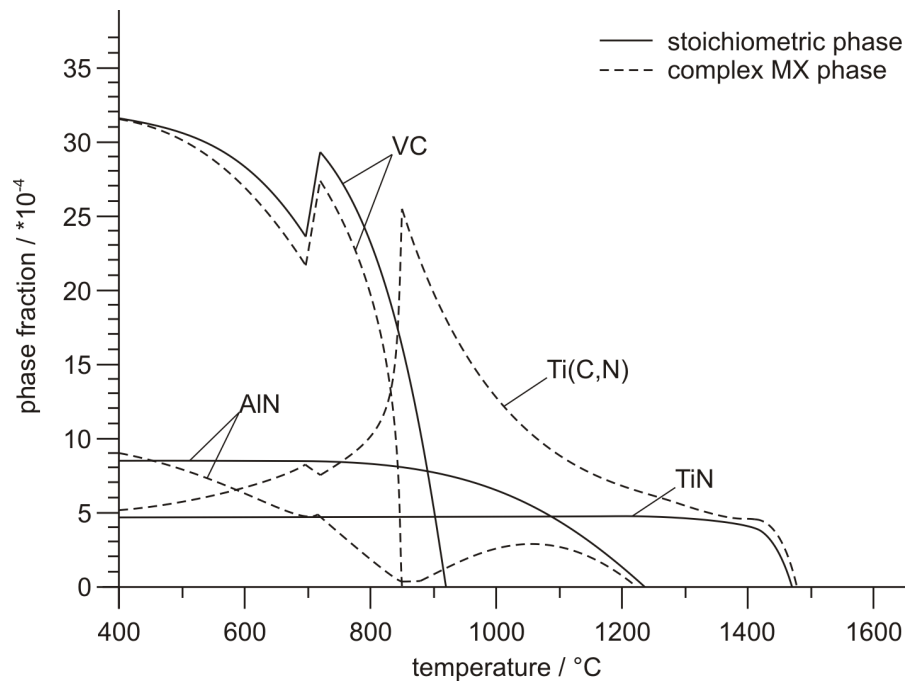


Figure 4.15: Equilibrium calculation of the phases VC, TiN and AlN in steel 3, using the description of simple phases and complex solution phases in 'mc\_steel'.

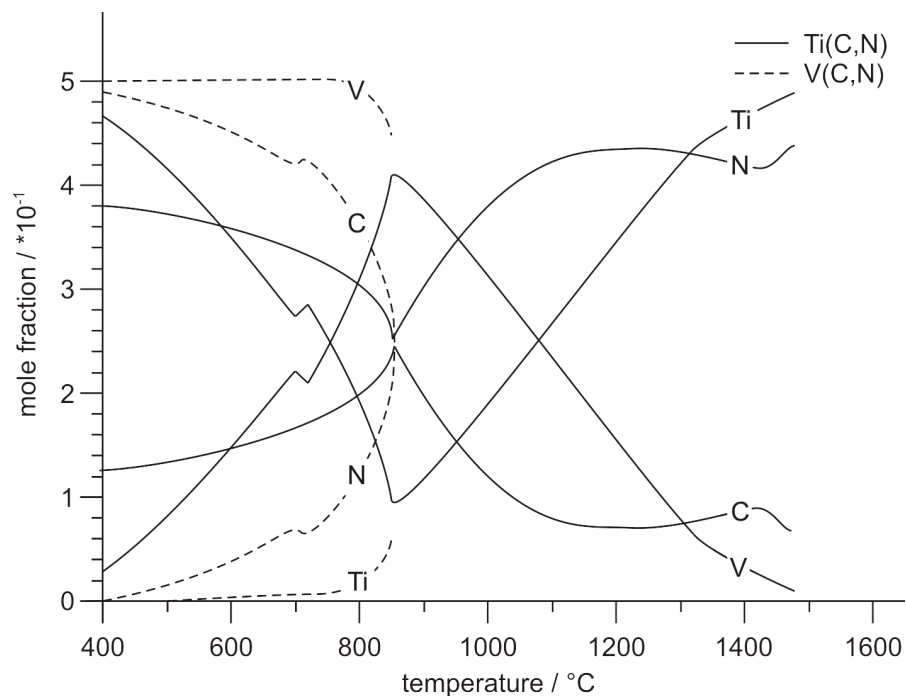


Figure 4.16: Equilibrium calculation of the chemical composition of the phases VC and Ti(C,N) in steel 3, using the description of complex solution phases in 'mc\_steel'.

tion predicts a (V,Nb)C, this phase is replaced by the two separate phases VC and NbC, when using the simple pure phases. Since the calculation of the chemical composition of the complex (V,Nb)C predicts a NbC at approximately 1000°C and a VC at temperatures below 700°C, it seems to be likely that both phases are stable depending on the temperature. In this special case it can not be clarified yet which description is more serious, since it is possible that there is a very sensitive miscibility gap which cannot be determined with the complex MX phase description.

The equilibrium calculations of the three alloys demonstrate, that the prediction of simple pure phases is in good agreement with the complex description in many cases. Thus, this simplification enables to run stable precipitation kinetics calculations.

## 4.4 Essential Ingredients for Modeling Precipitation Kinetics

For a consistent description of the precipitation kinetics of various nitrides and/or carbides, several important input parameters must be taken into account. Apart the main input quantities thermodynamic and diffusion data, it is necessary to account for several physical mechanisms that are often neglected in this type of simulations, among them the precipitate/matrix volumetric misfit and the temperature dependent Young's modulus, composition-, temperature- and size-dependent interfacial energies, as well as the ratio between bulk and grain boundary diffusion. The present section introduces these important quantities.

- **Thermodynamic input data:** The thermodynamic data is the basis of every precipitation kinetics calculation, providing chemical potentials and Gibb's free energies. Consequently the accuracy of the thermodynamic input data is of great importance, see section 4.2. For the kinetic calculations in the present work, the thermodynamic database 'mc\_steel' (version 1.50) [83] is used, unless stated otherwise.
- **Diffusion data:** The diffusion database is another very important basic ingredient, when doing precipitation kinetics calculations, providing mobility data of the elements in various phases. For the present kinetic calculations the information of the database 'mc\_sample\_fe' (version 1.03) [139] is used.
- **Interfacial energy:** The interfacial energy is calculated from the generalized  $n$ -next nearest-neighbour broken-bond approach [73], taking into account the influence of the precipitate size during nucleation [74], see also section 3.4.3.
- **Different nucleation sites:** Does precipitation occur heterogeneously at grain boundaries, dislocations or also homogeneously in the bulk? Depending on the volumetric misfit between precipitate and matrix at various nucleation sites, different types of precipitates nucleate on different nucleation sites.
- **Different diffusion geometries:** Different nucleation sites makes it necessary

to take into account different diffusion geometries. In the present work, a new model [78] for grain boundary precipitation is employed, which takes into account fast short-circuit diffusion along grain boundaries as well as slower bulk diffusion inside the grain, together with the classical treatment for randomly distributed precipitates with spherical diffusion fields.

- **Ratio between diffusion along grain boundaries  $D_{gi}$  and diffusion in the bulk  $D_{bi}$ :** Since grain boundaries represent very efficient short-circuit diffusion paths with diffusivities several orders of magnitude larger than in the bulk, it is necessary to take into account the ratio between diffusion at grain boundaries  $D_{gi}$  and in the bulk  $D_{bi}$ .
- **Grain size:** Another important input parameter is the grain size, which strongly influences the calculated results by applying the model for predominant precipitation at grain boundaries.
- **Dislocation densities:** If nucleation of precipitates occurs at dislocations, the dislocation density is of prior importance. For the present calculations dislocation densities are used according to [140].
- **Volumetric misfit:** As already mentioned, different precipitates nucleate on different nucleations sites depending on the volumetric misfit between particle and matrix. How big is the misfit between precipitate and matrix at different nucleation sites?
- **Young's Modulus:** In addition to the volumetric misfit, the calculation of the misfit strain energies requires a temperature dependent description of the Young's modulus of the fcc and bcc iron matrix phases. The precipitates are assumed to be a rigid body. The temperature dependence of the Young's modulus of the precipitates is therefore neglected.
- **Quenched in vacancies:** Depending on the heat treatment parameters (e.g. solution treatment and subsequent quenching) the effect of quenched in vacancies must not be neglected. Although, this effect is well-known and heavily utilized in aluminium alloys (e.g. [141]) it is assumed to be not relevant in the case of microalloyed steels.

These aspects are investigated in detail and descriptions are presented subsequently. Some of these descriptions are valid for the precipitation of different phases (e.g. ratio between grain boundary diffusion and bulk diffusion or temperature dependent Young's modulus), whereas others have to be studied for each phase (e.g. thermodynamic data or misfit between matrix and precipitate phase) or for each application (e.g. grain size) separately. The next sections deal with the introduction of temperature dependent functions, used independently of the precipitating phases.

#### 4.4.1 Ratio Between Grain Boundary Diffusion and Lattice Diffusion

An important ingredient for treatment of precipitation of nitrides and carbides with the new grain boundary precipitation model [78] is the grain boundary diffusional mobility.

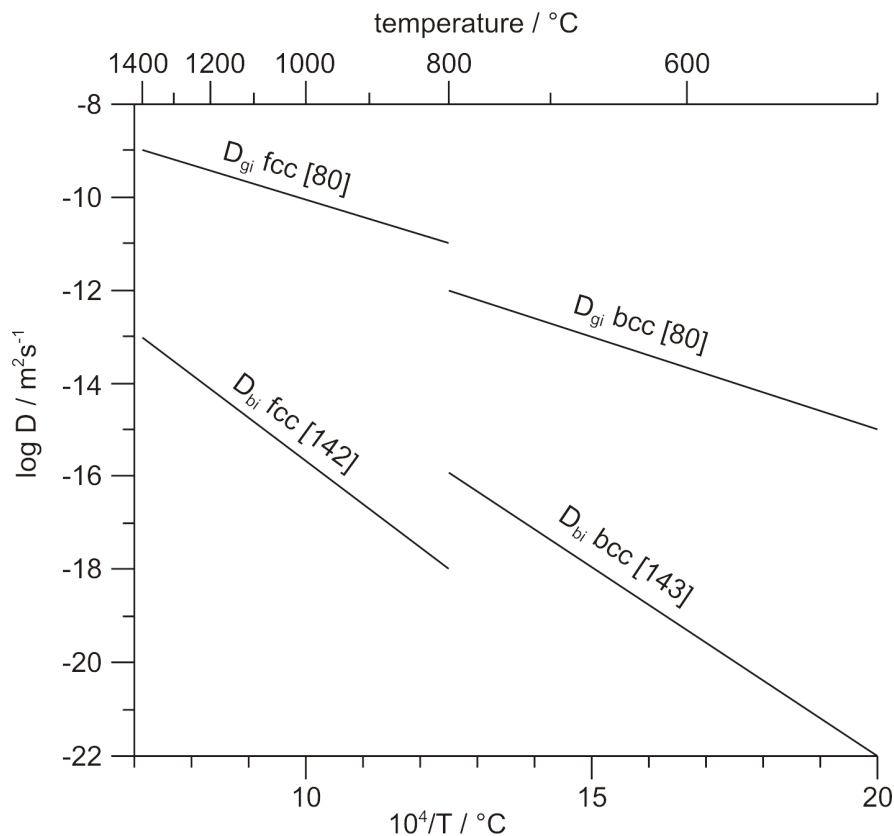


Figure 4.17: Temperature dependent self diffusion coefficients of Fe in the bulk ( $D_{bi}$ ) and at grain boundaries ( $D_{gi}$ ) [80, 142, 143].

Figure 4.17 shows the temperature dependent diffusion coefficients in the bulk and at the grain boundaries in bcc and fcc Fe. The ratios  $D_{gi}/D_{bi}$  are in the range of  $10^7$  to  $10^4$  for temperatures between 500 °C to 800 °C and 800 °C to 1400 °C, for ferrite and austenite, respectively. According to the approximated relations

$$\left(\frac{D_{gi}}{D_{bi}}\right)_{fcc} = 10^{(11-0.005\cdot T)} \quad (4.2)$$

and

$$\left(\frac{D_{gi}}{D_{bi}}\right)_{bcc} = 10^{(12-0.01\cdot T)} \quad (4.3)$$

these temperature dependent ratios  $D_{gi}/D_{bi}$  are taken into account for  $\gamma$  and  $\alpha$  iron, respectively. Both equations consider the temperature  $T$  in °C.

#### 4.4.2 Temperature Dependent Young's Modulus

Considering misfit strain energies in precipitation kinetics calculations makes it necessary to take into account the temperature dependence of the Young's modulus (see equation 3.16). Several authors have reported about this quantity in the ferrite phase field [144–153], whereas considerably less attention was paid to the austenite phase field [145, 153], see figure 4.18. Differences in the results are mainly attributed to different investigation methods [145], different chemical compositions [147] or different pre-treatment [148]. Most of the authors used the high frequency resonance method, e.g. [147, 148]. Fukuhara et al. [145] analyzed a plain low C steel using the ultrasonic pulse sing-around method. Both methods deliver accurate results, but Fukuhara et al. [145] mention in their discussion, that the ultrasonic pulse sing-around is the only technique for the determination of the elastic parameters in the higher temperature range beyond the recrystallization temperature.



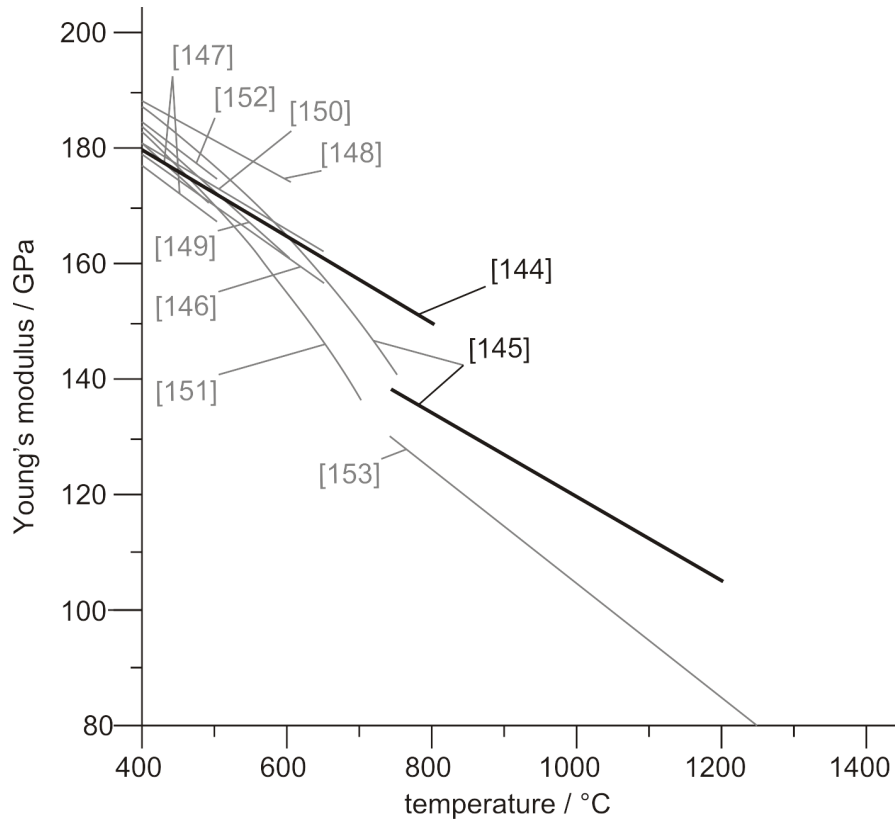


Figure 4.18: Temperature dependent Young's modulus.

In the present work, the Young's modulus of austenite is used according to

$$E_{fcc} = -75 \cdot T + 210000, \quad (4.4)$$

following the description of Fukuhara et al. [145]. The description for the ferrite phase field reads

$$E_{bcc} = -73.33 \cdot T + 193000, \quad (4.5)$$

following the results of Peil and Wichers [144], which are similar to the observations of Takeuti et al. [150], in between the lower and upper limits of Date [147] and a linear regression of the ultrasonic pulse sing-around investigations of Fukuhara et al. [145]. Also these equations for the temperature-dependent Young's modulus, accept the temperature  $T$  in °C. The relations are indicated by bold lines in figure 4.18.

## 4.5 Precipitation at Grain Boundaries

This section deals with the treatment of predominant precipitation of particles at grain boundaries, utilizing a newly developed model [78]. Therefore, the model characteristics are investigated in the present section on the example of AlN precipitation in microalloyed steels. The iron-rich corner of the Fe-Al-N system in the fcc + AlN phase field of the Fe-Al-N phase diagram represents a classical example, where nucleation of precipitates occurs predominantly at grain boundaries and is almost entirely suppressed elsewhere [26]. This model system has, therefore, been selected for investigation of the features and characteristics of the present model. The analysis is carried out in the form of a parameter study. The comparison of the calculations with experimental data will be subject of section 4.7.

### 4.5.1 Growth of Equi-Sized Precipitates

First, the growth characteristics of individual precipitates are explored utilizing an ensemble of identical particles located at the grain boundaries of unit volume of polycrystalline material with given grain radius. In this setup, the kinetics of a single particle is representative for the kinetics of the whole system.

The first set of simulations is initialized with a constant number of precipitates ( $N_0 = 10^{18} \text{ m}^{-3}$ ) with supercritical nucleation radius according to classical nucleation theory [71]. The precipitates are assumed to be homogeneously distributed over the entire grain boundary area. The simulations are carried out for different austenite grain radii between  $R = 1 \text{ }\mu\text{m}$  and  $R = 250 \text{ }\mu\text{m}$  at a temperature  $T = 1000 \text{ }^\circ\text{C}$ . The composition of the system is taken to be 0.05 wt% Al, 0.005 wt% N, balance Fe, which are typical values for microalloyed steel. It is further assumed that the diffusion of elements in the grain boundaries is  $10^4$  times faster than diffusion in the matrix.

Figure 4.19 shows the calculated growth kinetics of a constant number of equi-sized AlN precipitates ( $N_0 = 10^{18} \text{ m}^{-3}$ ) using the new grain boundary diffusion geometry (GBDG) formalism at hand. Plot (a) displays the evolution of the phase fraction of AlN precipitates versus time, whereas plot (b) shows the evolution of the precipitate radius.

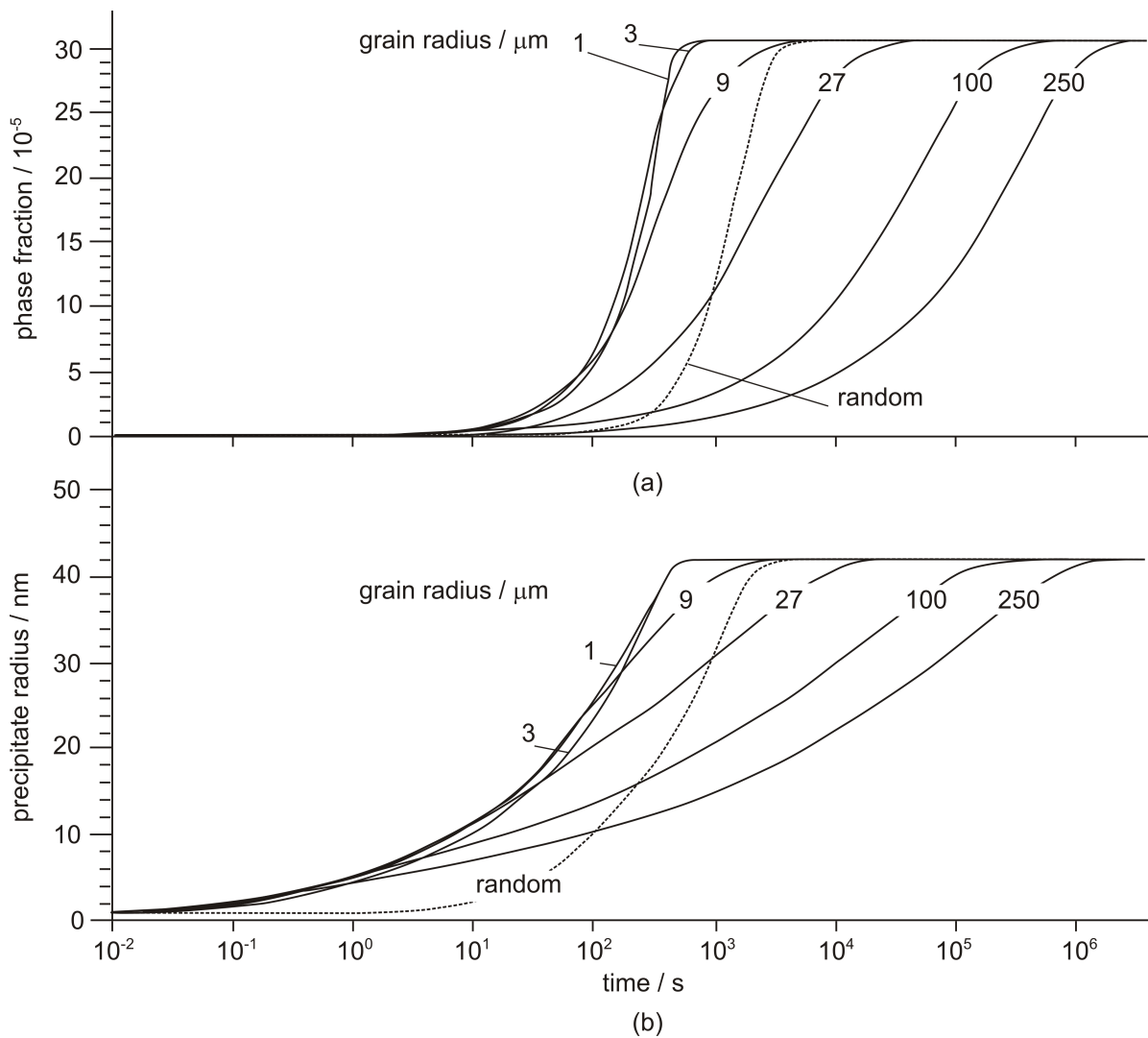


Figure 4.19: Evolution of phase fraction and radius of a constant number of identical AlN precipitates for different grain radii. Al=0.05 wt%, N=0.005 wt%, T=1000 °C,  $D_{gi} = 10000 \cdot D_{bi}$ .

In accordance with expectation, the growth kinetics of grain boundary precipitates is the faster the smaller  $R$  is. The leftmost curve in the phase fraction plot is computed for a grain radius of  $R = 1 \mu\text{m}$ . The corresponding particle density at the grain boundary is 8 precipitates per grain with a mean diffusion field radius  $\Lambda$  of approximately 700 nm. The precipitate radius evolution shows a rapid increase in the early stages of the growth process and an almost constant growth rate until soft impingement hinders further growth. If  $R$  is increased from  $R = 1 \mu\text{m}$  to  $9 \mu\text{m}$ , while holding the total

number of precipitates constant, the precipitate density at the grain boundary increases from 8 precipitates per grain to a value of 5200. This increase leads to a significant decrease of the mean diffusion field width  $\Lambda$  from 700 nm to 250 nm and, together with the increased diffusion distance from the grain center  $R$ , to a slower growth rate in the final stages of the growth process. This effect is even more pronounced, if  $R$  is further increased to  $R = 250 \mu\text{m}$ , where the decreasing growth rate is already visible for medium size precipitates. For the largest  $R$ , a precipitate density of  $9 \cdot 10^7$  precipitates per grain is observed with a mean diffusion field radius of  $\Lambda \approx 50 \text{ nm}$ .

In the phase fraction evolution shown in plot (a) of figure 4.19, the different growth kinetics with increasing grain radius and precipitate density are clearly reflected in a change of the slope of the phase fraction versus time curves. This becomes most obvious in comparing also to the classical parabolic growth behavior observed in the random-distribution spherical diffusion field geometry (RSDG). This curve is shown in dashed line for comparison.

The second set of simulations is carried out for variable grain radius and constant precipitate density of 10000 precipitates per grain. This constraint is maintained by proper adaptation of the total number of precipitates with variable grain radius. The results are summarized in figure 4.20. Image (a) displays the evolution of the phase fraction versus time, showing that the time to reach full precipitation is strongly retarded in larger grains. Since the diffusion field radius  $\Lambda$  is approximately constant in this simulation setup (one precipitate always shares the same amount of grain boundary area), this effect is mainly due to the increased diffusion distances from the grain center to the grain boundary as well as the fact that the precipitates grow to considerably larger size (mass conservation). For the six values  $R = 1, 3, 9, 27, 100$  and  $250 \mu\text{m}$  used in figure 4.20, the corresponding number densities are  $2 \cdot 10^{21}$ ,  $7 \cdot 10^{19}$ ,  $3 \cdot 10^{18}$ ,  $1 \cdot 10^{17}$ ,  $2 \cdot 10^{15}$  and  $1.5 \cdot 10^{14} \text{ m}^{-3}$ . These results reflect the well-known fact that grain boundary precipitation (viewed in terms of phase fraction) can be effectively controlled by the parent phase grain radius.

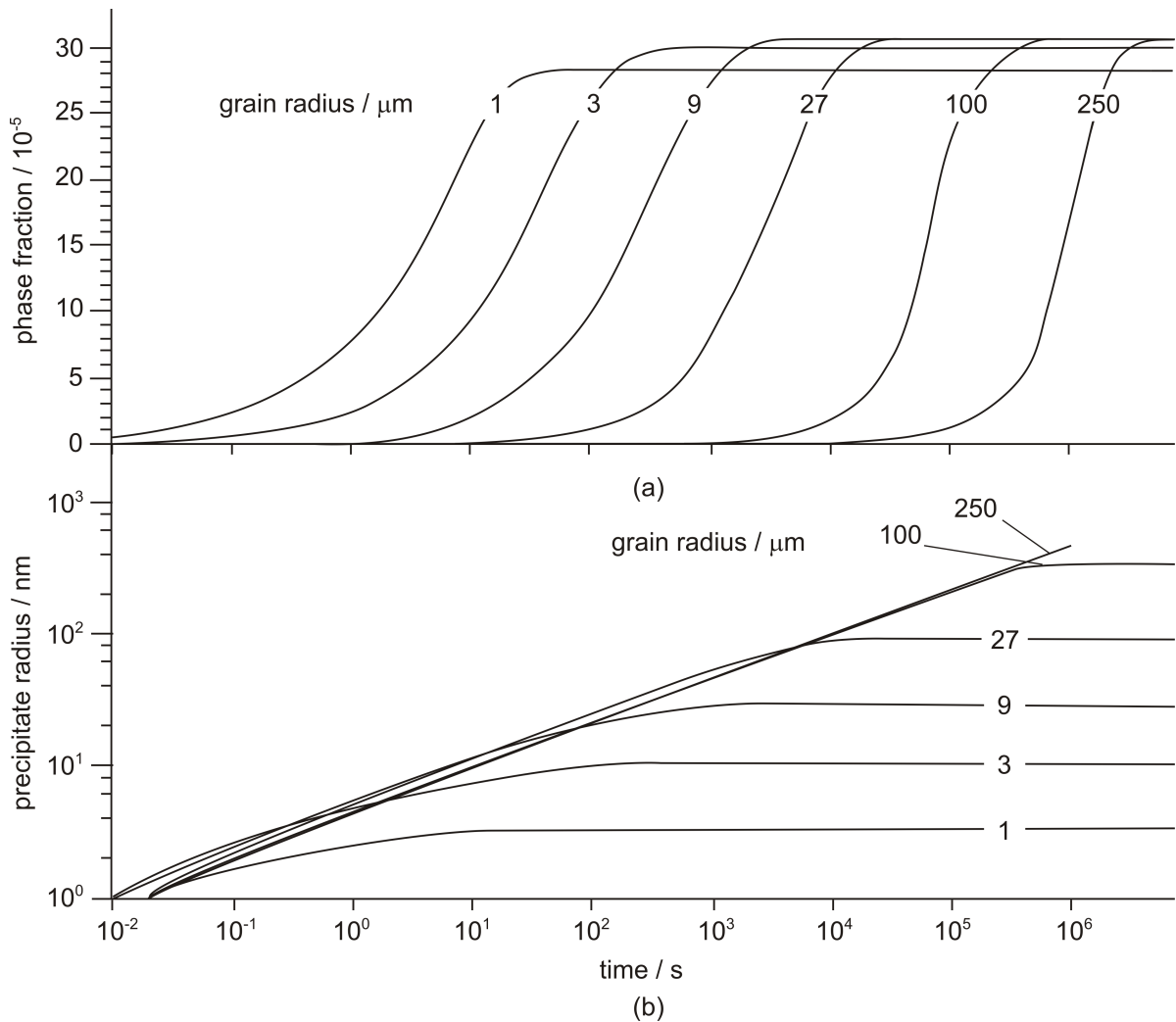


Figure 4.20: Evolution of phase fraction and radius computed in grain boundary diffusion geometry (GBDG) for constant precipitate density of 10000 precipitates per grain. Al=0.05 wt%, N=0.005 wt%, T=1000 °C,  $D_{gi} = 10000 \cdot D_{bi}$ .

#### 4.5.2 Evolution of Non-Equi-Sized Precipitates

All simulations start with a pre-defined initial size distribution of precipitates (Gaussian distribution) and a total precipitate phase fraction of 1/1000 of the equilibrium value. The simulations are carried out for a composition of 0.1 wt% Al, 0.01 wt% N, balance Fe, and a temperature T=1000 °C.

Figure 4.21 summarizes the results of the simulations for various grain radii between  $R = 1$  and  $R = 1000 \mu\text{m}$ . The major difference given by these simulation conditions

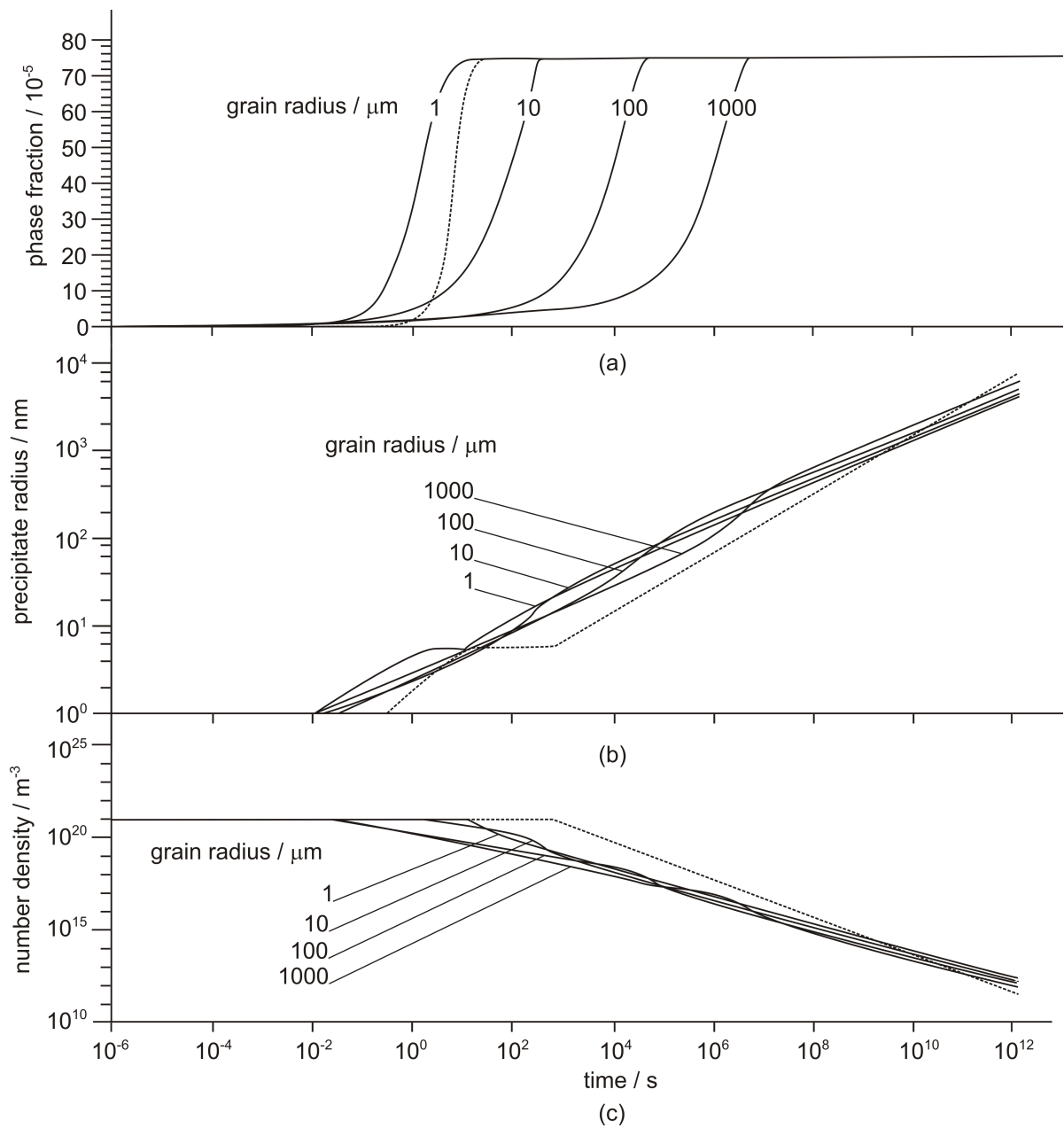


Figure 4.21: Phase fraction, mean radius and number density of AlN precipitates as a function of grain radius ( $R = 1, 10, 100, 1000 \mu\text{m}$ ). Dashed line indicates simulation with random precipitate distribution and radial diffusion fields (RSDG). Al=0.1 wt%, N=0.01 wt%,  $T = 1000 \text{ }^\circ\text{C}$ ,  $D_{gi} = 10000 \cdot D_{bi}$ .

is the average number of precipitates per grain, which varies from approximately  $10^5$  precipitates per grain for  $R = 1 \mu\text{m}$  to  $10^{15}$  precipitates per grain for  $R = 1000 \mu\text{m}$ . Investigating the evolution of the phase fraction for this simulation setup in figure 4.21,

similar characteristics are observed as in the simulations shown in figure 4.19. The evolution of the radius, however, is more complex since it involves interaction of precipitates with different sizes. The curves for the smallest grain radius of  $R = 1 \mu\text{m}$  as well as the curve for the RSDG treatment (dashed lines) exhibit a clear distinction of the growth stage from the coarsening regime. The two regions are separated by a plateau indicating the transient region after completion of growth and before onset of coarsening. For grain radii  $R = 10 \mu\text{m}$  and more, this transient region disappears due to a significant overlapping of precipitate growth and coarsening by short-circuit diffusion in the grain boundary. This observation is discussed in the following section.

Figure 4.22 displays two typical precipitate distributions observed at different stages of coarsening. Since the number of precipitates per grain changes continuously during coarsening, no stationary precipitate distribution is achieved for the entire simulation time. Instead, for sufficiently long time, a distribution corresponding to classical Ostwald ripening of grain boundary precipitates [154–159] is observed, with a time exponent of approximately  $1/4.2 - 1/4.0$  for the evolution of the mean precipitate radius. This is in good accordance with theory, predicting a value of  $1/4$  for purely grain boundary diffusion driven coarsening. The left plot in figure 4.22 shows a typical numerical distribution obtained with the present model compared to the theoretical precipitate distribution for the limit of zero phase fraction, given as [156]

$$g(u) = \left(\frac{3}{4}\right)^4 u^3 \left(1 - \frac{3u}{4}\right)^{-19/6} \left(\frac{3u^2}{16} + \frac{u}{2} + 1\right)^{-23/12} \cdot \exp\left\{\frac{1}{2} - \frac{\frac{2}{3}}{\left(\frac{4}{3} - u\right)} - \frac{1}{6\sqrt{2}} \cdot \left[\tan^{-1}\left(\frac{1 + \frac{3u}{4}}{\sqrt{2}}\right) - \tan^{-1}\frac{1}{\sqrt{2}}\right]\right\}; \quad u < \frac{4}{3}, \quad (4.6)$$

where the normalized number density  $g(u)$  is plotted over the normalized precipitate radius  $u$ .

Higher values of the time exponent up to  $1/3$  are predicted [157, 158], if solute transfer occurs by volume and grain boundary diffusion simultaneously. This occurs typically at higher temperatures. Lower values down to  $1/5$  are observed experimentally [160–162]

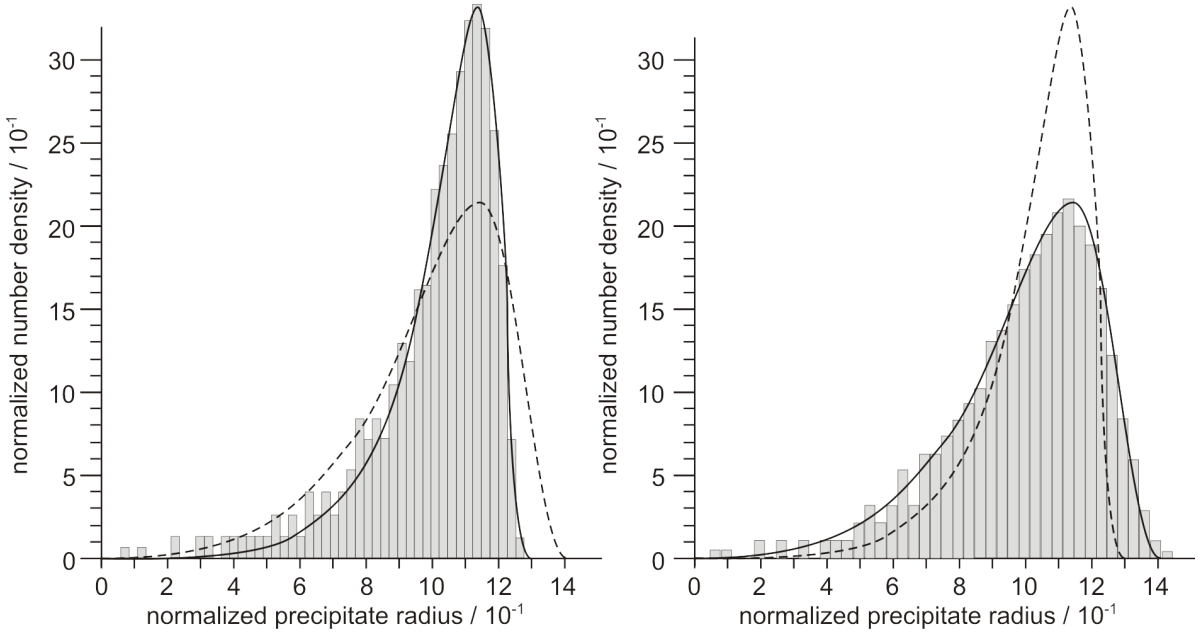


Figure 4.22: Typical precipitate distributions observed during coarsening. Left image:  $t^{1/4}$ -coarsening distribution. Right image: classical  $t^{1/3}$ -LSW distribution.

and theoretically [156], if the diffusional transport occurs via pipe diffusion through a dislocation network of, e.g., a low-angle grain boundary. In the simulations, the time exponent for the evolution of the number density is observed with approximately  $1/1.4$ . For approximately two orders of magnitude in time, the distribution changes its shape and closely resembles the classical LSW-distribution [163,164] for coarsening of randomly distributed precipitates, theoretically expressed as

$$g(u) = \frac{4u^2}{9} \left[ \frac{3}{3-2u} \right]^{-11/3} \left[ \frac{3}{3+u} \right]^{-7/3} \cdot \exp \left[ \frac{-2u}{3-2u} \right]; \quad u < \frac{3}{2}. \quad (4.7)$$

This region is characterized by a 'hump' in the mean radius and the total number of precipitates curves, as clearly visible in figure 4.21. The right plot in figure 4.22 compares a corresponding distribution with the theoretical LSW distribution. In our simulations, the particular region, where this progression occurs, is always found at the position, where the precipitate density adopts values of approximately  $10^4 - 10^5$  precipitates per grain.



### 4.5.3 Influence of Grain Boundary and Bulk Diffusion

In this section, the growth kinetics and interactions of an ensemble of grain boundary precipitates of different size are investigated. By changing the ratio between the diffusional mobilities of all elements in the bulk and the grain boundaries,  $D_{gi}/D_{bi}$ , the important influence of this quantity on the overall precipitation kinetics is explored. In contrast to the numerical analyses in the previous sections, the present simulations are carried out without pre-existing nuclei. The formation of new precipitates follows a numerical procedure based on classical nucleation theory, see section 3.4. The simulations are performed with a maximum of 500 size classes. A constant grain radius of  $R = 250 \mu\text{m}$  is used, and the ratio between grain boundary and bulk diffusional mobility of all elements  $D_{gi}/D_{bi}$  is varied between 1 and  $10^6$ . If the diffusivities in the grain boundary and the bulk are assumed to be similar, this would correspond to temperatures close to the melting point. If they are assumed to be different by several orders of magnitude, this would correspond to situations commonly observed at relatively low temperatures, when the bulk diffusion is already sluggish and most diffusional transport occurs along short-circuit diffusion paths such as dislocations and grain boundaries [80]. Figure 4.23 shows the evolution of the AlN phase fraction, the mean radius and the precipitate density for various values of  $D_{gi}/D_{bi}$  while leaving the values of  $D_{bi}$  unchanged. The maximum precipitate density is limited to approximately  $10^{20} \text{ m}^{-3}$  according to the saturation of nucleation sites. The slowest precipitation kinetics is observed for the lowest values of  $D_{gi}/D_{bi}$ , which is easily attributed to the slow transport velocity of atoms towards the precipitate. With increasing ratio  $D_{gi}/D_{bi}$ , this transport is speeded up together with the entire precipitation process. In the diagram for the phase fraction of AlN versus time (figure 4.23 (a)), the kinetics of precipitation is increased significantly, if the grain boundary diffusivity is increased starting at low values. Interestingly, at a given point, this increase comes to a stop, and only the shape of the diagram changes. The faster the diffusion in the grain boundary is, the flatter is the slope of the phase fraction increase. This observation can be understood on the basis of the evolution of the other precipitation parameters shown in figures 4.23 (b) and 4.23 (c). The curves

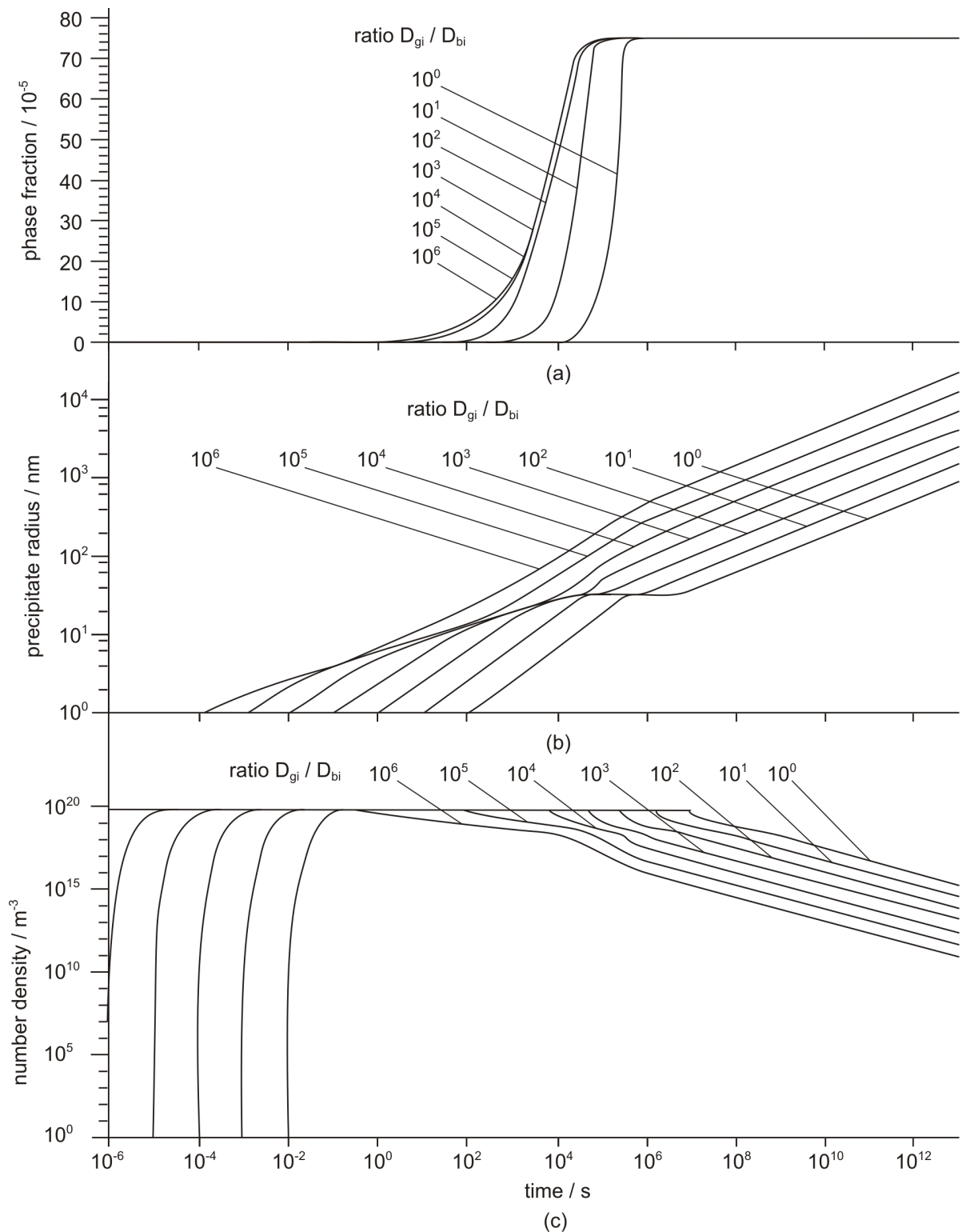


Figure 4.23: Phase fraction, mean radius and number density of AlN precipitates as a function of the ratio between grain boundary and bulk diffusion coefficient  $D_{gi}/D_{bi}$ .

for the precipitate mean radius and number density are continuously shifted to shorter times with increasing ratio  $D_{gi}/D_{bi}$ . This observation is attributed to the fact that the nucleation process of precipitates is mainly controlled by the diffusional mobility of elements within the grain boundary. With increasing  $D_{gi}/D_{bi}$ , the transport of atoms inside the grain boundary becomes faster, and growth of precipitates speeds up. An increase of  $D_{gi}/D_{bi}$  thus simultaneously shifts the start time of nucleation and the curves for mean radius and number density.

On closer investigating the evolution of the precipitate mean radius, we find that the growth of precipitates occurs approximately with a slope of 1/2 in the case of a low ratio of  $D_{gi}/D_{bi}$ . This corresponds to the well-known parabolic growth laws for diffusion-controlled transformations. If the grain boundary diffusivity becomes faster, the coarsening contribution to  $\dot{\rho}_k$  starts to dominate over the growth contribution (see section 3.4.2.2 for the two parallel processes), and the slope of the precipitate mean radius evolution starts to decrease to a value of approximately 1/3 to 1/4 for values of  $D_{gi}/D_{bi} > 10$ . A comparable decrease in the time exponent for the precipitate mean radius evolution is also observed in the case of classical precipitate coarsening [154, 158, 163, 164]. The higher the ratio  $D_{gi}/D_{bi}$  is, the earlier occurs coarsening inside the grain boundary and the smaller becomes the contribution from diffusional transport through the bulk to the grain boundaries. This fact is reflected in the phase fraction curves for higher ratio  $D_{gi}/D_{bi}$ . The curves are not shifted to shorter times. Instead, the slope of the curves changes, once a critical value of  $D_{gi}/D_{bi}$  is reached. In this case, growth and coarsening occur simultaneously, with increasing dominance of the coarsening contribution with increasing ratio  $D_{gi}/D_{bi}$ . It should be emphasized that this latter effect of change of slope occurs in addition to the effect described in the previous section based on the variation of the precipitate diffusion zone radius  $\Lambda$  in the grain boundary.

#### 4.5.4 Time-Temperature-Precipitation (TTP) Diagram

This section deals with the comparison of the time-temperature-precipitation diagram for AlN precipitation in austenite using the GBDG model and the RSDG model.

Figure 4.24 shows that the TTP plot calculated with the GBDG model deviates significantly from the TTP plot evaluated with the RSDG model. A grain radius  $R = 50 \mu\text{m}$  and a dislocation density of  $10^{12} \text{ m}^{-2}$  have been assumed in these calculations. Although both models predict approximately the same upper temperature limit for the start of AlN precipitation, the nose of the C-curve for the grain boundary diffusion model is shifted to higher temperature by almost  $100^\circ\text{C}$ . At lower temperatures, the model with radial diffusion fields predicts faster kinetics by several orders of magnitude in time for the end of precipitation. This effect is attributed to the long diffusion distances from the grain interior to the grain boundary precipitates for completion of the precipitation process.

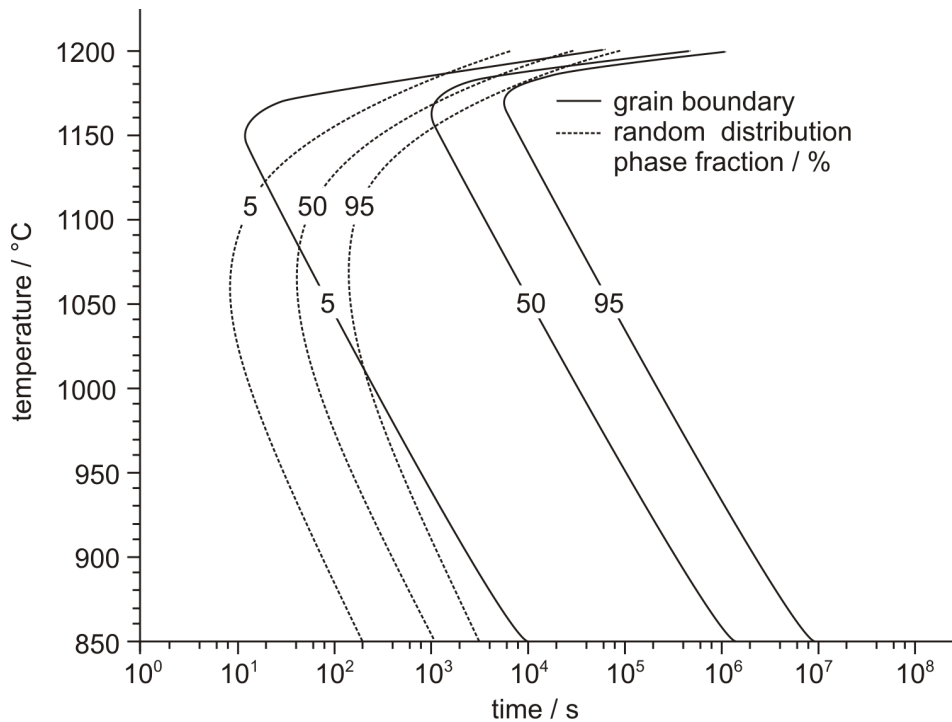


Figure 4.24: Calculated time-temperature-precipitation diagram for AlN precipitation in austenite. The solid lines represent the result using the GBDG model, while the dashed lines indicate the treatment with radial diffusion fields (RSDG model).

## 4.6 Influence of Volumetric Mismatch and Young's Modulus

Precipitation at dislocations is simulated using the classical model assuming randomly distributed precipitates surrounded by spherical diffusion fields (RSDG model), see section 3.4.2.1. Due to differences in lattice parameters, in some cases (e.g. AlN), high volumetric misfits between microalloy precipitates and the steel matrix are observed. This makes it necessary to take into account the elastic strain energy  $G_s$ . Especially, in the nucleation stage, this energy contribution can substantially reduce the total free energy change (see equation 3.15) and, thus, the nucleation rate (equation 3.11). In addition to the volumetric misfit, the Young's modulus is an important ingredient for the calculation of the elastic strain energy  $G_s$ , see equation 3.16.

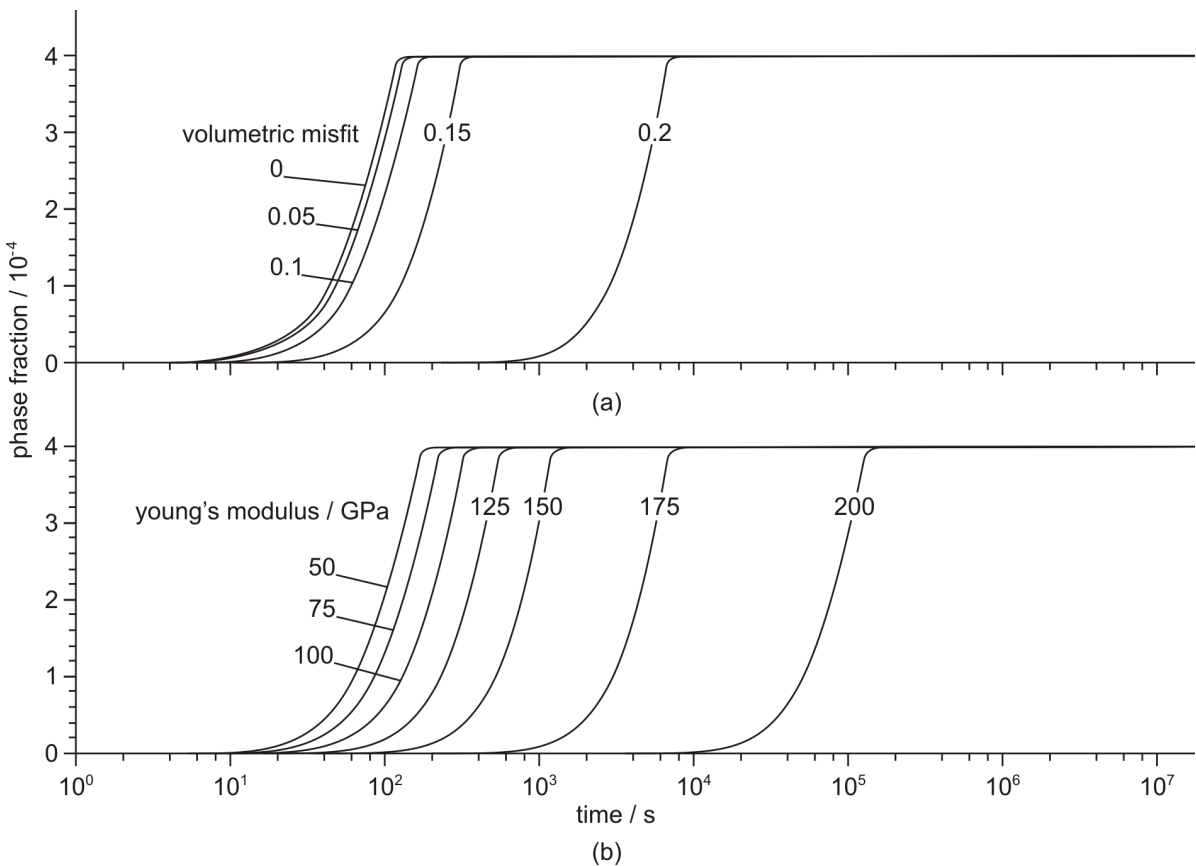


Figure 4.25: Sensitivity of the calculations of precipitation at dislocations on the parameters (a) volumetric misfit and (b) Young's modulus.

For the analysis of the sensitivity of the calculation of precipitation at dislocation on the volumetric misfit energy contribution, an alloy containing 0.05 % Al and 0.005 % N is isothermally heat-treated at a temperature of 700 °C. Figure 4.25 indicates the influence of the values of the volumetric misfit and the Young's modulus on the simulation results. The volumetric misfit is varied between 0 and 0.2 in 0.05 steps (figure 4.25 (a)) and the Young's modulus is varied between 50 GPa and 200 GPa in 25 GPa steps (figure 4.25 (b)). Both, the volumetric misfit as well as the Young's modulus show a high influence on the precipitation kinetics calculations. An increase of these parameters can retard the precipitation to longer times by several orders of magnitude. While the temperature-dependent Young's moduli of the matrix phases fcc and bcc are well established (see section 4.4.2), the volumetric misfit has to be analysed for each microalloying phase at each nucleation site.

## 4.7 Kinetics of AlN Precipitation in Microalloyed Steel

It is already mentioned in section 3.2.1.1 that Al is an essential ingredient for the production of microalloyed high strength steels. The predominant precipitation of AlN is made responsible for grain refinement, which increases strength as well as toughness, see section 3.1.3. To achieve the desired mechanical properties, it is essential to control the precipitation of AlN carefully during the production process, and thus to understand the precipitation of these important precipitates.

Several authors have reported about theoretical treatments of AlN precipitation in microalloyed steel. Duit et al. [165] carried out calculations of AlN precipitation in ferrite based on the Avrami equation. Cheng et al. [41, 95, 166] model the precipitation of AlN in microalloyed steel. In [41, 95], these authors treat the coarsening and growth stages not accounting for the nucleation process. In reference [166], they also consider the nucleation stage, but they verify their model only on a single chemical composition. Technology-oriented treatments are available for AlN precipitation during compact strip production (CSP) [93], thin slab casting and rolling (TSCR) [167] or coiling of hot strips of microalloyed steels [168]. Zolotarevsky et al. [169] as well as Kozeschnik et al. [170] investigated the competing processes of AlN precipitation and recrystallization during batch annealing of low C steel. Biglari et al. [47] model the precipitation kinetics of AlN during internal nitridation with special emphasis on the crystallographic structure. Using thermodynamic calculations, they demonstrate that, in recrystallized specimens, the precipitation of incoherent, hexagonal AlN is favoured, whereas the precipitation of coherent cubic precipitates is preferred in cold rolled specimens.

However, despite this high number of computational studies available in literature, there is no comprehensive and rigorous description of the precipitation kinetics available for AlN precipitation. The present section deals with the numerical simulation of the precipitation kinetics of AlN, taking into account that AlN precipitation can occur simultaneously at grain boundaries and on dislocations [171–175]. The computed results are compared to independent experimental data from literature, spanning a wide range of

different chemical compositions and various isothermal annealing temperatures.

### 4.7.1 Materials

The numerical simulations are compared to experimental data from nine literature sources [45, 87–89, 137, 165, 167, 176, 177]. The Al and N contents are given in table 4.9 and table 4.10 for ferrite and austenite, respectively, and sorted according to the value of their solubility product. These alloys cover a wide range of different Al and N contents representative for conventional microalloyed steel.

Table 4.9: Chemical composition of the alloys investigated in ferrite.

alloy	Al / wt%	N / wt%	Log $k_s$	$T_{sol}$ / °C	grain size / $\mu\text{m}$	ref.
1	0.046	0.0067	-3.51	1187	15	[137]
2	0.059	0.0051	-3,52	1187	10	[176]
3	0.058	0.0050	-3,54	1181	7	[45]
4	0.058	0.0038	-3,66	1150	7	[177]
5	0.030	0.0070	-3.67	1143	15	[167]
6	0.047	0.0043	-3,69	1142	10	[165]
7	0.016	0.0065	-3,98	1072	5	[176]
8	0.035	0.0029	-3,99	1072	10	[176]
9	0.060	0.0017	-3,99	1074	15	[176]
10	0.019	0.0039	-4,13	1040	7	[176]

Table 4.10: Chemical composition of the alloys investigated in austenite.

alloy	Al / wt%	N / wt%	Log $k_s$	$T_{sol}$ / °C	grain size / $\mu\text{m}$	ref.
11	0.055	0.0240	-2,88	1417	50	[88]
12	0.120	0.0057	-3,16	1327	100	[89]
13	0.079	0.0072	-3,25	1262	50	[87]
14	0.058	0.0058	-3,47	1198	100	[87]
15	0.030	0.0070	-3.67	1143	150	[167]



### 4.7.2 Weight Percent Versus Volume Percent

Most of the experimental references [45, 87–89, 137, 165, 167, 176, 177] report about measured phase fractions in weight percent or the amount of N in wt% bound to AlN, see e.g. König et al. [88]. However, performing numerical simulations with MatCalc delivers phase fractions of carbides and/or nitrides in mole percent. Therefore, the experimental data has to be transformed from measured N bounded to AlN in wt% to the phase fraction of AlN in mol%. This section deals with a short explanation on how to handle this transformation on the example of AlN.

Let us assume that a numerical simulation with MatCalc delivers a calculated phase fraction for AlN of 0.056 mol%. This means that 0.056% of all atoms are Al and N atoms. Let us further assume that AlN is a stoichiometric compound. Thus, 0.056 mol% AlN consist of 0.028 mol% N and 0.028 mol% Al. The major amount of atoms in a microalloyed steel is represented by Fe. Therefore, the atomic weight of the measured N ( $14 \text{ g mol}^{-1}$ ) has to be related to the atomic weight of Fe ( $56 \text{ g mol}^{-1}$ ). Finally the relation of calculated AlN in mol% and measured N in wt% bound to AlN can be expressed as

$$AlN [\text{mol}\%] = \frac{m_{Fe}}{m_N} \cdot x_{stoch} \cdot N[\text{wt}\%], \quad (4.8)$$

where  $m_{Fe}/m_N$  represents the differences in atomic weight and  $x_{stoch}$  takes into account the stoichiometric compound. Thus a factor of 8 ( $4 \cdot 2$ ) must be considered when comparing the calculated phase fraction of AlN in mol% to the measured amount of N in wt% bound to AlN.

### 4.7.3 Simulation Setup

Precipitation of AlN in ferrite and austenite is assumed to occur simultaneously on grain boundaries and dislocations. Therefore, two separate populations of AlN precipitates are considered in the simulations, different in the type of nucleation site and the value for the volumetric misfit. Figure 4.26 shows the volumetric misfit (solid line) along with

the molar volumes of the bcc, fcc and AlN phases (dashed lines) versus temperature.

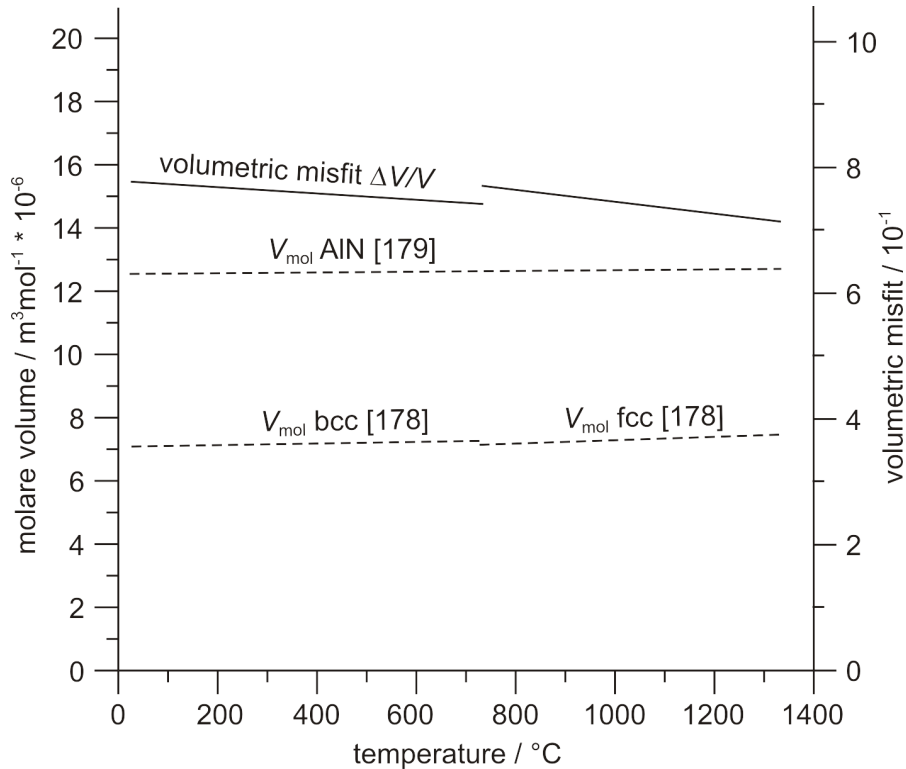


Figure 4.26: Volumetric misfit of the phases bcc, fcc, AlN and the calculated volumetric misfit [178, 179].

Accordingly, the misfit is relatively constant over the entire temperature range and varies between 71 % and 77 %. However, these values represent the precipitate/matrix mismatch in an undisturbed crystal lattice and represent the so-called unconstrained misfit, which corresponds to the difference in lattice parameters of isolated bcc, fcc and AlN phases. In a precipitate/matrix composite, the matrix and precipitates will be elastically stressed, which leads to a so-called constrained misfit [180]. This effect reduces the effective volumetric mismatch that must be inserted for calculation of the misfit energy. Moreover, the misfit strain energy between AlN nuclei and matrix is strongly influenced by lattice defects, such as dislocations and grain boundaries. For the former, precipitates can accommodate some of the volumetric misfit strain in the dislocation regions with either compressive or tensile local stresses. For the latter, we assume that the grain boundary represents a heavily disturbed lattice region with almost amorphous

structure. The precipitate nuclei can form with incoherent phase boundary, where lattice vacancies can easily be created or annihilated, and relaxation of misfit strains [181] can occur almost simultaneously with nucleus formation.

The simulations have shown that a constant value of 19\* % effective volumetric misfit for precipitates located at dislocations delivers good results for the calculation of precipitation of AlN in ferrite as well as in austenite. No volumetric mismatch is considered for the precipitates at grain boundaries due to the assumed immediate stress relaxation caused by vacancy creation or annihilation at the precipitate/matrix interface. Using a value of 70 % mismatch for coherent bulk nucleation entirely suppresses the nucleation of AlN in all situations considered in the present work, consistent with experimental evidence. For the calculation of the misfit strain energy, temperature dependent Young's moduli are used according to equations 4.4 and 4.5 for austenite and ferrite, respectively. For the calculations of precipitates at dislocations, dislocation densities of  $10^{12} \text{ m}^{-2}$  and  $10^{11} \text{ m}^{-2}$  are assumed for annealed ferrite and austenite, respectively [140].

Furthermore, figure 4.17 shows the temperature dependent ratios between diffusion at grain boundaries and in the bulk  $D_{gi}/D_{bi}$ , used for grain boundary precipitation [78]. Equations 4.2 and 4.3 are incorporated in the present simulation setup.

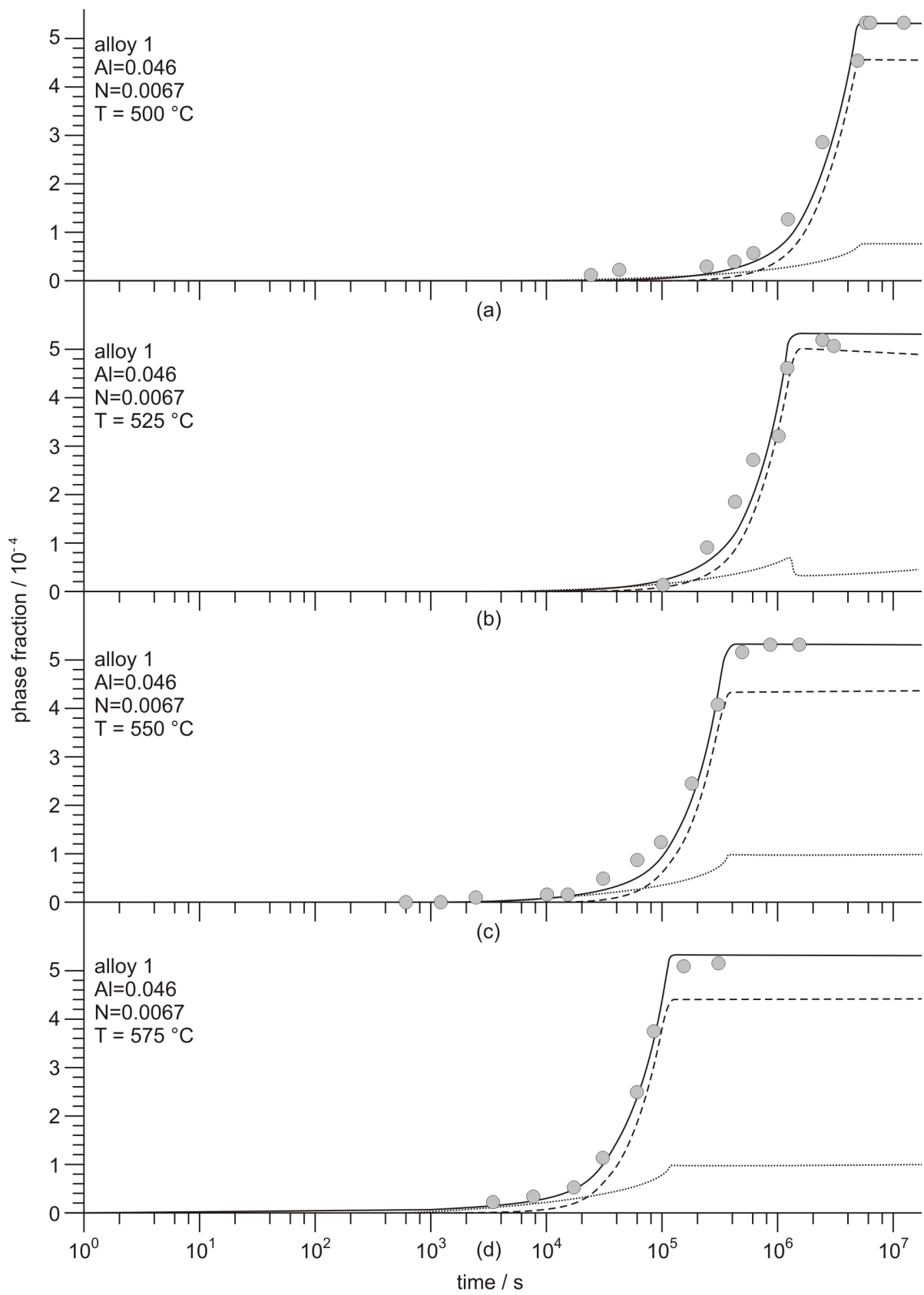
The only variable input parameter in the simulations is the grain size, which is unfortunately widely unknown for the various experiments. The grain sizes assumed in the calculations are summarized in table 4.9 and table 4.10 and are between 5 and 15  $\mu\text{m}$  for ferrite and 50 and 150  $\mu\text{m}$  for austenite. All simulations presented subsequently are otherwise performed with the same identical set of simulation parameters.

#### 4.7.4 Phase Fraction of AlN Versus Time

The following figures show the calculated phase fraction versus time curves of AlN precipitated in 15 different alloys at various temperatures.

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\*The effective volume misfit of AlN at dislocations, originally determined with 0.19 in ref. [171] and as used in this thesis, must be replaced by 0.27 when using MatCalc versions 5.40 and higher. Due to a software bug, the misfit value has erroneously been taken twice in older versions.



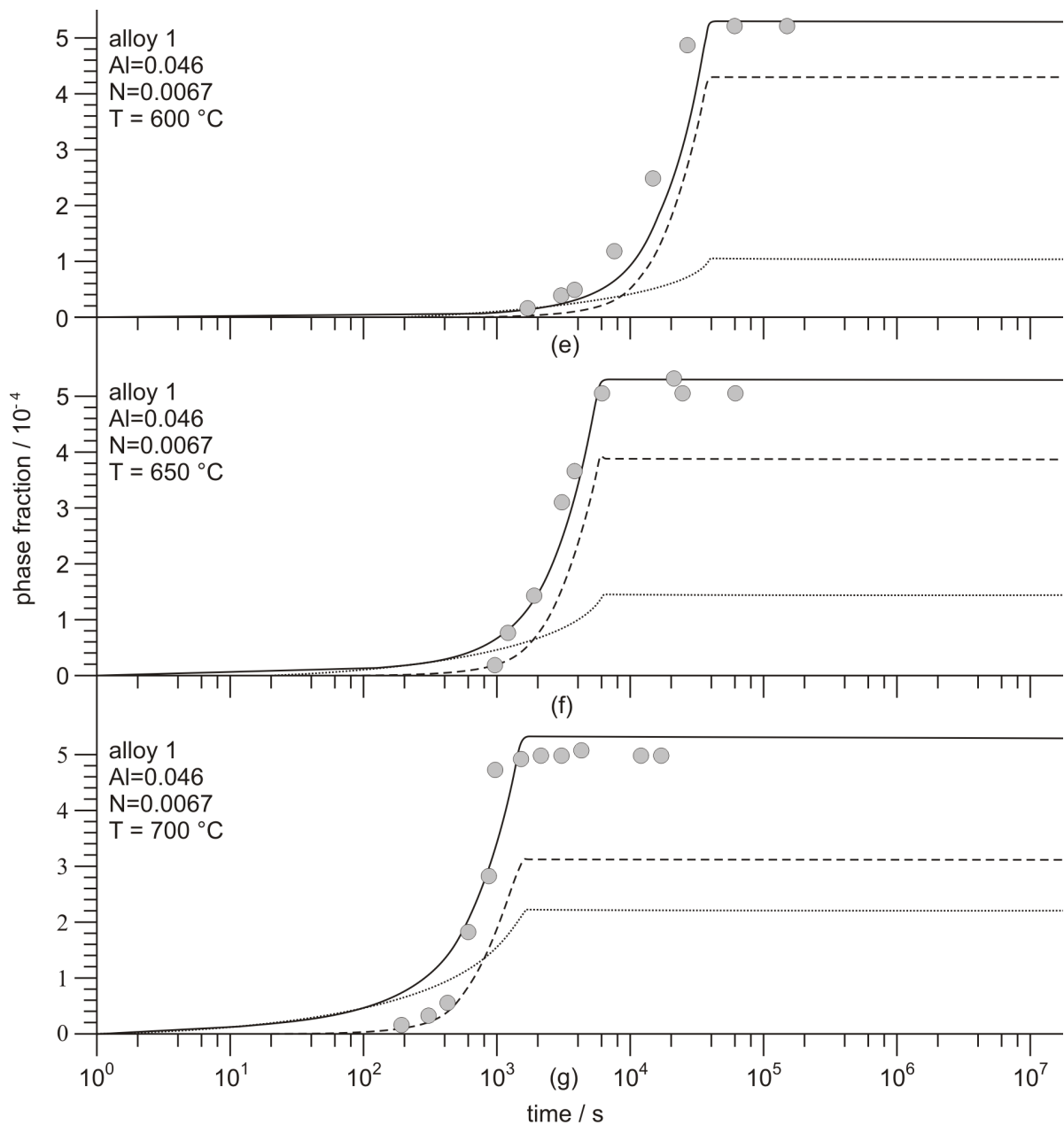


Figure 4.27: Calculated and experimental precipitation kinetics of AlN in ferrite for alloy 1 at several temperatures.

For ten alloys, precipitation occurred in the ferrite phase field, see figures 4.27 - 4.36. The simulation of the precipitation kinetics in the austenite phase field is shown in figures 4.37 - 4.41, using five different chemical compositions.

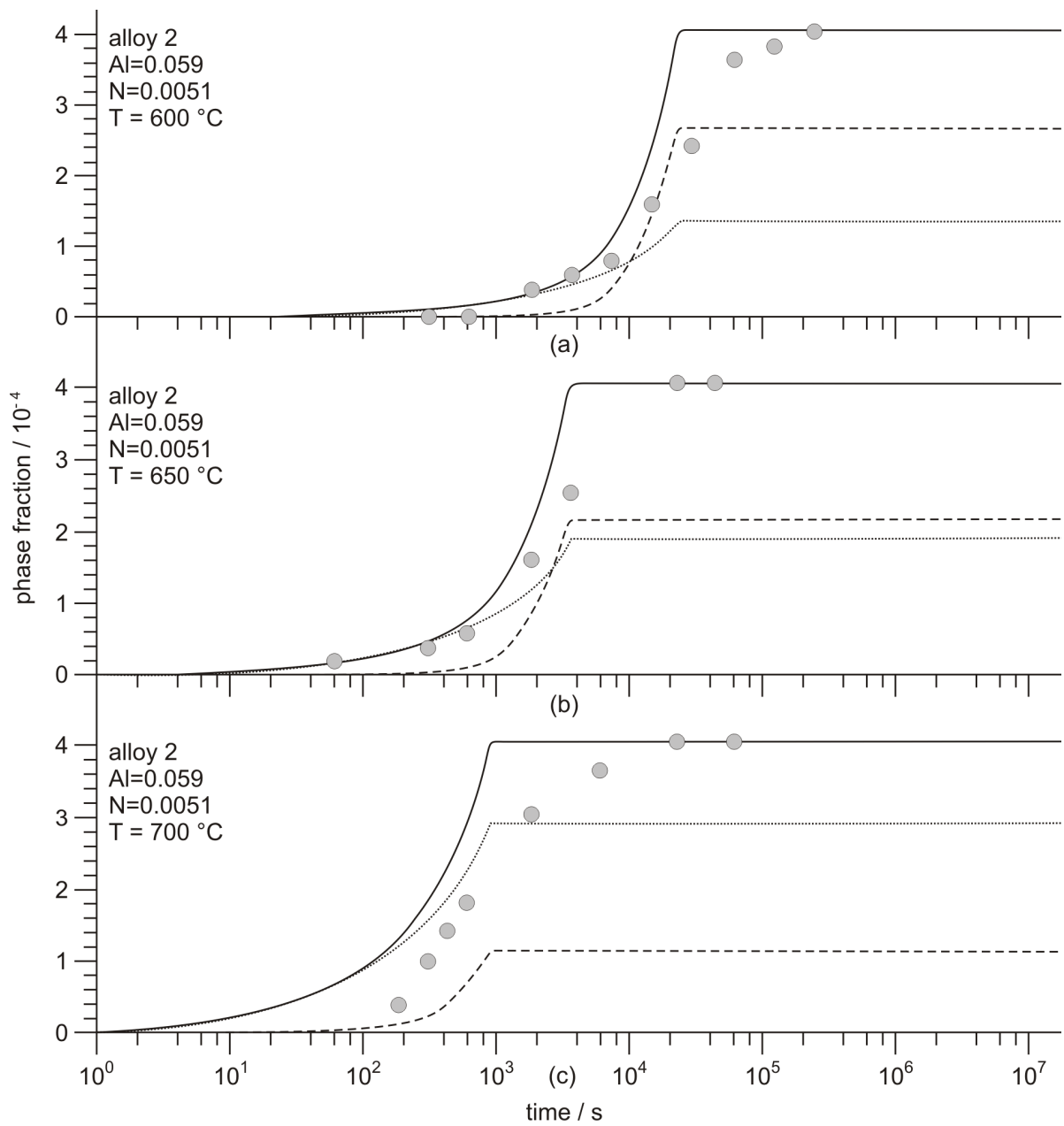


Figure 4.28: Calculated and experimental precipitation kinetics of AlN in ferrite for alloy 2 at the temperatures 600 °C, 650 °C and 700 °C.

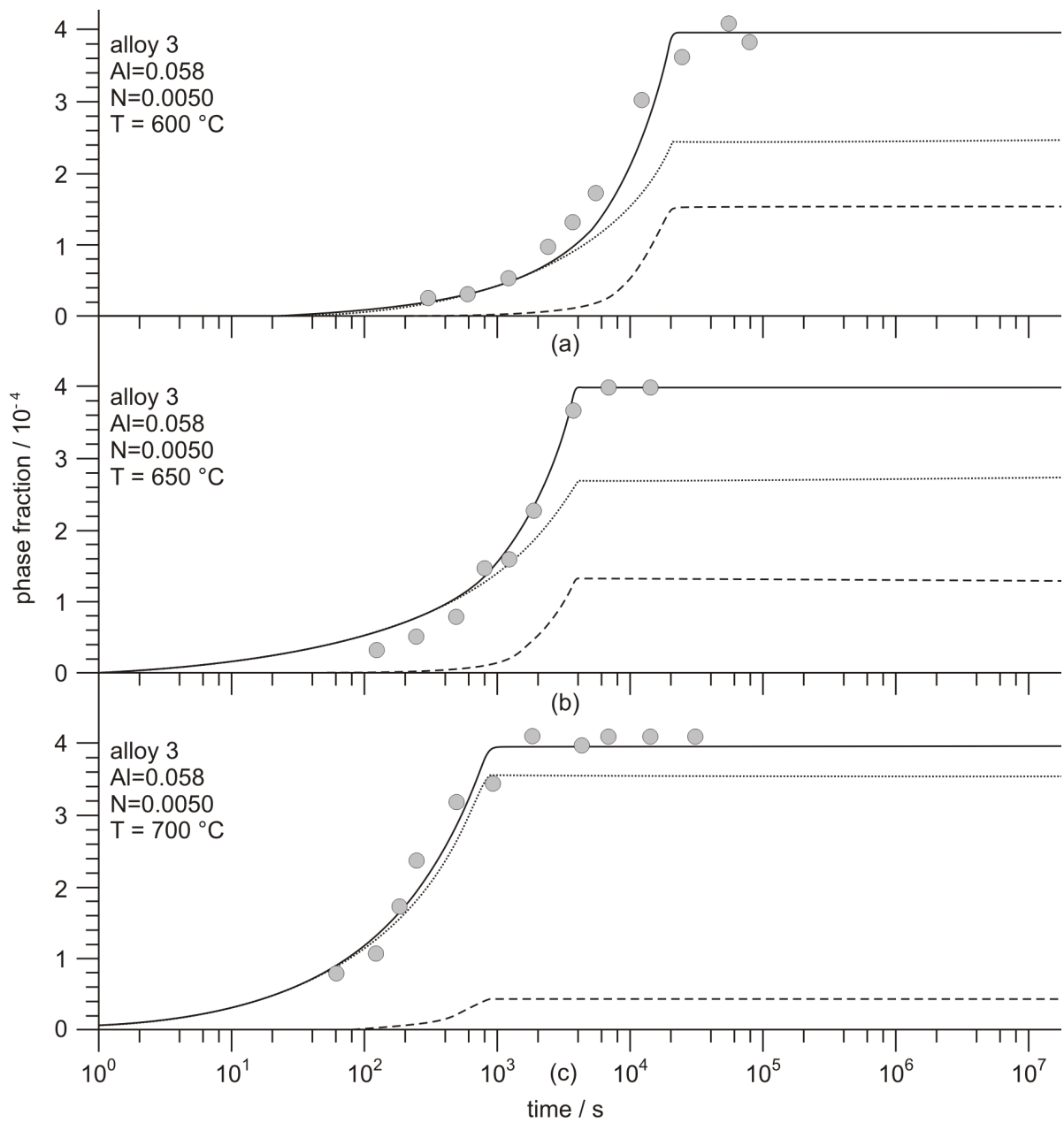


Figure 4.29: Calculated and experimental precipitation kinetics of AlN in ferrite for alloy 3 at the temperatures 600 °C, 650 °C and 700 °C.

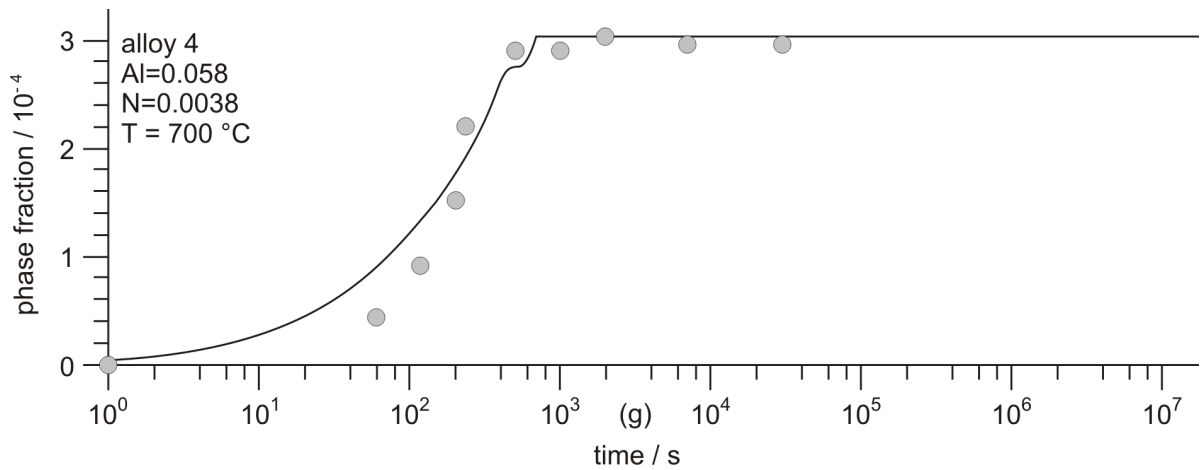


Figure 4.30: Calculated and experimental precipitation kinetics of AlN in ferrite for alloy 4 at a temperature of 700 °C.

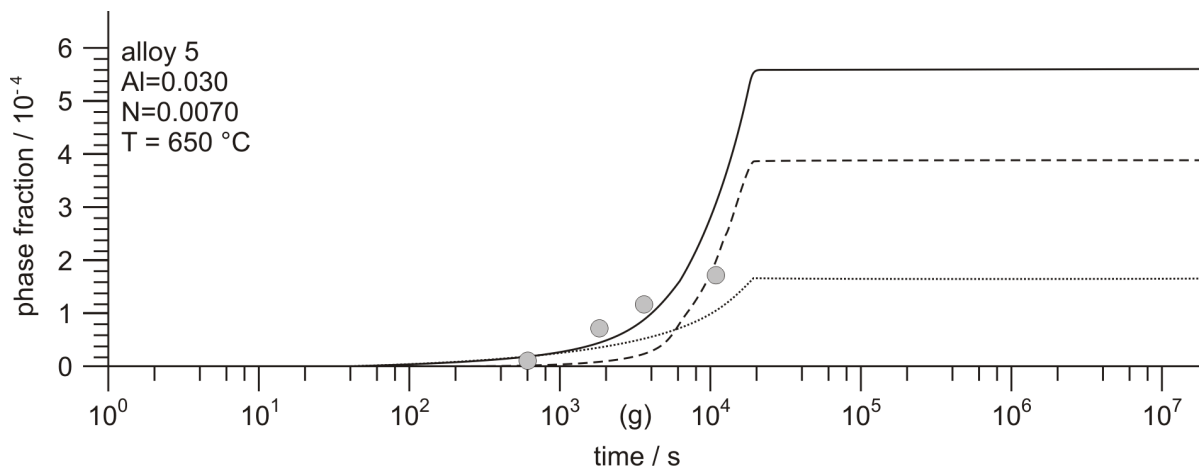
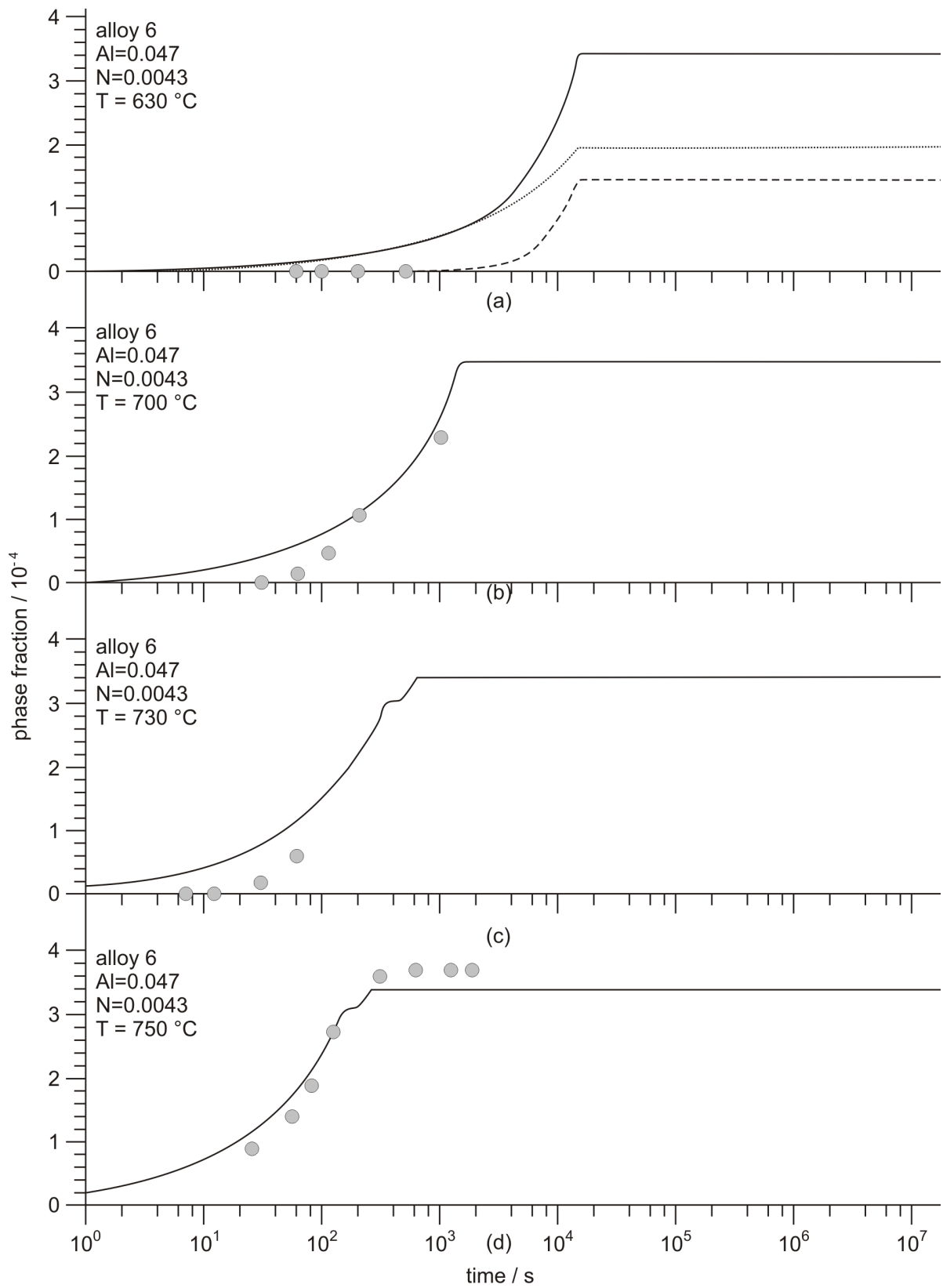


Figure 4.31: Calculated and experimental precipitation kinetics of AlN in ferrite for alloy 5 at a temperature of 650 °C.

The grain sizes used for the calculations are given in table 4.9 and table 4.10. The dashed lines represent precipitates formed at dislocations, whereas the dotted lines represent precipitates formed at grain boundaries. The sum of both phase fractions corresponds to the solid lines. Again, the alloys are sorted according to decreasing Al and N contents, which decrease from alloy 1 to 10 and from alloy 11 to 15 for the calculations in ferrite and austenite, respectively.





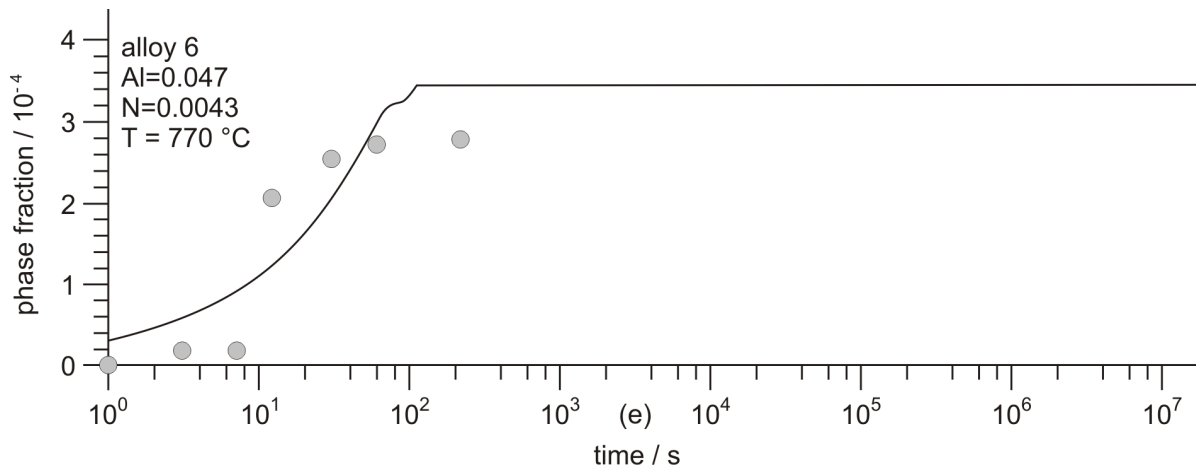


Figure 4.32: Calculated and experimental precipitation kinetics of AlN in ferrite for alloy 6 at at several temperatures.

Depending on chemical composition and annealing temperature, the calculations show only precipitation at grain boundaries or simultaneous precipitation at grain boundaries and dislocations. Especially in alloys with high Al and N contents, precipitation is observed at both nucleation sites, which is shown in figures 4.27 - 4.31, 4.32(a), 4.33(a), 4.34(a), 4.35(a) for precipitation in ferrite, and in figures 4.37(a,b) for precipitation in austenite.

Moreover, a strong dependency of nucleation site selection on the isothermal annealing temperature is observed. Whereas, at low temperatures and with higher driving forces, the majority of precipitates is located at dislocations, at higher annealing temperatures, precipitation at grain boundaries becomes predominant, see figures 4.27 - 4.29. For steels with low Al and N additions (see alloy 6-10), which have solubility products smaller than approximately  $-3.7$ , a clear preference for precipitation at ferrite grain boundaries is observed, see figures 4.32 - 4.36.

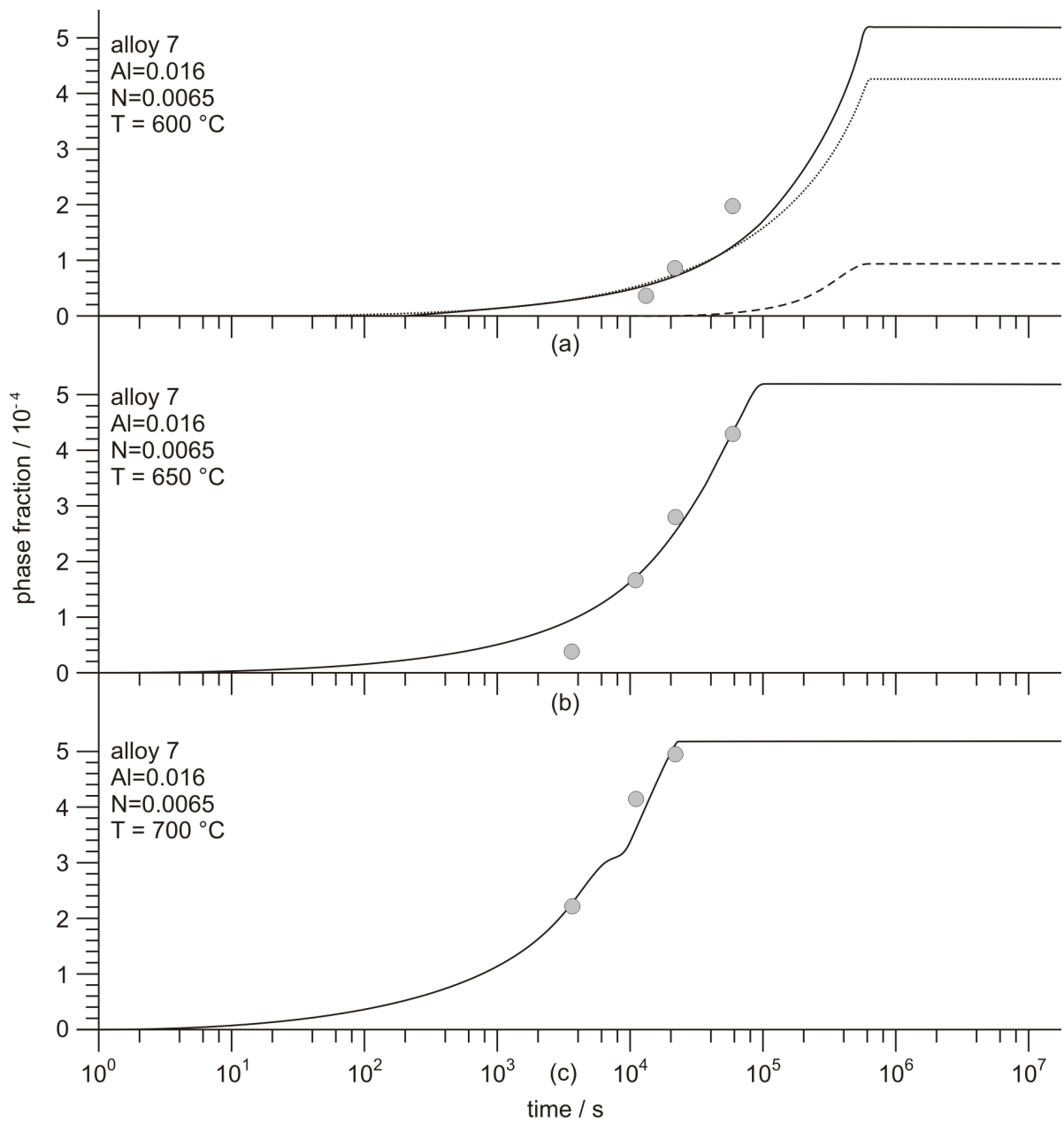


Figure 4.33: Calculated and experimental precipitation kinetics of AlN in ferrite for alloy 7 at the temperatures 600 °C, 650 °C and 700 °C.

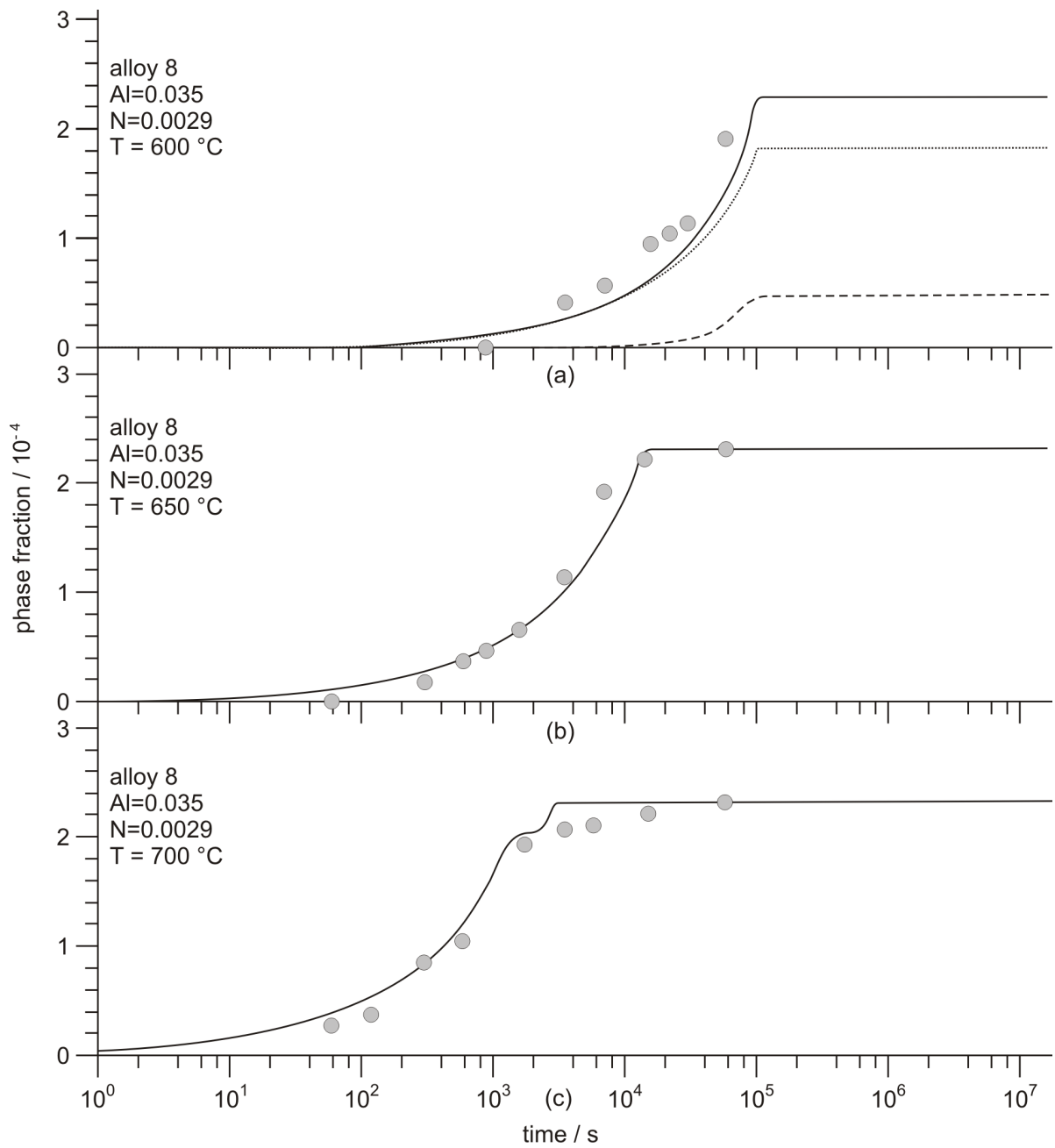


Figure 4.34: Calculated and experimental precipitation kinetics of AlN in ferrite for alloy 8 at the temperatures 600 °C, 650 °C and 700 °C.

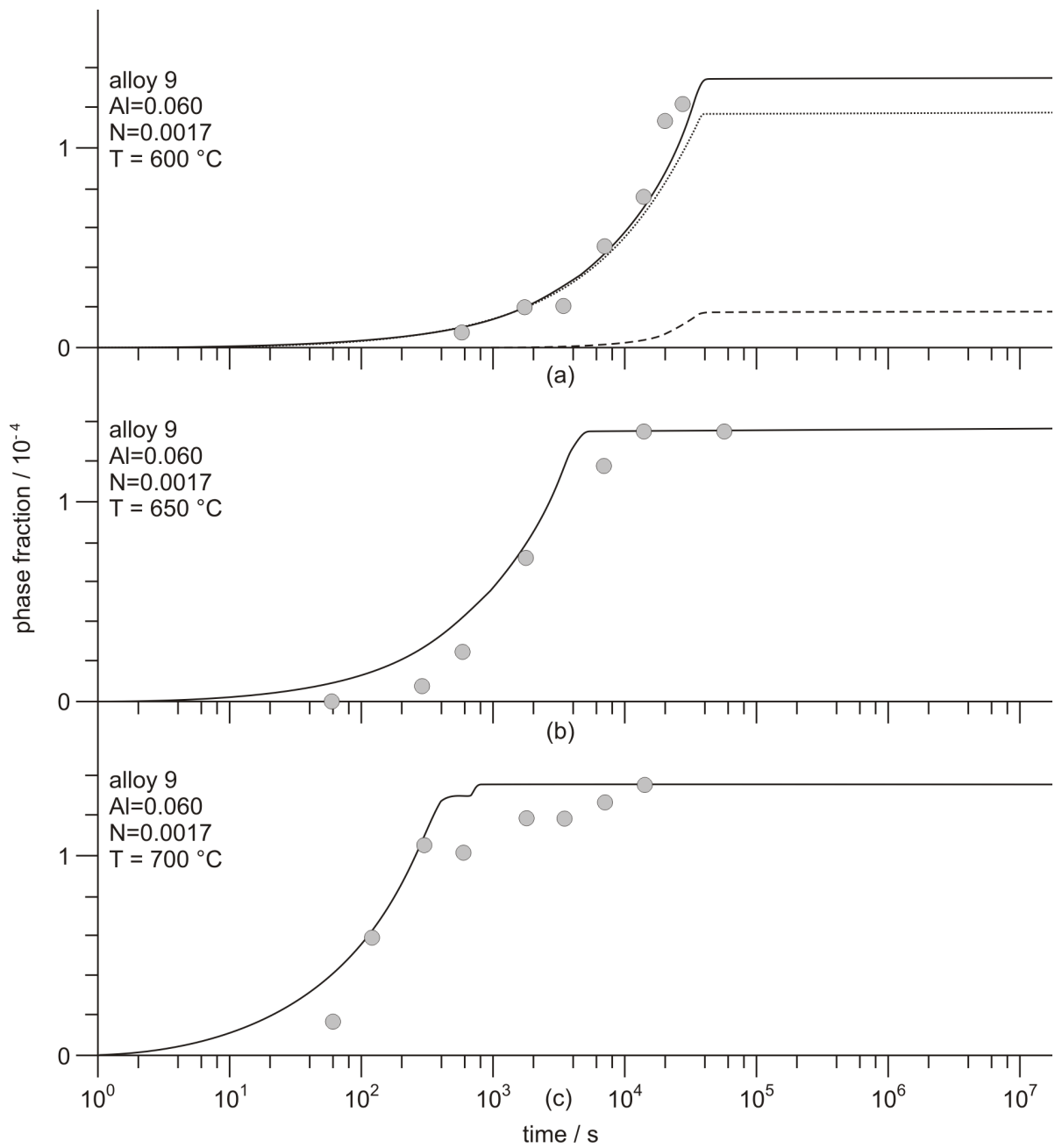


Figure 4.35: Calculated and experimental precipitation kinetics of AlN in ferrite for alloy 9 at the temperatures 600 °C, 650 °C and 700 °C.

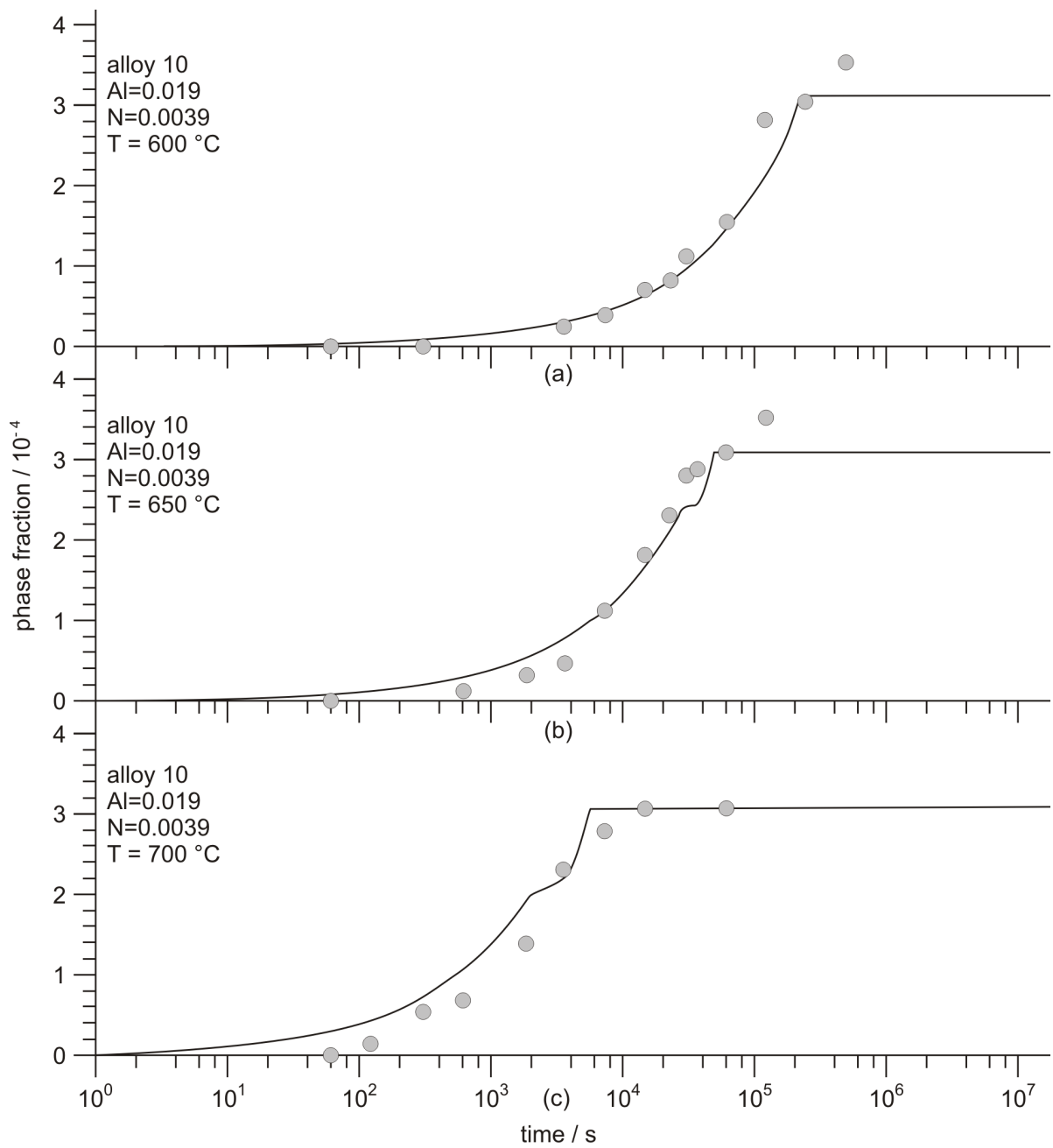
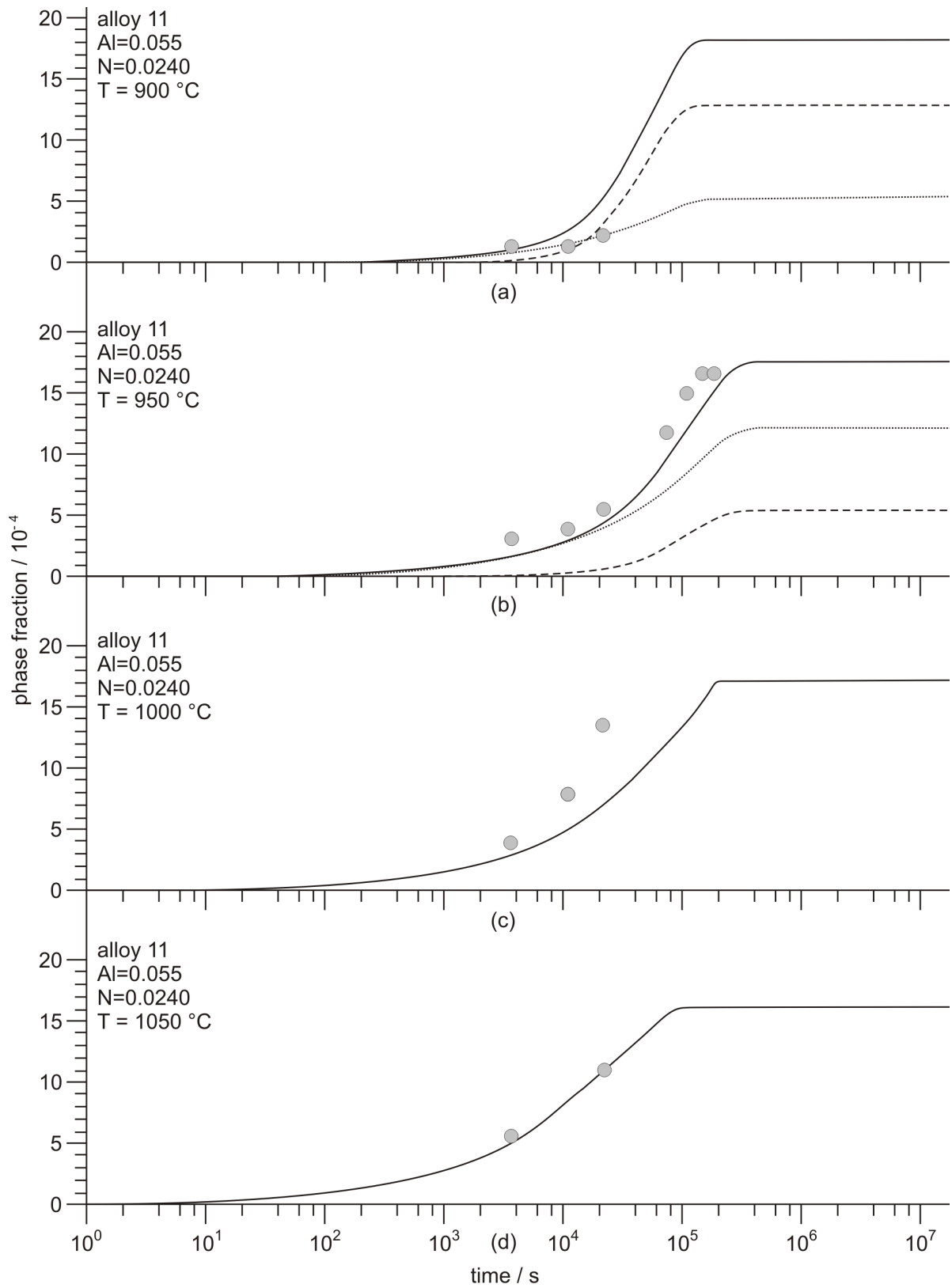


Figure 4.36: Calculated and experimental precipitation kinetics of AlN in ferrite for alloy 10 at the temperatures 600 °C, 650 °C, 700 °C.



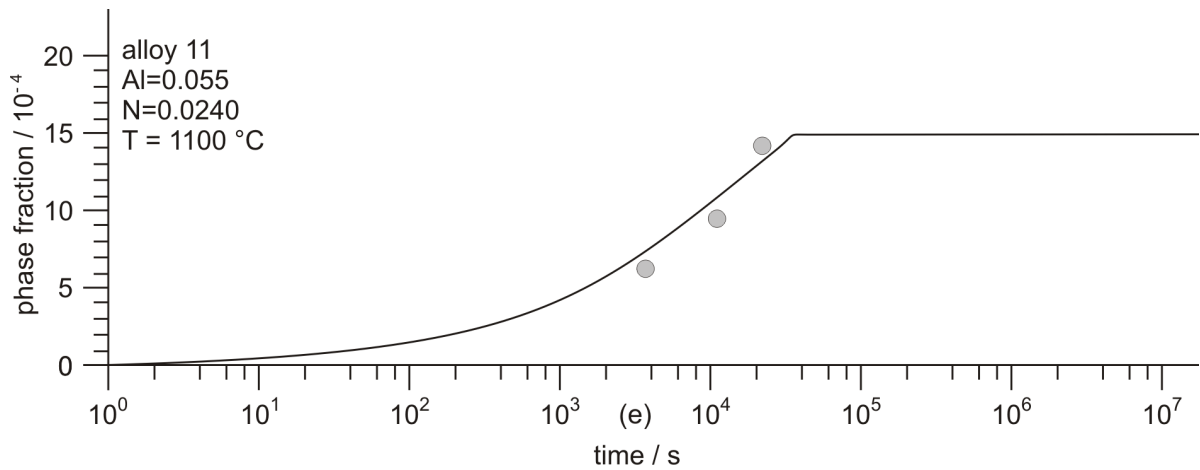


Figure 4.37: Calculated and experimental precipitation kinetics of AlN in austenite for alloy 11 at at several temperatures.

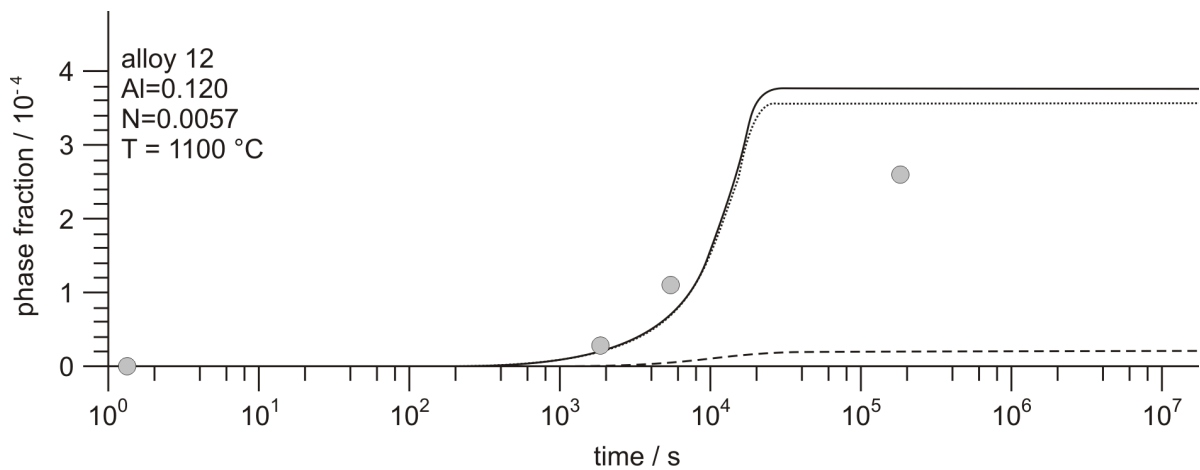


Figure 4.38: Calculated and experimental precipitation kinetics of AlN in austenite for alloy 12 at a temperature of 1100 °C.

This observation is even more pronounced in the case of precipitation in austenite, in agreement with thermodynamic calculations of Cheng et al. [166], who have also showed that  $G^*$  is much higher for homogeneous nucleation compared to nucleation at grain boundaries. These authors concluded that AlN nucleates predominantly at austenite grain boundaries. Apart from alloy 11, this is also observed in this work, see figures 4.37 - 4.41. However, alloy 11 has a very high solubility product of -2.88. This leads to the predicted additional precipitation at dislocation for temperatures below 950 °C.



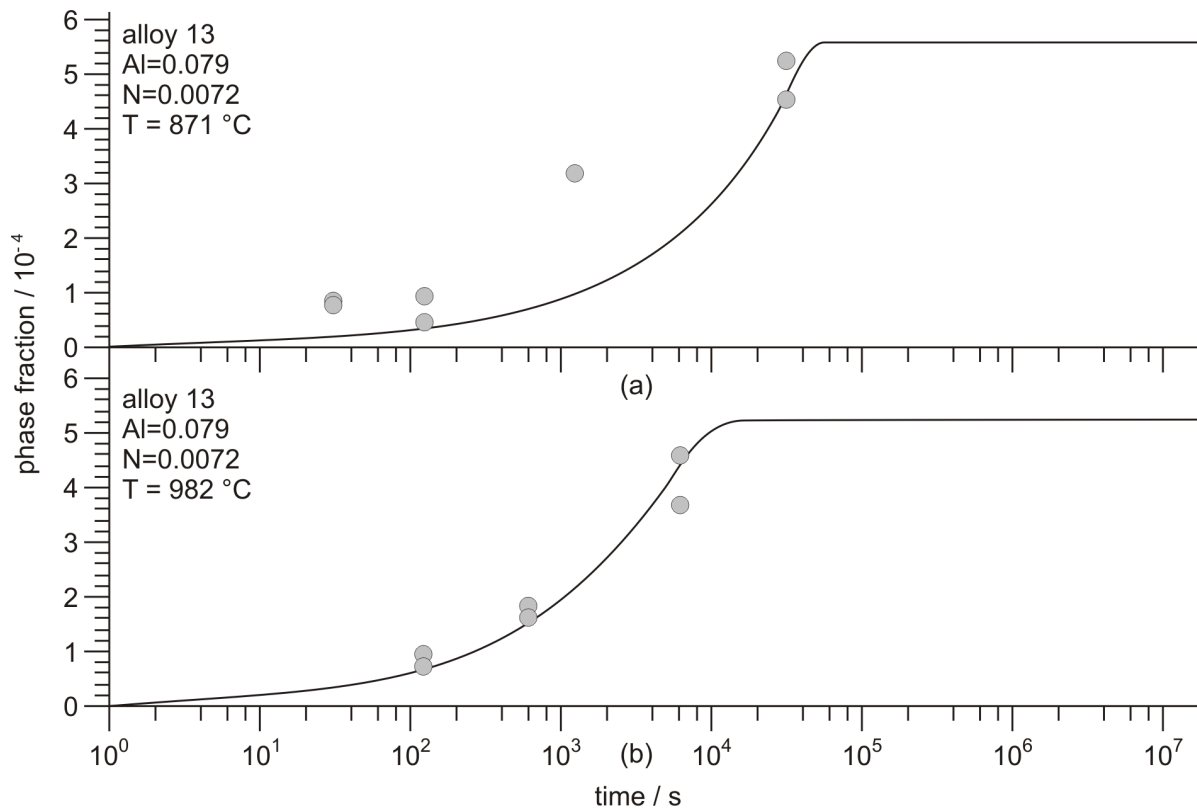


Figure 4.39: Calculated and experimental precipitation kinetics of AlN in austenite for alloy 13 at the temperatures 871 °C and 982 °C.

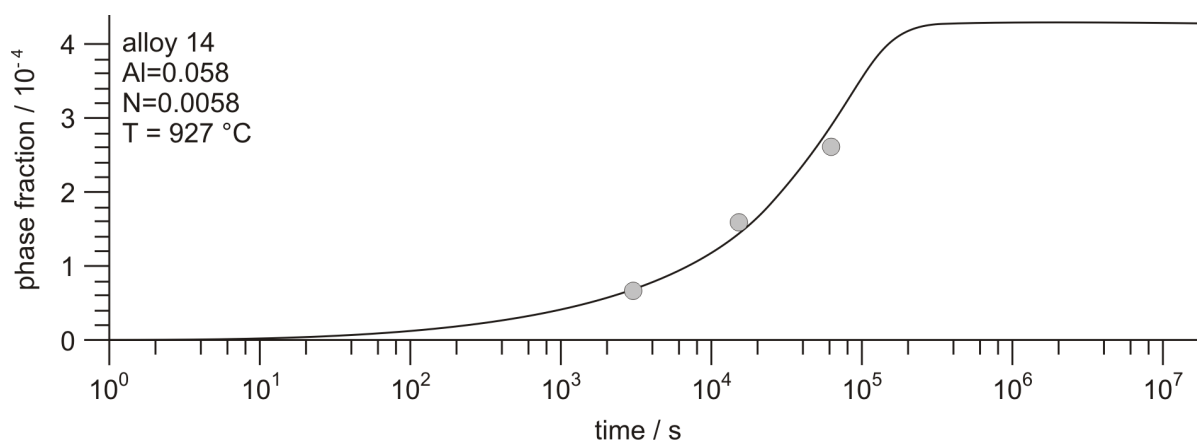


Figure 4.40: Calculated and experimental precipitation kinetics of AlN in austenite for alloy 14 at a temperature of 927 °C.

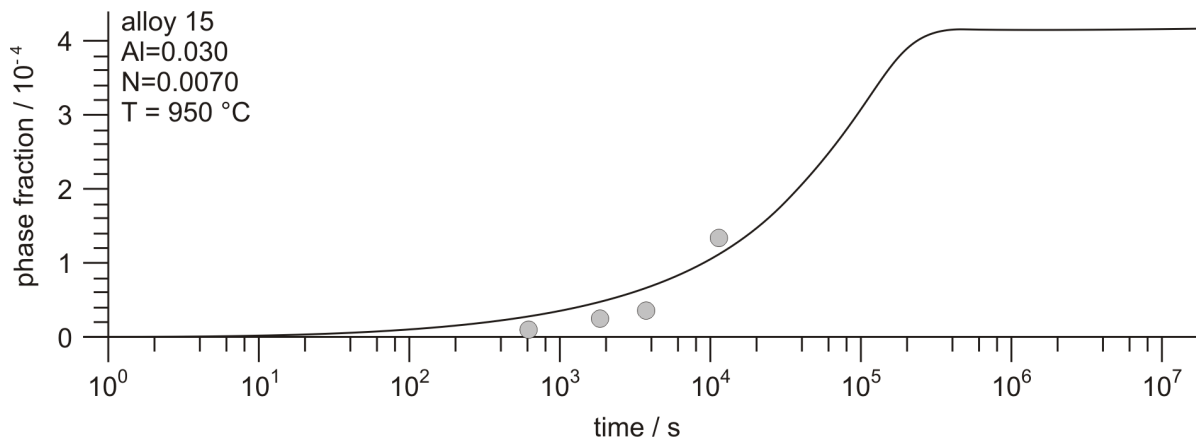


Figure 4.41: Calculated and experimental precipitation kinetics of AlN in austenite for alloy 15 at a temperature of 950 °C.

#### 4.7.5 Time-Temperature-Precipitation Plots for AlN in Steel

Figures 4.42 and 4.43 show a comparison of the precipitation kinetics of the alloys with lowest (alloy 10) and highest (alloy 11) amount of Al and N, depicted in time-temperature-precipitation (TTP) diagrams. In the simulations, the allotropic transformation from austenite to ferrite is assumed to occur at 800 °C. The calculations are performed with a grain size of 10  $\mu\text{m}$  for the ferritic matrix and 100  $\mu\text{m}$  for the austenitic matrix, respectively. Similar to figures 4.27 - 4.41, the grain boundary precipitates are represented by dotted lines, whereas dashed lines indicate the precipitates formed at dislocations. For both populations, the time for 5 % and 95 % precipitation (relative phase fraction) is displayed.

The two TTP diagrams reflect the features discussed for the previous figures. Whereas, for alloy 10, the solution temperature of AlN is calculated with 1040 °C, the solution temperature for alloy 11 is calculated with 1417 °C. As a consequence, in alloy 10, the precipitation of AlN is only observed at temperatures below 900 °C. Especially in austenite, precipitation occurs after long annealing times, e.g.  $t_{0.05} = 2 * 10^4$  s with a nose temperature of 825 °C. Moreover, above 600 °C, precipitation of AlN is observed only at grain boundaries (dotted lines). Nucleation at dislocations is only predicted for temperatures below 600 °C (dashed lines) with rather small volume fractions.

In alloy 11 (figure 4.43), the nose temperature for precipitation in austenite is calculated

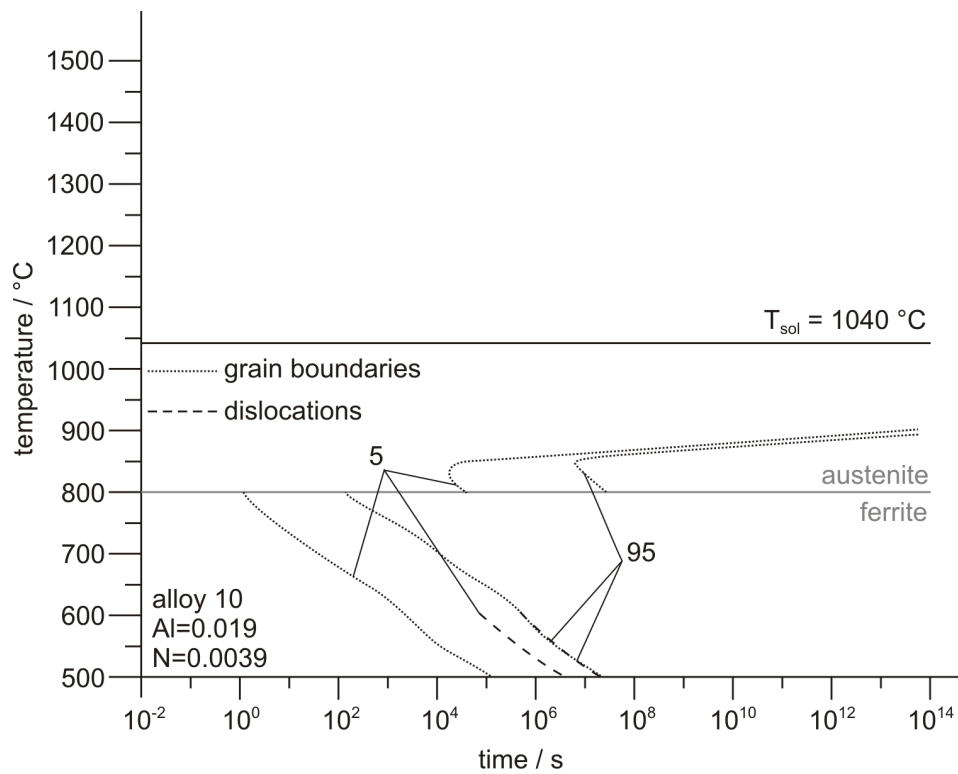


Figure 4.42: Calculated TTP diagram for alloy 10.

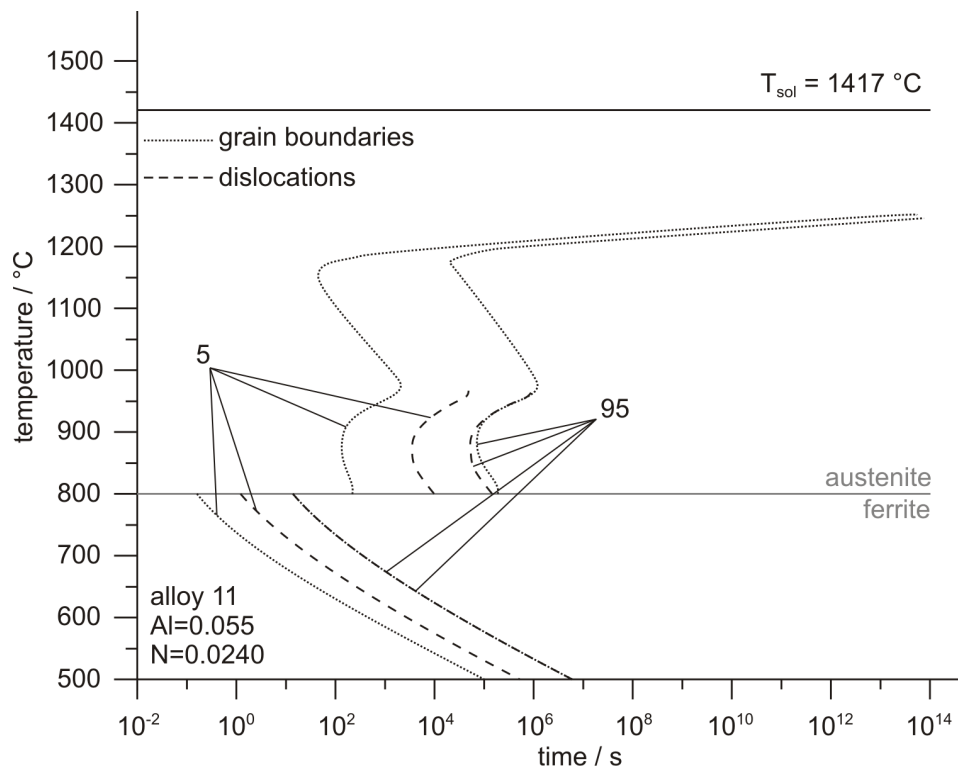


Figure 4.43: Calculated TTP diagram for alloy 11.

with 1150 °C. At this temperature, precipitation at grain boundaries occurs after 43 s. The discontinuity in the C-curve of the time-temperature-precipitation diagram for precipitation at austenite grain boundaries is attributed to the simultaneous precipitation of AlN at dislocations below a temperature of 1000 °C. Due to the change from major precipitation at grain boundaries to favored precipitation at dislocation, the total phase fraction of grain boundary precipitates is strongly reduced. Thus, the 5 % and 95 % phase fraction lines are shifted to shorter times. The fastest AlN precipitation kinetics is observed at a temperature of 800 °C for both alloys, which is the temperature of the assumed austenite to ferrite transformation. At this temperature, the precipitation of 5 % of the equilibrium phase fraction is completed after 1.5 s and 0.15 s for alloy 10 and alloy 11, respectively. Consequently, the present calculations suggest that AlN precipitation occurs almost immediately after austenite has transformed into ferrite.

#### 4.7.6 Discussion

In order to get good agreement of our simulations with the available experimental data based on a single, identical set of input parameters, several physical mechanisms and several simulation steps had to be considered. First, the thermodynamic basis for AlN formation has been reassessed and a revised Gibbs energy expression has been developed, see section 4.2.1. Comparison of this expression with solubility products reported in literature shows that the new description delivers solubility products, which are close to the 'recommended' solubility products of Wilson and Gladman [26]. Although there exists some discrepancy with the assessment of Hillert and Jonsson [94] in terms of the entropy of AlN, the revised expression is capable of reproducing the necessary driving forces for AlN precipitation in austenite and ferrite in a satisfying manner. This agreement could not be achieved with any of the alternative thermodynamic descriptions of other databases. It is not clear whether there exists any influence of the crystal structure (simple cubic or hexagonal) of the AlN nuclei on the precipitation kinetics.

The present analysis demonstrates that, for the numerical description of AlN precipitation in microalloyed steel, it is essential to account for the predominant precipitation of

AlN at grain boundaries as well as at dislocations. The numerical simulations show that precipitation at grain boundaries starts very quickly in the early stages of the precipitation process due to fast diffusion of atoms along grain boundaries. However, the kinetics becomes significantly more sluggish in the later stages because of increasingly wide diffusion distances from the grain interior to the grain boundaries, see section 4.5 [78]. This effect is clearly reflected in the shape of the volume fraction versus time curves. Especially in alloys with high Al and N contents, simultaneous precipitation of AlN at grain boundaries and dislocations is observed in the present simulations. The reason for this is founded in the assumption that the misfit strain energy for the critical nucleus at dislocations is higher compared to the values at grain boundaries. Thus, the activation barrier for dislocation nucleation is only overcome in alloys with higher Al and N contents. The preferred precipitation at dislocations at lower temperatures can be attributed to higher driving forces caused by higher undercooling and faster growth of dislocation precipitates characterized by approximately spherical diffusion fields. This issue has been discussed in section 4.5, where the model characteristics of the GBDG model and the RSDG model are compared.

## 4.8 Concurrent Precipitation of AlN and VN in Microalloyed Steel

The goal of the present section is the investigation of the simultaneous precipitation of AlN and VN in austenite [182, 183]. The presence of both types of precipitates combines the positive effects of AlN (section 3.2.1.1) and VN (section 3.2.1.2). The predominant precipitation of AlN at grain boundaries is made responsible for grain size control (section 3.1.3), whereas VN increases strength by precipitation of fine dispersed distributed particles in the grain interior (section 3.1.4). Therefore, a proper control of these precipitation reactions is necessary to achieve the desired mechanical properties. Hence, it is essential to understand the precipitation process of both, AlN and VN, as well as their kinetic interactions.

### 4.8.1 Materials

Two different alloys are used for the analysis of the concurrent precipitation of AlN and VN in the austenite phase field of microalloyed steels. The chemical compositions as well as the calculated solution temperatures for the phases AlN and VN are given in table 4.11.

Table 4.11: Chemical composition of the investigated alloys [88].

alloy	Al / wt%	V / wt%	N / wt%	$T_{sol}^{AlN} / ^\circ\text{C}$	$T_{sol}^{VN} / ^\circ\text{C}$
16	0.052	0.14	0.023	1404	1018
17	0.010	0.14	0.024	1197	1177

While both alloys contain similar amounts of V and N, the additions of Al differ significantly.

### 4.8.2 Simulation Setup

The calculation of the precipitation kinetics of AlN is carried out according to the description given in section 4.7.3. A default grain diameter (gs) of 50  $\mu\text{m}$  is used,

unless stated otherwise. In addition to the precipitation of AlN particles, a separate population of VN precipitates is introduced. The nucleation of VN is assumed to occur at dislocations, mainly, with an assumed dislocation density in austenite of  $10^{11} \text{ m}^{-2}$  [140]. The interfacial energy is evaluated from the generalized  $n$ -nearest-neighbour broken-bond approach [73, 74]. A brief description of this method is given in section 3.4.3. Figures 4.44 and 4.45 illustrate the calculated interfacial energy for a planar sharp interface as well as the corrections regarding the interfacial curvature. The solid lines represent the values for the interfacial energy of AlN, the dashed lines that of VN. The filled circles stand for the calculated interfacial energy of a planar sharp interface, whereas the squares indicate the corrections due to the spherical shape of the nucleus.

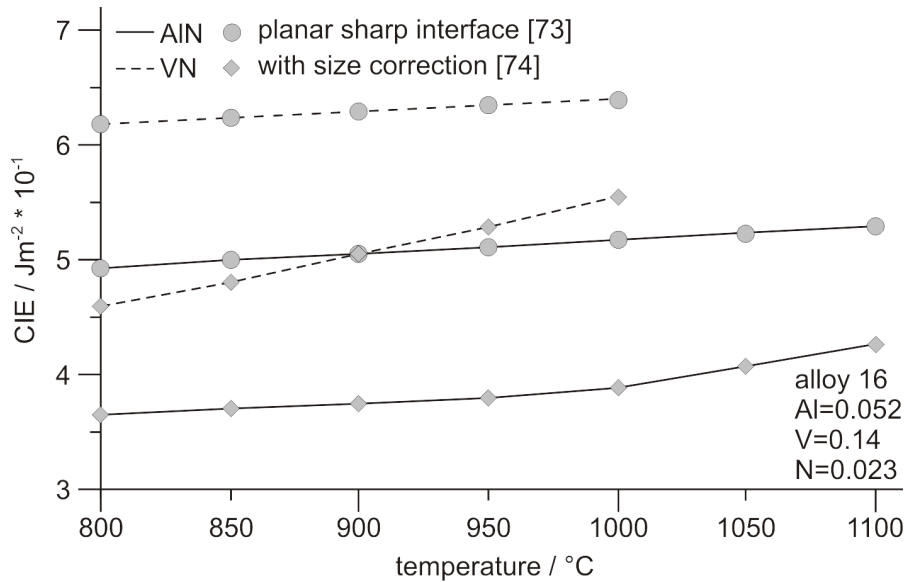


Figure 4.44: Calculated interfacial energy (CIE) for a planar sharp interface [73] and including corrections for interfacial curvature [74] for alloy 16.

In addition to the precipitate size correction, corrections accounting for entropic contributions from finite interface thickness must be considered in the calculation of the interfacial energy for the precipitate/matrix interface between VN and austenite. Particularly, close to the solution temperature, these contributions can become prominent. Therefore, an additional multiplicative temperature-dependent correction factor  $\beta$  is introduced, as described in section 3.4.3 and shown in figure 4.46. No corrections are made

for the precipitation of AlN at the austenite grain boundaries, except the corrections concerning the interfacial curvature [74].

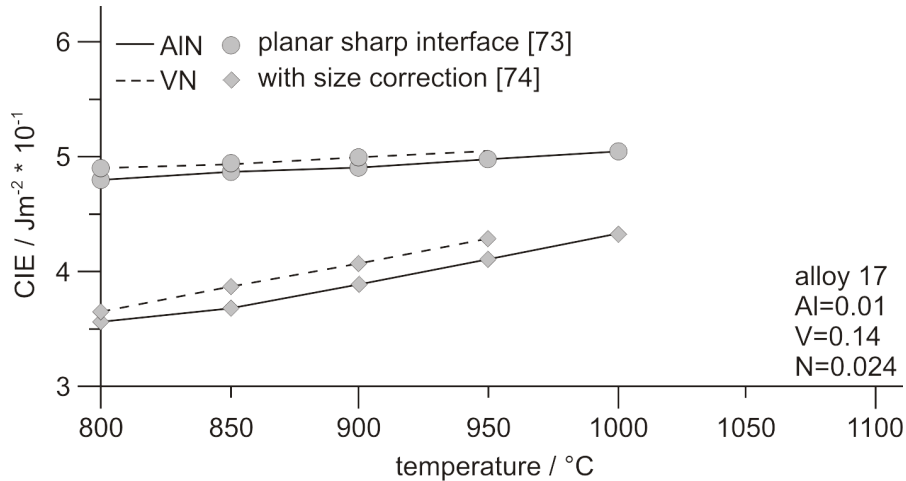


Figure 4.45: Calculated interfacial energy (CIE) for a planar sharp interface [73] and including corrections for interfacial curvature [74] for alloy 17.

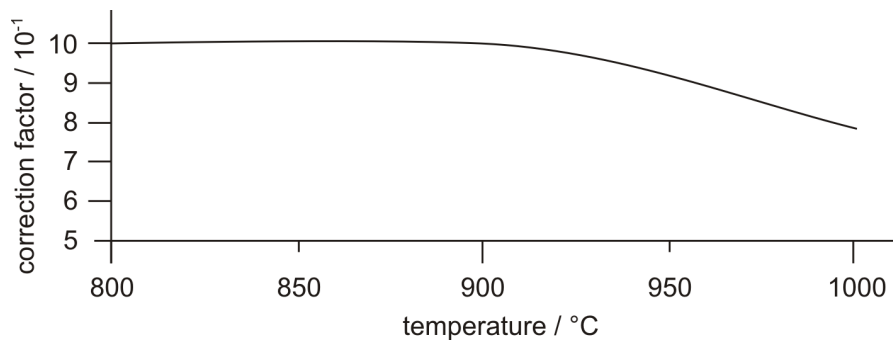


Figure 4.46: Temperature-dependent correction factor  $\beta$  to the calculated sharp interface energy of VN, in addition to the size correction  $\alpha$  for interfacial curvature (see also section 3.4.3).

### 4.8.3 Phase Fraction of AlN and VN Versus Time

The precipitation kinetics of AlN and VN is simulated for two different alloys at various different temperatures. While the importance of using the new GBDG model for the precipitation of AlN was demonstrated in sections 4.5 and 4.7, figure 4.47 shows now the



simultaneous precipitation of AlN and VN in alloy 16 using the two different approaches. Figure 4.47(a) presents the treatment with randomly distributed precipitates, whereas figure 4.47(b) shows the calculated results assuming predominant precipitation of AlN at austenite grain boundaries. The solid lines represent the phase fraction of AlN, whereas the dashed lines represent the phase fraction of VN.

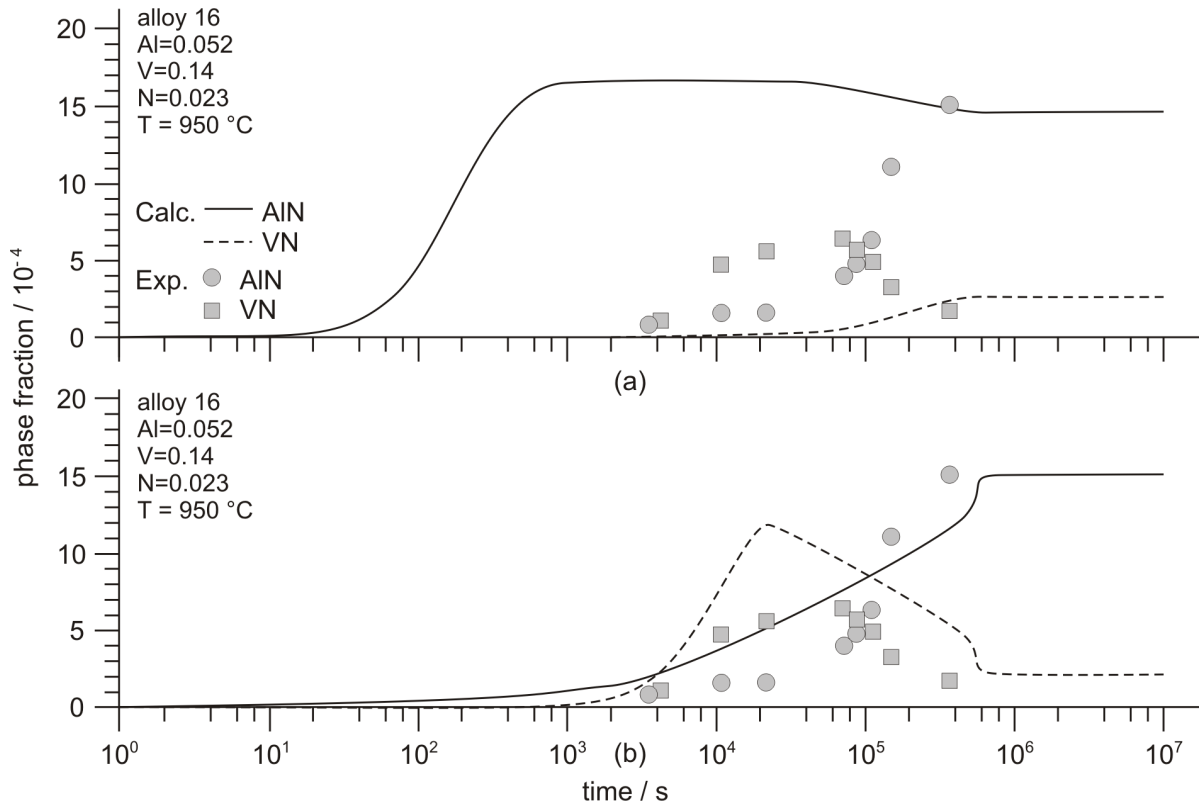


Figure 4.47: Comparison of the calculated phase fractions versus time and experimental data [88] for alloy 16 at 950 °C (a) assuming randomly distributed precipitates or (b) assuming AlN precipitation at grain boundaries.

Figures 4.48 and 4.49 show the calculated phase fraction versus time plots for both alloys and the comparison with experimental results from König et al. [88].

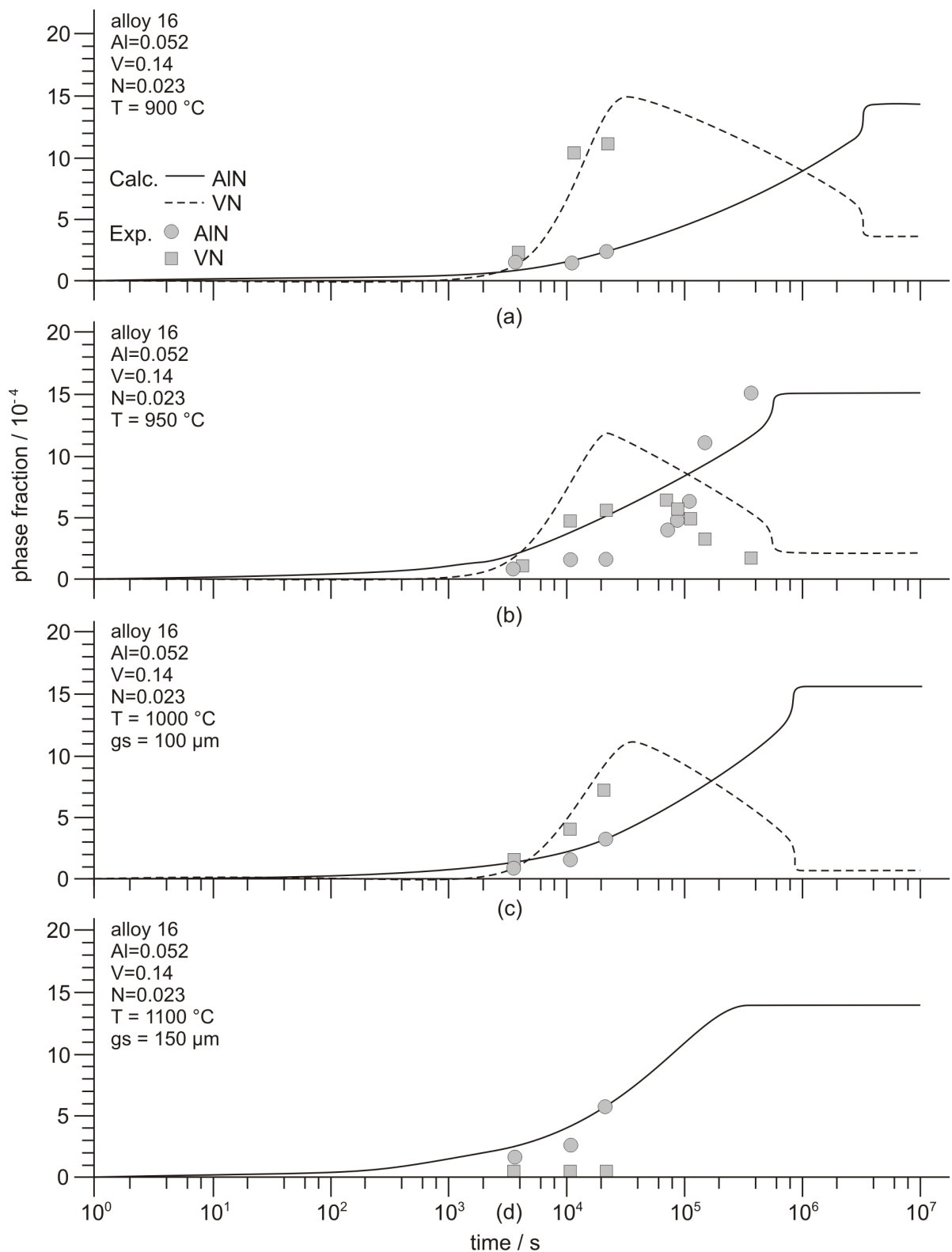


Figure 4.48: Calculated and experimental precipitation kinetics of AlN and VN in austenite for alloy 16 at the temperatures 900 °C, 950 °C, 1000 °C and 1100 °C.

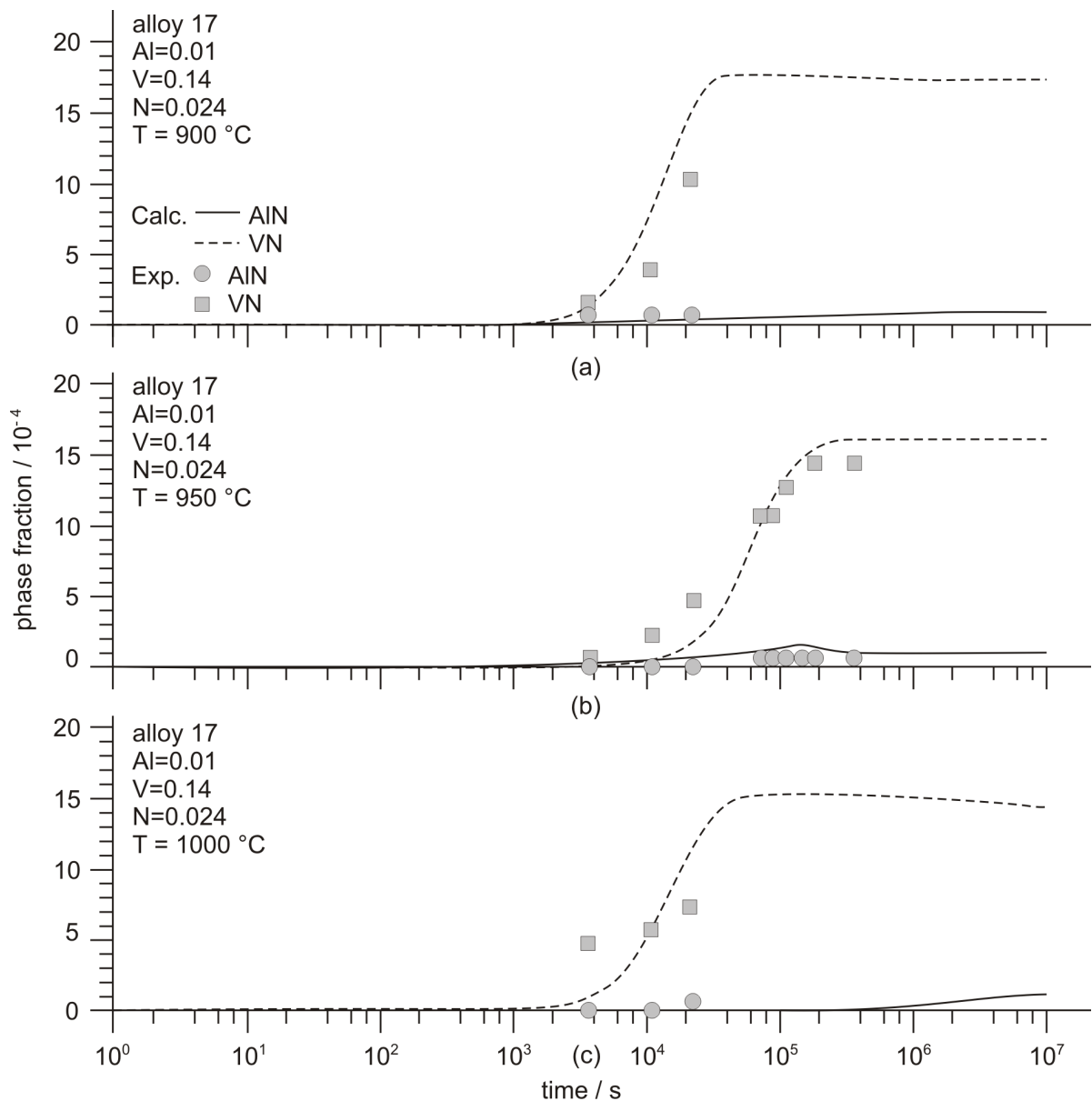


Figure 4.49: Calculated and experimental precipitation kinetics of AlN and VN in austenite for alloy 17 at the temperatures 900 °C, 950 °C and 1000 °C.

#### 4.8.4 Time-Temperature-Precipitation Plots for AlN and VN in Steel

Figures 4.50 and 4.51 shows the time-temperature-precipitation (TTP) plots for AlN and VN precipitation in both investigated alloys. The solid lines represent the precipitation of AlN, whereas the dashed lines stand for VN.

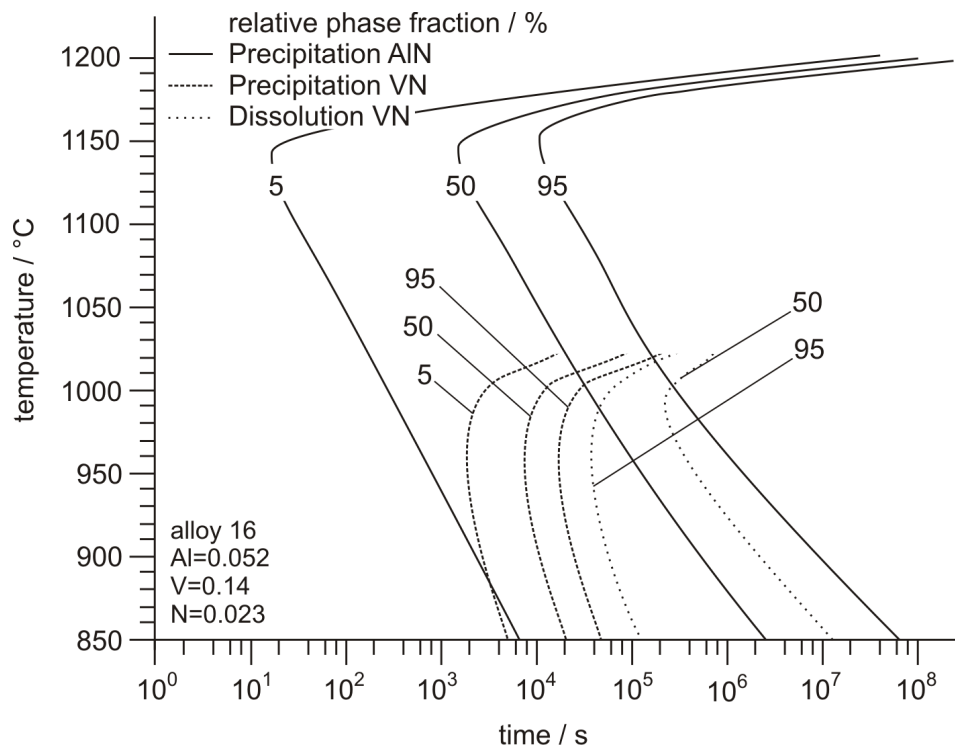


Figure 4.50: Calculated TTP diagram of AlN and VN for alloy 16.

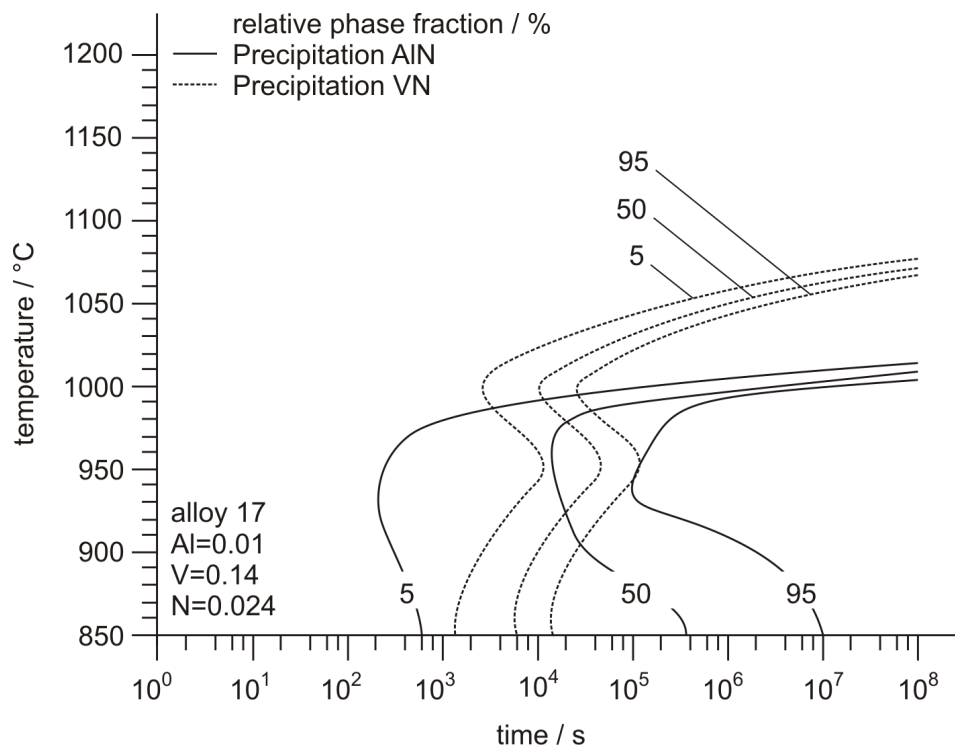


Figure 4.51: Calculated TTP diagram of AlN and VN for alloy 17.

Figure 4.51 additionally shows the predicted dissolution lines of VN, indicated by dotted lines. The lines are labelled with numbers corresponding to the relative phase fractions at each temperature.

#### 4.8.5 Discussion

The present section emphasizes the importance of considering the predominant precipitation of AlN on austenite grain boundaries, not only for the calculation of the precipitation kinetics of AlN (see also sections 4.5 and 4.7), but also for the correct prediction of the precipitation kinetics of VN. Figure 4.47 show that the classical model with randomly distributed precipitates and spherical diffusion fields strongly overestimates the precipitation kinetics of AlN. In addition, the precipitation of VN is completely suppressed if grain boundary precipitation is not taken into account correctly, see figure 4.47(a). Taking into account AlN precipitation at austenite grain boundaries leads to a good agreement between calculation and experimental observation, see figures 4.48 and 4.49. Moreover, in alloy 16, due to the decelerated precipitation of AlN at the grain boundaries, pronounced precipitation of metastable VN inside the austenite grains is observed, dissolving again, however, in favour of the thermodynamically more stable AlN. Figures 4.50 and 4.51 show the predicted TTP-plots for AlN and VN precipitation in the investigated alloys. Below 950°C, the precipitation of AlN is shifted to longer times, which is indicated by the change of the slope of the TTP curves. Although, there are only small quantities of Al present in alloy 17, the TTP diagram clearly reproduces the deceleration of the VN precipitation kinetics below approximately 1000°C. Furthermore, the TTP plots for both alloys show that the distances between the curves, which represent 5, 50 and 95 % of the relative phase fraction, are larger in the case of AlN precipitation compared to VN. This is again an indication for the experimentally observed flatter slope of the phase fraction curves for AlN precipitation at grain boundaries, see also figure 4.48.

## 4.9 Precipitation Kinetics Considering Microsegregation During Solidification

During the production process of steel, e.g. continuous casting, solidification of steel occurs within a given temperature interval. After cooling below solidification temperature, first, solid crystals nucleate, which finally represent the cores of new dendrites. During further solidification, a composition gradient from the regions where the solidification started (centre of the dendrite) to the areas where the last liquid solidified (outer shell of the dendrite) develops [184]. Consequently, in as-cast material, the precipitation behavior of nitrides and carbides is expected to be totally different at different regions of the solidification microstructure. The precipitation kinetics is very sluggish in the regions corresponding to the dendrite core, while it is many orders of magnitude faster in the outer shell of the dendrite. Pudar et al. [185, 186] studied the segregation of alloying elements as well as the precipitation behavior in the dendrite core and in the residual liquid after final solidification. However, his precipitation kinetics calculations concentrated on the prediction of maximum precipitate sizes in the dendrite core and in between the dendrites. The authors did not take into account the phase fraction of precipitates in the region between the dendrite center and the region of final solidification separately. This first approach is extended in this section to be applicable for the prediction of phase fractions in segregated microalloyed steels, too.

The new approach, which is presented subsequently, deals with the prediction of phase fractions, particle sizes and number densities locally at every point of the primary solidification microstructure as well as the accumulated phase fractions, particle sizes and number densities globally in the whole microstructure.

### 4.9.1 Simulation Procedure

The numerical procedure of the calculation of the precipitation process during solidification is subdivided into two calculations [187]. First, the calculation of the segregation of alloying elements during solidification, and, second, the calculation of the precipi-

tation kinetics of the nitrides and carbides in the regions of different chemistry. The calculations are performed using the thermodynamic database 'mc\_steel' [83] (version 1.20). The numerically computed results are compared to independent experimental data [188]. The chemical compositions of the investigated alloy is given in table 4.12.

Table 4.12: Chemical composition of the investigated alloy [188].

alloy	C	Si	Mn	Al	V	Ti	Nb	N
18	0.175	0.014	1.46	0.032	0.006	0.0151	0.013	0.004

#### 4.9.1.1 Segregation of Alloying Elements During Solidification

For the calculation of the segregation of alloying elements during solidification, 'Scheil-Gulliver' simulations are performed [189,190]. In the simulations, the phases LIQUID, BCC, FCC, AlN, NbN, NbC, TiN, TiC, VN and VC are considered. These calculations deliver the fraction and composition of all phases during solidification step by step from the liquidus temperature to the temperature where final solidification of the residual liquid phase occurs.

On the basis of these results (see section 4.9.2.2), the composition profile of the dendrites is divided into eight different zones, one core, six shells and 5% residual liquid, which represent the composition gradient from the dendrite center to the region of final solidification. Figure 4.52 shows a schematic image of such dendrites. The darker the zone, the more alloying elements are in solution and available for precipitation of nitrides and carbides. The content of alloying elements is drawn as a function of the phase fraction of the solid phase in figure 4.53 for the case of Ti. With ongoing solidification, the content of Ti increases in the solidifying shells as well as in the residual liquid phase. The chemical compositions of the different zones are weighted according to figure 4.53. Thus, the core and the first shell are weighted with 25%, the following four shells with 10% and the last shell and the residual liquid are weighted with 5%, respectively.

To determine the mean composition of each zone, the mean content of each alloying element is calculated and multiplied with the weighting coefficient. Summarizing these

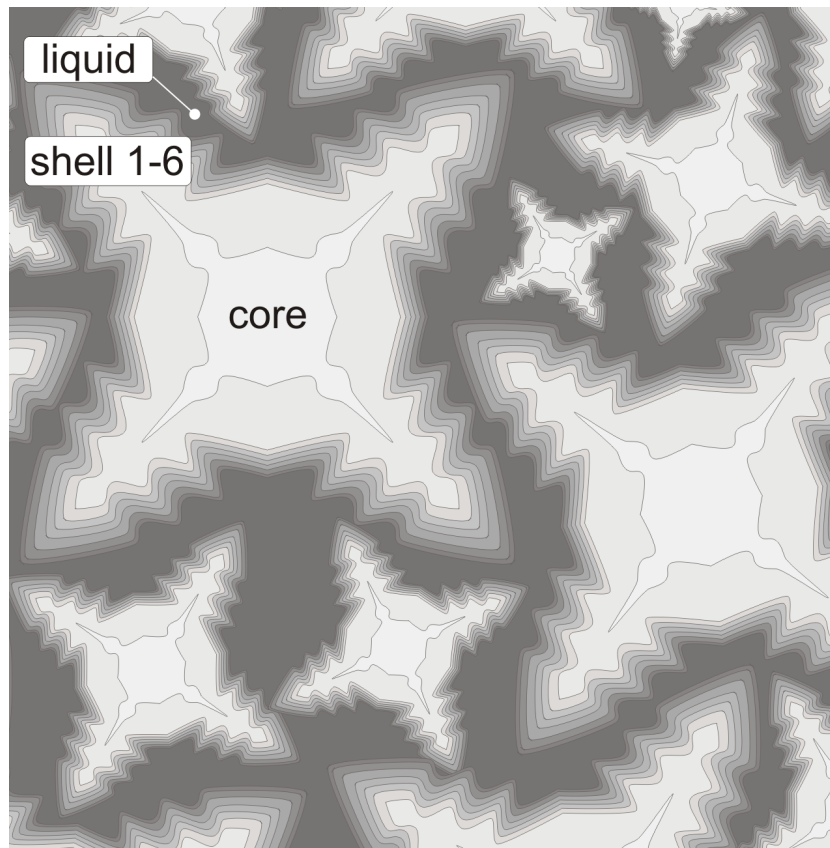


Figure 4.52: Dendrites with zones of different chemical compositions.

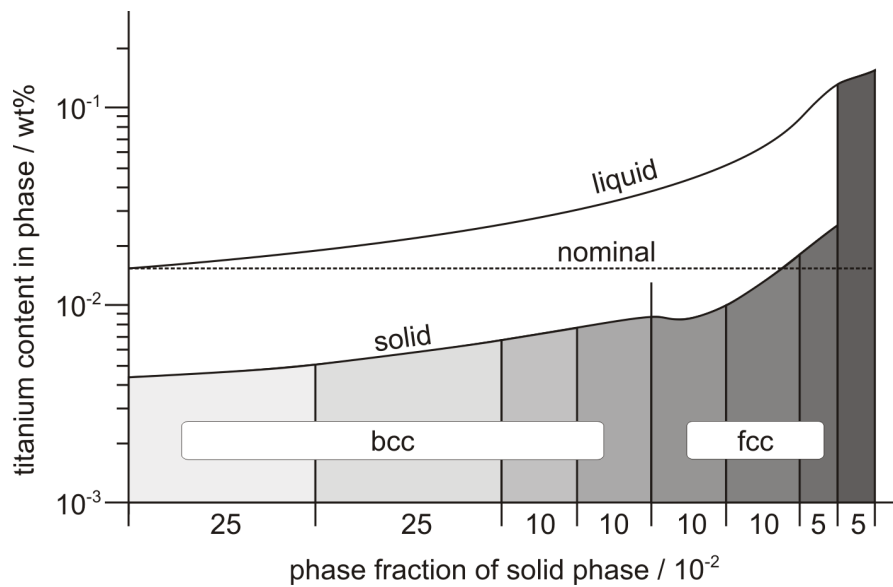


Figure 4.53: Ti content in different zones as function of the phase fraction of the solid phase.



weighted chemical compositions of all zones deliverers the nominal composition of the steel.

#### 4.9.1.2 Precipitation Kinetics in Regions of Different Chemistry

The heat treatment used in the present simulations is shown in figure 4.54. After a linear cooling from 1300°C to 900°C the material is isothermally annealed for 4 h. Subsequently, a series of calculations with different chemical compositions, according to the results of the 'Scheil-Gulliver' simulations (see section 4.9.2.2) are performed. Thus, the precipitation kinetics for every zone of the dendrite is determined.

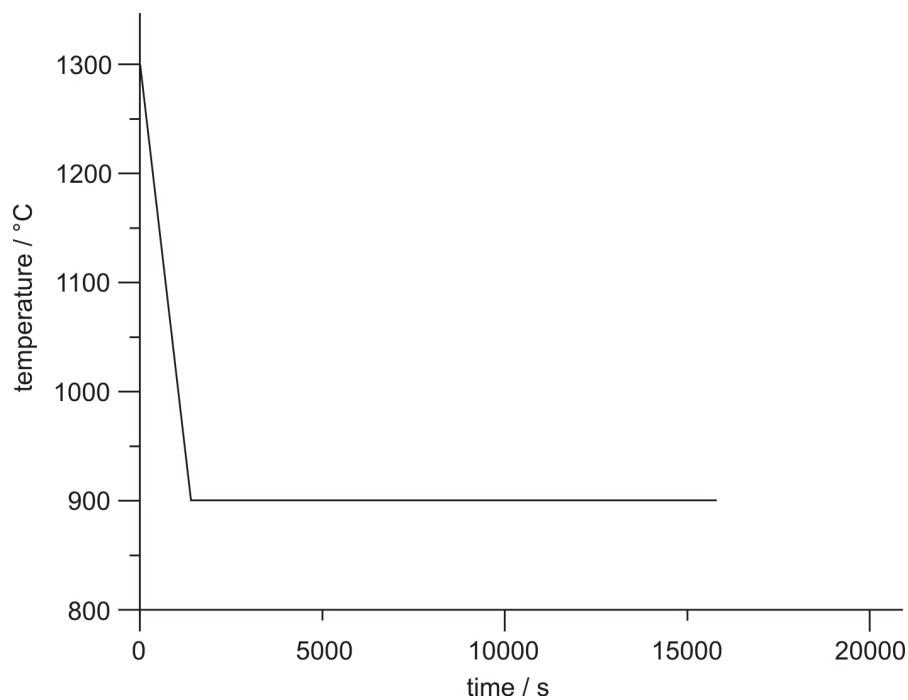


Figure 4.54: Heat treatment used for the precipitation kinetics simulations.

Finally, the individual phase fractions of all zones are summarised and compared to the experimental investigations. This is done after accumulating the amount of the carbides and nitrides of one microalloying element because the amount of Ti and Nb delivered in the experimental studies is also an accumulated quantity and gives no information whether the alloying element is bonded to a nitride or carbide. The comparison of the calculated volume fractions of TiN, TiC and NbC with the measured weight fractions

of Ti and Nb is done according to the procedure described in section 4.7.2.

## 4.9.2 Results

After some equilibrium considerations of alloy 18 the segregation of alloying elements is calculated. With the results of these prior segregation calculations, precipitation kinetics calculations are performed in different chemical regions of the dendrites.

### 4.9.2.1 Equilibrium Calculations

For a rough estimation about stable nitrides and carbides in alloy 18, an equilibrium calculation is performed using the nominal composition, given in table 4.12.

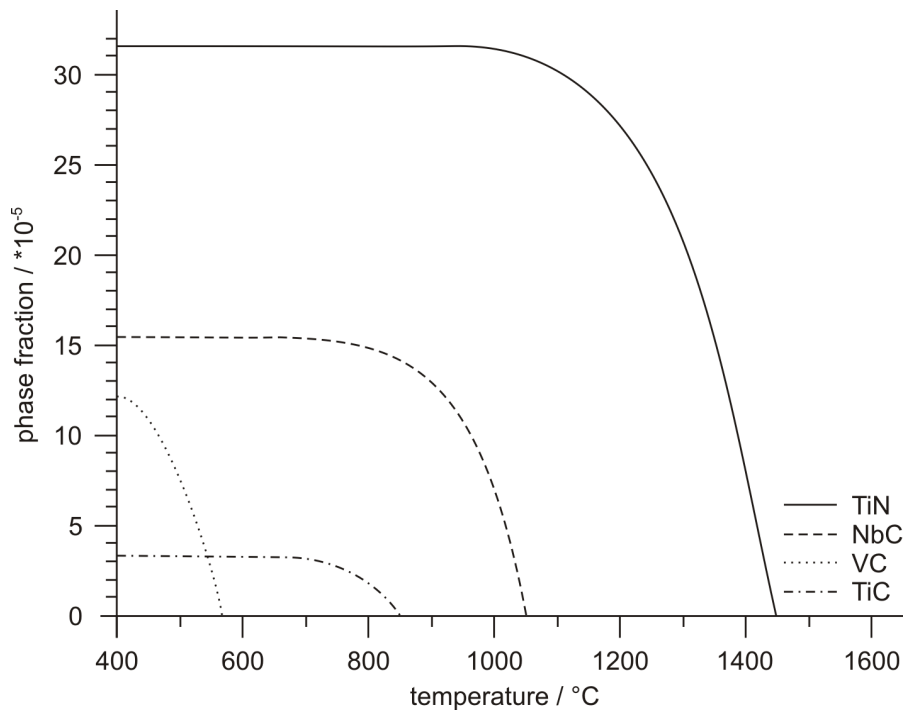


Figure 4.55: Equilibrium phase fractions versus temperatures of alloy 18.

Several nitrides (AlN, NbN, TiN, VN) and carbides (NbC, TiC, VC) are considered in the simulation. Figure 4.55 shows that, in equilibrium, the second phase precipitates TiN, NbC, TiC and VC are stable below solution temperatures of 1446°C, 1052°C, 850°C and 568°C, respectively. Thus, after isothermal annealing of alloy 18 for infinitely long time at a temperature of 900°C, only the precipitation of TiN and NbC should occur.

#### 4.9.2.2 Segregation of Alloying Elements During Solidification

After casting, the samples are quenched and reheated to 1300°C, afterwards directly cooled from 1300°C to 900°C and isothermally annealed at this temperature for 4h [188].

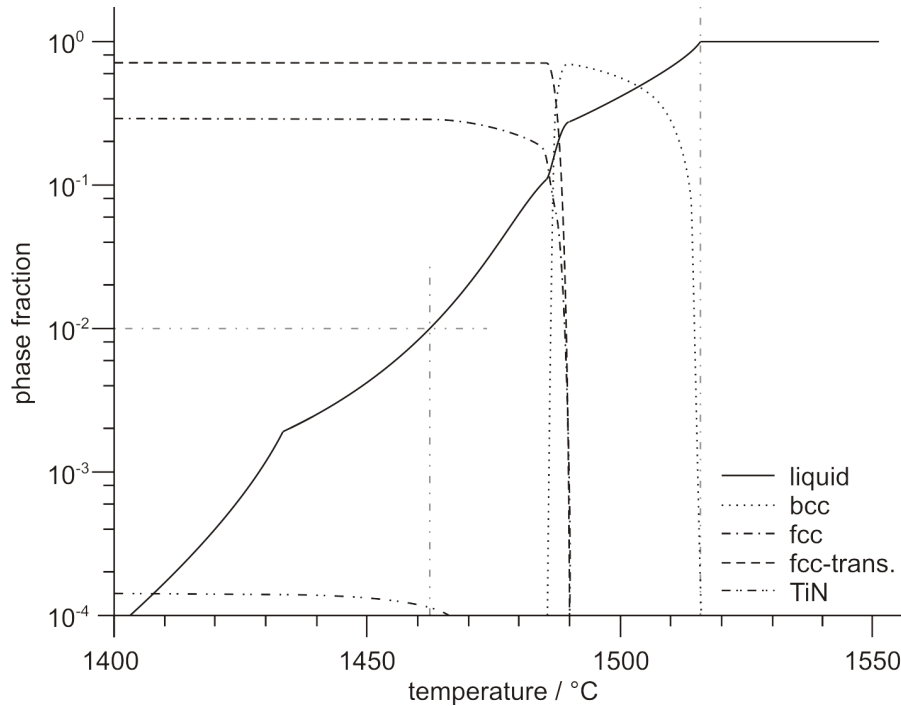


Figure 4.56: Scheil-Gulliver simulation of alloy 18.

Since we expect strong segregation during solidification and no time for diffusion to level out the composition gradients, the precipitation kinetic calculations for the as-cast material cannot be carried out with the nominal compositions of the steels. Instead, 'Scheil-Gulliver' simulations are performed (4.9.1.1) to estimate the segregation of the alloying elements during solidification. Figure 4.56 shows the result of the 'Scheil-Gulliver' calculation for alloy 18. The solidification range is between 1517°C and 1463°C, if solidification of the residual liquid is assumed to occur at 1%. If this is the case, some coarse primary TiN precipitates are likely to be present in this steel.

Table 4.13 shows the chemical composition of each zone in the dendrite obtained from the 'Scheil-Gulliver' simulation. The mean values of alloying elements at different residual liquid states are calculated in the different matrix phases.

Table 4.13: Calculated amount of alloying elements (wt-ppm) in the different zones (dendrite core, shells, residual liquid) of the primary solidification microstructure.

shell no.	residual liquid	measured in phase	amount of alloying elements in shell / wt-ppm							
			Mn	C	Nb	Ti	Si	Al	V	N
7	0.05	liquid	31400	5321	1592	1262	317	336	174	67
6	0.076	fcc	20940	1512	160	211	205	291	82	39
5	0.177	fcc	17246	1442	80	111	162	264	59	43
4	0.256	fcc	15776	1404	59	83	146	252	51	46
3	0.354	bcc	13810	621	51	79	133	325	54	21
2	0.457	bcc	12938	525	42	68	123	331	52	19
1	0.612	bcc	12020	426	34	57	114	336	49	15
0	0.867	bcc	11000	326	26	46	103	342	46	12

Table 4.14: Corrected and weighted amount of alloying elements (wt-ppm) in the different zones (dendrite core, shells, residual liquid) of the primary solidification microstructure.

shell no.	residual liquid	weight	amount of alloying elements in shell / wt-ppm							
			Mn	C	Nb	Ti	Si	Al	V	N
7	0.05	0.05	1598	502	82	71	16	17	9	5
6	0.076	0.05	1066	143	8	12	11	15	4	3
5	0.177	0.1	1756	272	8	13	17	27	6	7
4	0.256	0.1	1606	265	6	9	15	25	5	7
3	0.354	0.1	1406	117	5	9	14	33	6	3
2	0.457	0.1	1317	99	4	8	13	33	5	3
1	0.612	0.25	3059	201	9	16	29	85	13	6
0	0.867	0.25	2800	153	7	13	27	86	12	5
sum			14607	1751	130	150	140	321	60	40
correction factor			1.018	1.885	1.035	1.132	1.03	1.008	1.04	1.62

Using mean values of the alloying elements leads to some discrepancies between the sum of alloying additions of all zones and the nominal composition. Therefore, correction factors are introduced, which are multiplied with the amount of each element in every area, see table 4.14. Generally, these factors are very small, except for the interstitial elements C and N. These achieve equilibrium (liquid and solid phases) due to rapid back-diffusion from the liquid phase into the solid phase [189]. In these cases, the correction factors are approximately in the order between 1,5 and 2. However, using these factors and weighting the amount of the elements with the width of the zones (dendrite core, shells and residual liquid), ensures mass conservation with respect to the nominal composition (see table 4.14).

#### 4.9.2.3 Precipitation Kinetics in Regions of Different Chemistry

Several precipitation kinetics calculations are performed using the chemical composition of each zone presented in section 4.9.2.2 as input parameter. Figure 4.57 displays the summarized results of the calculations. The dashed lines indicate the individual phase fractions of TiN in the different zones (dendrite core, shells, residual liquid) of the dendrite. Whereas the precipitation of TiN is predicted in each zone of the dendrites, precipitation of TiC and NbC only occurs in the interdendritic regions. Finally, the sum of the individual phase fractions of TiN in each zone, as well as the sum of the phase fractions of TiN and TiC is depicted. The accumulation of the phase fractions of TiN and TiC is necessary, due to the fact that the amount of Ti and Nb provided from the experimental studies is given in total quantities and give no information whether the alloying element is bonded to a nitride or carbide. These total quantities are indicated by solid circles and squares for Ti and Nb, respectively.

#### 4.9.3 Discussion

The precipitation kinetics calculations of alloy 18 performed in this study show that the proposed method leads to fair agreement between calculations and experimental observations. This could not be achieved with any calculations using the nominal com-

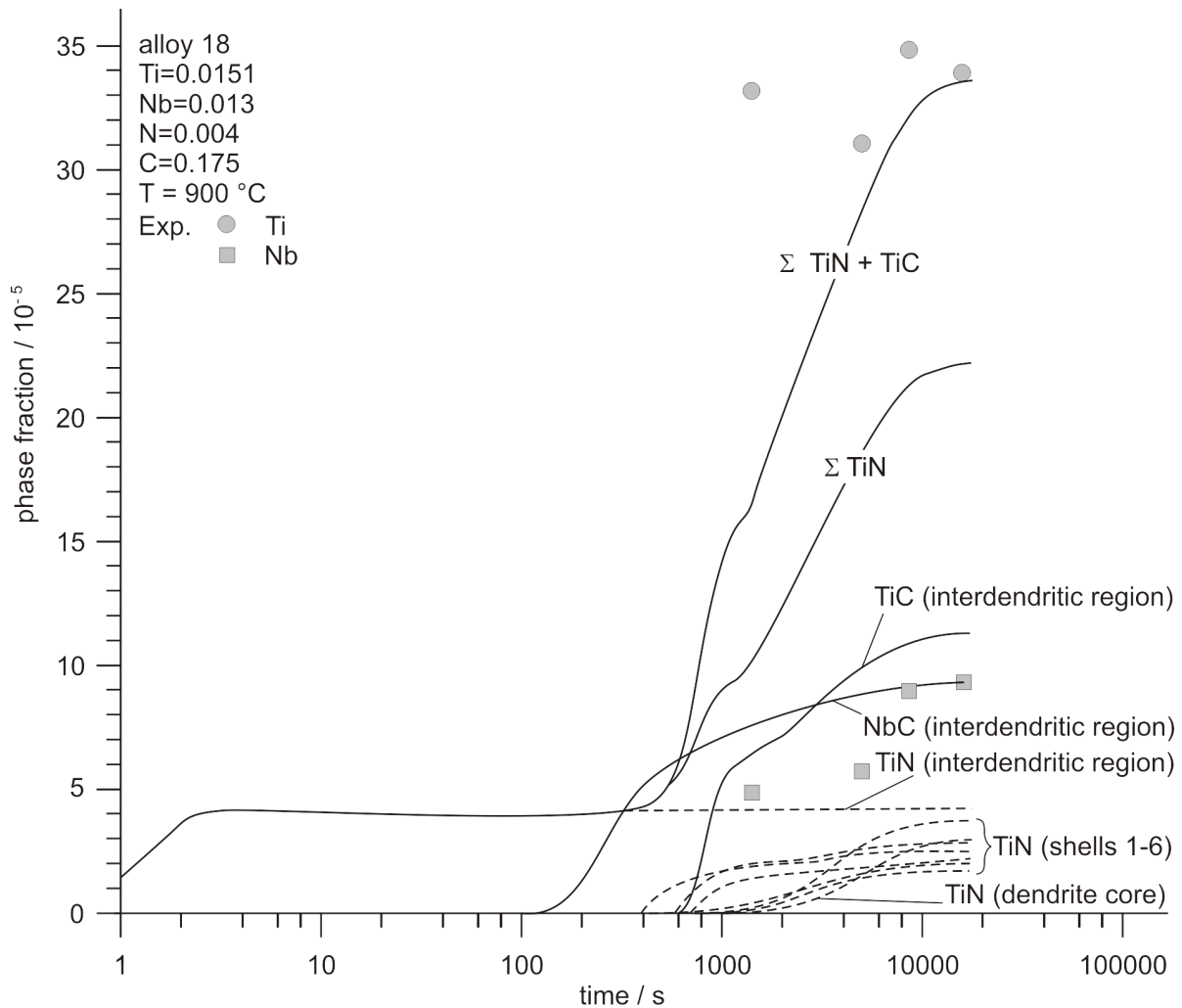


Figure 4.57: Calculated phase fractions of nitrides and carbides in alloy 18 at a temperature of 900°C.

position of alloy 18. Moreover, figure 4.55 shows that no TiC is stable in equilibrium when using the nominal composition. However, the measured amounts of Ti can only be reached, when considering both, TiN and TiC precipitated in different zones of the segregated dendrites, see figure 4.57. While the precipitation of TiN is observed in all regions of the dendrite (core, shells and residual liquid), precipitation of TiC and NbC is only observed in the interdendritic region.

It should be emphasized that the experimental results for Ti are compared to the sum of the results of the numerical calculations of two Ti containing phases. Thus, to match the experimental results, it is necessary that all the critical parameters, such as thermody-

namics, interfacial energies or diffusion constants, of the two phases must be reasonably accurate. Furthermore, it should be noted, that no adaptations of any of these parameters have been made. In this sense, the calculations can be considered as predictions and must be evaluated as good.

The calculations based on the proposed method also provide some information about the spatial arrangement of the particles. Particle sizes, number densities and phase fractions can be determined as a function of the distance from the core of a dendrite.

# Chapter 5

## Summary

The present thesis deals with the numerical simulation of the precipitation kinetics of various nitrides and carbides in (HSLA) microalloyed steel. The solubility data of all microalloy carbides and nitrides are reviewed and compared to different thermodynamic assessments. Reasonable descriptions of the thermodynamic properties of the nitrides and carbides are incorporated in the thermodynamic database 'mc\_steel' in the form of simple phase descriptions.

Employing the software package MatCalc and the corresponding thermodynamic and diffusion databases, numerical precipitation kinetics simulations are performed using two different models. On the one hand, a new model for grain boundary precipitation is employed, which takes into account fast short-circuit diffusion along grain boundaries as well as slower bulk diffusion inside the grain. On the other hand, the classical treatment for randomly distributed precipitates with spherical diffusion fields is used. In the case of randomly distributed precipitates located at dislocations, volumetric misfit strain energies between precipitate/matrix interfaces are fully accounted for in the present calculations.

Detailed parameter studies are presented, describing the characteristics of both models. It is shown that the grain boundary model reproduces the well-known effect of grain size on the grain boundary precipitation kinetics, where large grain sizes significantly retard the overall precipitation process due to large inner grain diffusion distances. Moreover, depending on the ratio of grain boundary diffusion to bulk diffusion, simultaneous



growth and coarsening of precipitates can occur. In the case of high grain boundary diffusivity, the initial stage of precipitation is accelerated by fast nucleation and growth kinetics. Simultaneously, the phase fraction increase is significantly slower compared to the random distribution case. In the case of the classical treatment of randomly distributed precipitates, the parameter studies demonstrate the high sensitivity of the calculations on the input quantities volumetric misfit and Young's modulus.

The precipitation kinetics of AlN and VN is analyzed by the numerical simulation of isothermal annealing of microalloyed steels covering various chemical compositions. The simulations clearly demonstrate that AlN formation occurs by simultaneous precipitation at grain boundaries and at dislocations. Depending on chemical composition, grain size and annealing temperature, predominant precipitation at grain boundaries and/or at dislocations is predicted. With the application of the two models representing these two mechanisms, an excellent agreement between prediction and experimental data is achieved. Moreover, it is demonstrated, that for certain ratios between Al and V, after a temporary precipitation of VN, these precipitates dissolve again, in favour of the thermodynamically more stable AlN. It is also shown that, for a consistent description of these precipitation reactions, it is necessary to account for several physical mechanisms that are often neglected in this type of simulations, among them the precipitate/matrix volumetric misfit, the temperature dependent Young's modulus, as well as the ratio between bulk and grain boundary diffusion.

Finally, the precipitation kinetics of TiN and NbC is simulated, taking into account the microsegregation of these elements during solidification. A new approach is presented, which allows calculating the precipitation parameters locally at every point of the dendrite. The chemical concentration profile from the dendrite core to the segregated regions is simulated with 'Scheil-Gulliver' calculations of the microsegregation process during solidification. Using these chemical compositions, the precipitation parameters are calculated in each individual zone. Accumulation of phase fractions, particle sizes and number densities leads to the overall precipitation products in the whole material. The calculations and comparison with experimental data clearly demonstrate that these segregation effects must be taken into account.

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# Appendix

## A MatCalc Script for the Calculation of the Precipitation Kinetics of Nitrides and Carbides in Steel

Employing the software package MatCalc, precipitation kinetics calculations can be performed by executing script files, which also ensure the documentation of the used parameters. In the following a script file is presented for the calculation of the precipitation of nitrides and carbides during the production process of a hypothetical microalloyed steel, according to a hypothetical heat treatment shown in figure A.1.

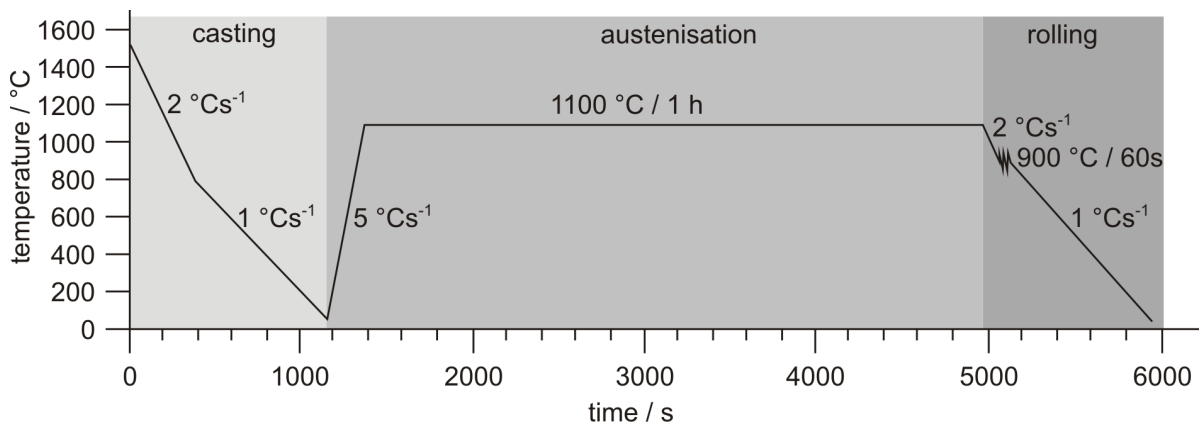


Figure A.1: Heat treatment used for the simulation of the production process of a hypothetical steel during the stages casting, austenisation and rolling.

The simulation of AlN and VN is performed according to the descriptions given in sections 4.7 and 4.8, whereas the calculation of the precipitation kinetics of other precipitates must be considered as a first estimation and is not verified yet.

```

$ *****
$ ***** GENERAL INFORMATION *****
$ *****
$ Simulation of the production process of a hypothetical microalloyed steel
$ Database: mc_steel_1.86.tdb, mc_sample_fe.ddb
$ Author: R. Radis
$ Creation date: 2010-05-27
$ This is a script for MatCalc version 5.30
$ *****
$ ***** SETUP INFORMATION *****
$ *****

new-workspace                               $open new workspace

set-workspace-info +Selected elements:       $give some information about the script
set-workspace-info +C, Si, Mn, Al, V, Nb, N, Cr, Mo, Ni
set-workspace-info +
set-workspace-info +Matrix phases: Ferrite (BCC_A2), Austenite (FCC_A1)
set-workspace-info +Precipitate phases: AlN, VN, VC, NbN, NbC, M23C6, M7C3 and M3C (Cementite)
set-workspace-info +
set-workspace-info +-----
set-workspace-info +More information about MatCalc on
set-workspace-info + http://www.matcalc.at
set-workspace-info +-----

use-module core                             $select core module for kinetic simulation

$ *****
$ ***** SYSTEM SETUP *****
$ *****

$ *****
$ ***** DATABASES, CHEMICAL COMPOSITION, SELECTED PHASES *****
$ *****

open-thermo-database mc_steel_1.86.tdb      $open thermodynamic database
select-elements Al C Cr Fe Mn Mo N Nb Si V  $select elements
select-phase bcc_a2 fcc_a1 m23c6 m7c3 cementite aln vn vc nbn nbc Cr2N $select phases
set-reference-element fe                    $set fe as reference element
read-thermodyn-database                     $read thermodynamic database
enter-composition wp al=0,02 c=0,1 cr=1,5 mn=1 n=0,01 nb=0,03 si=0,3 v=0,2 $system setup
read-mobility-database mc_sample_fe.ddb     $read diffusion data

$ *****
$ ***** HEAT TREATMENTS, VARIABLES, TABLES & FUNCTIONS *****
$ *****

$ ----- HEAT TREATMENTS -----

create-heat-treatment ht                    $create a new heat treatment
append-ht-segment ht                       $append new segment to heat treatment
edit-ht-segment ht . d n austenite         $set precipitation domain to austenite
edit-ht-segment ht . s 1550                $set start temperature of heat treatment
edit-ht-segment ht . l 800 2              $cool to 800°C with a cooling rate of 2°C/s

edit-ht-segment ht . r + set-precipitation-parameter aln_p0 n s g $nucleation is defined to occur in austenite
edit-ht-segment ht . r + set-precipitation-parameter aln_p1 n s d $nucleation of aln at gb in austenite
edit-ht-segment ht . r + set-precipitation-parameter aln#01_p0 n s n $no nucleation of aln at gb in ferrite
edit-ht-segment ht . r + set-precipitation-parameter aln#01_p1 n s n $no nucleation of aln at d in ferrite

```



edit-ht-segment ht . r + set-precipitation-parameter vn_p0 n s d	\$nucleation of vn at d in austenite
edit-ht-segment ht . r + set-precipitation-parameter vn#01_p0 n s n	\$no nucleation of vn in ferrite
edit-ht-segment ht . r + set-precipitation-parameter vc_p0 n s d	\$nucleation of vc at d in austenite
edit-ht-segment ht . r + set-precipitation-parameter vc#01_p0 n s n	\$no nucleation of vc in ferrite
edit-ht-segment ht . r + set-precipitation-parameter nbn_p0 n s d	\$nucleation of nbn at d in austenite
edit-ht-segment ht . r + set-precipitation-parameter nbn#01_p0 n s n	\$no nucleation of nbn in ferrite
edit-ht-segment ht . r + set-precipitation-parameter nbc_p0 n s d	\$nucleation of nbc at d in austenite
edit-ht-segment ht . r + set-precipitation-parameter nbc#01_p0 n s n	\$no nucleation of nbc in ferrite
edit-ht-segment ht . r + set-precipitation-parameter m7c3_p0 n s n	\$no nucleation of m7c3 in ferrite
edit-ht-segment ht . r + set-precipitation-parameter m23c6_p0 n s n	\$no nucleation of m23c6 in ferrite
edit-ht-segment ht . r + set-precipitation-parameter Cr2N_p0 n s n	\$no nucleation of cr2n in ferrite
edit-ht-segment ht . r + set-precipitation-parameter cementite_p0 n s n	\$no nucleation of cementite in ferrite
append-ht-segment ht	\$append new segment to heat treatment
edit-ht-segment ht . d n ferrite	\$change precipitation domain
edit-ht-segment ht . 1 25 1	\$cool to 25°C with a cooling rate of 1°C/s
edit-ht-segment ht . r + set-precipitation-parameter aln_p0 n s n	\$change nucleation to occur in ferrite
edit-ht-segment ht . r + set-precipitation-parameter aln_p1 n s n	\$no nucleation of AIN at gb in austenite
edit-ht-segment ht . r + set-precipitation-parameter aln#01_p0 n s g	\$no nucleation of AIN at d in austenite
edit-ht-segment ht . r + set-precipitation-parameter aln#01_p1 n s d	\$nucleation of AIN at d in ferrite
edit-ht-segment ht . r + set-precipitation-parameter vn_p0 n s n	\$no nucleation of VN in austenite
edit-ht-segment ht . r + set-precipitation-parameter vn#01_p0 n s d	\$nucleation of VN at d in ferrite
edit-ht-segment ht . r + set-precipitation-parameter vc_p0 n s n	\$no nucleation of VC in austenite
edit-ht-segment ht . r + set-precipitation-parameter vc#01_p0 n s d	\$nucleation of VC at d in ferrite
edit-ht-segment ht . r + set-precipitation-parameter nbn_p0 n s n	\$no nucleation of NbN in austenite
edit-ht-segment ht . r + set-precipitation-parameter nbn#01_p0 n s d	\$nucleation of NbN at d in ferrite
edit-ht-segment ht . r + set-precipitation-parameter nbc_p0 n s n	\$no nucleation of NbC in austenite
edit-ht-segment ht . r + set-precipitation-parameter nbc#01_p0 n s d	\$nucleation of NbC at d in ferrite
edit-ht-segment ht . r + set-precipitation-parameter m7c3_p0 n s d	\$nucleation of M7C3 at d in ferrite
edit-ht-segment ht . r + set-precipitation-parameter m23c6_p0 n s d	\$nucleation of M23C6 at d in ferrite
edit-ht-segment ht . r + set-precipitation-parameter Cr2N_p0 n s d	\$nucleation of Cr2N at d in ferrite
edit-ht-segment ht . r + set-precipitation-parameter cementite_p0 n s d	\$nucleation of cementite at d in ferrite
append-ht-segment ht	\$append new segment to heat treatment
edit-ht-segment ht . d y	\$use inherent precipitation domain
edit-ht-segment ht . 1 800 5	\$heat to 800°C with a heating rate of 5°C/s
append-ht-segment ht	\$append new segment to heat treatment
edit-ht-segment ht . d n austenite	\$change precipitation domain
edit-ht-segment ht . 1 1100 5	\$heat to 1100°C with a heating rate of 5°C/s
edit-ht-segment ht . r + set-precipitation-parameter aln_p0 n s g	\$change nucleation to occur in austenite
edit-ht-segment ht . r + set-precipitation-parameter aln_p1 n s d	\$nucleation of AIN at gb in austenite
edit-ht-segment ht . r + set-precipitation-parameter aln#01_p0 n s n	\$nucleation of AIN at d in austenite
edit-ht-segment ht . r + set-precipitation-parameter aln#01_p1 n s n	\$no nucleation of AIN at d in ferrite
edit-ht-segment ht . r + set-precipitation-parameter vn_p0 n s d	\$nucleation of VN at d in austenite
edit-ht-segment ht . r + set-precipitation-parameter vn#01_p0 n s n	\$no nucleation of VN in ferrite
edit-ht-segment ht . r + set-precipitation-parameter vc_p0 n s d	\$nucleation of VC at d in austenite
edit-ht-segment ht . r + set-precipitation-parameter vc#01_p0 n s n	\$no nucleation of VC in ferrite
edit-ht-segment ht . r + set-precipitation-parameter nbn_p0 n s d	\$nucleation of NbN at d in austenite
edit-ht-segment ht . r + set-precipitation-parameter nbn#01_p0 n s n	\$no nucleation of NbN in ferrite
edit-ht-segment ht . r + set-precipitation-parameter nbc_p0 n s d	\$nucleation of NbC at d in austenite
edit-ht-segment ht . r + set-precipitation-parameter nbc#01_p0 n s n	\$no nucleation of NbC in ferrite
edit-ht-segment ht . r + set-precipitation-parameter m7c3_p0 n s n	\$no nucleation of M7C3 in ferrite
edit-ht-segment ht . r + set-precipitation-parameter m23c6_p0 n s n	\$no nucleation of M23C6 in ferrite
edit-ht-segment ht . r + set-precipitation-parameter Cr2N_p0 n s n	\$no nucleation of Cr2N in ferrite
edit-ht-segment ht . r + set-precipitation-parameter cementite_p0 n s n	\$no nucleation of cementite in ferrite

```

append-ht-segment ht                                $append new segment to heat treatment
edit-ht-segment ht . d y                            $use inherent precipitation domain
edit-ht-segment ht . 3 1100 3600                    $isothermal annealing at 1100°C for 1h

append-ht-segment ht                                $append new segment to heat treatment
edit-ht-segment ht . d y                            $use inherent precipitation domain
edit-ht-segment ht . 1 900 2                        $cool to 900°C with a cooling rate of 2°C/s

append-ht-segment ht                                $append new segment to heat treatment
edit-ht-segment ht . d y                            $use inherent precipitation domain
edit-ht-segment ht . 3 900 60                       $isothermal annealing at 900°C for 60s
edit-ht-segment ht . r + set-precipitation-parameter austenite d e 1e15 $increase dislocation density due to rolling

append-ht-segment ht                                $append new segment to heat treatment
edit-ht-segment ht . d y                            $use inherent precipitation domain
edit-ht-segment ht . 1 800 2                        $cool to 800°C with a cooling rate of 2°C/s

append-ht-segment ht                                $append new segment to heat treatment
edit-ht-segment ht . d n ferrite                    $change precipitation domain
edit-ht-segment ht . 1 25 1                          $cool to 25°C with a cooling rate of 1°C/s

edit-ht-segment ht . r + set-precipitation-parameter aln_p0 n s n $change nucleation to occur in ferrite
edit-ht-segment ht . r + set-precipitation-parameter aln_p1 n s n $no nucleation of aln at gb in austenite
edit-ht-segment ht . r + set-precipitation-parameter aln#01_p0 n s g $no nucleation of aln at d in austenite
edit-ht-segment ht . r + set-precipitation-parameter aln#01_p1 n s d $nucleation of aln at gb in ferrite
edit-ht-segment ht . r + set-precipitation-parameter vn_p0 n s n $no nucleation of aln at d in ferrite
edit-ht-segment ht . r + set-precipitation-parameter vn#01_p0 n s d $no nucleation of vn in austenite
edit-ht-segment ht . r + set-precipitation-parameter vc_p0 n s n $nucleation of vn at d in ferrite
edit-ht-segment ht . r + set-precipitation-parameter vcn#01_p0 n s d $no nucleation of vc in austenite
edit-ht-segment ht . r + set-precipitation-parameter nbn_p0 n s n $nucleation of vc at d in ferrite
edit-ht-segment ht . r + set-precipitation-parameter nbn#01_p0 n s d $no nucleation of nbn in austenite
edit-ht-segment ht . r + set-precipitation-parameter nbc_p0 n s n $nucleation of nbn at d in ferrite
edit-ht-segment ht . r + set-precipitation-parameter nbc#01_p0 n s d $no nucleation of nbc in austenite
edit-ht-segment ht . r + set-precipitation-parameter m7c3_p0 n s d $nucleation of nbc at d in ferrite
edit-ht-segment ht . r + set-precipitation-parameter m23c6_p0 n s d $no nucleation of m7c3 at d in ferrite
edit-ht-segment ht . r + set-precipitation-parameter Cr2N_p0 n s d $nucleation of m23c6 at d in ferrite
edit-ht-segment ht . r + set-precipitation-parameter cementite_p0 n s d $nucleation of cr2n at d in ferrite
edit-ht-segment ht . r + set-precipitation-parameter cementite_p0 n s d $nucleation of cementite at d in ferrite

$ ----- VARIABLES -----

set-variable-value npc 20                            $number of precipitate classes

$ ----- TABLES -----

create-global-table cie_vn                            $correction factors for CIE of VN
add-table-entry cie_vn                                1073                1                    $MC interpolates between defined values
add-table-entry cie_vn                                1173                1
add-table-entry cie_vn                                1223                0,95
add-table-entry cie_vn                                1273                0,78
add-table-entry cie_vn                                1373                0,65
add-table-entry cie_vn                                1473                0,65

$ *****
$                                     PRECIPITATION DOMAINS, PRECIPITATES
$ *****

$ ----- PRECIPITATION DOMAINS -----

create-precipitation-domain austenite                $austenite is precipitation domain = matrix

```

set-precipitation-parameter austenite x fcc_a1	\$matrix phase of domain „austenite“
set-precipitation-parameter austenite d e 1e11	\$dislocation density (=1e11 m <sup>-2</sup> )
set-precipitation-parameter austenite g 100e-6	\$austenite grain size (=100e-6 m)
set-precipitation-parameter austenite t (193000-73,333*T\$C)*1e6	\$temperature dependent young's modulus
create-precipitation-domain ferrite	\$ferrite is precipitation domain = matrix
set-precipitation-parameter ferrite x bcc_a2	\$matrix phase of domain „ferrite“
set-precipitation-parameter ferrite d e 1e12	\$dislocation density (=1e12 m <sup>-2</sup> )
set-precipitation-parameter ferrite g 10e-6	\$ferrite grain size (=100e-6 m)
set-precipitation-parameter ferrite t (210000-75*T\$C)*1e6	\$temperature dependent young's modulus

§ ----- PRECIPITATE PHASES IN AUSTENITE -----

create-new-phase aln p	\$new precipitate phase aln_p0 in austenite
set-precipitation-parameter aln_p0 c npc	\$use variable npc for number of prec. classes
set-precipitation-parameter aln_p0 d austenite	\$austenite is precipitation domain
set-precipitation-parameter aln_p0 n s g	\$nucleation sites are grain boundaries
set-precipitation-parameter aln_p0 i y	\$use calculated interfacial energy (CIE)
set-precipitation-parameter aln_p0 z y	\$use automatic size correction for CIE
set-precipitation-parameter aln_p0 s g (10^(11-0,005*T\$C))	\$ratio between diffusion at gb and bulk
set-precipitation-parameter aln_p0 s n (10^(11-0,005*T\$C))	\$ratio between diffusion at gb and bulk
create-new-phase aln p	\$new precipitate phase aln_p1 in austenite
set-precipitation-parameter aln_p1 c npc	\$use variable npc for number of prec. classes
set-precipitation-parameter aln_p1 d austenite	\$austenite is precipitation domain
set-precipitation-parameter aln_p1 n s d	\$nucleation sites are dislocations
set-precipitation-parameter aln_p1 i y	\$use calculated interfacial energy (CIE)
set-precipitation-parameter aln_p1 z y	\$use automatic size correction for CIE
set-precipitation-parameter aln_p1 t m n 0,19	\$use volumetric misfit between prec./matrix
set-precipitation-parameter aln_p1 n f y	\$use 0.27 instead 0.19 for MC versions ≥ 5.40
set-precipitation-parameter aln_p1 n f y	\$take into account volume misfit in nucleation
create-new-phase vn p	\$ new precipitate phase vn_p0 in austenite
set-precipitation-parameter vn_p0 c npc	\$use variable npc for number of prec. classes
set-precipitation-parameter vn_p0 d austenite	\$austenite is precipitation domain
set-precipitation-parameter vn_p0 n s d	\$nucleation sites are dislocations
set-precipitation-parameter vn_p0 i N cie\$vn*cie_vn(T)	\$correction of CIE due to entropic contr.
set-precipitation-parameter vn_p0 z y	\$use automatic size correction for CIE
create-new-phase vc p	\$new precipitate phase vc_p0 in austenite
set-precipitation-parameter vc_p0 c npc	\$use variable npc for number of prec. classes
set-precipitation-parameter vc_p0 d austenite	\$austenite is precipitation domain
set-precipitation-parameter vc_p0 n s d	\$nucleation sites are dislocations
set-precipitation-parameter vc_p0 i y	\$use calculated interfacial energy (CIE)
set-precipitation-parameter vc_p0 z y	\$use automatic size correction for CIE
create-new-phase nbn p	\$ new precipitate phase nbn_p0 in austenite
set-precipitation-parameter nbn_p0 c npc	\$use variable npc for number of prec. classes
set-precipitation-parameter nbn_p0 d austenite	\$austenite is precipitation domain
set-precipitation-parameter nbn_p0 n s d	\$nucleation sites are dislocations
set-precipitation-parameter nbn_p0 i y	\$use calculated interfacial energy (CIE)
set-precipitation-parameter nbn_p0 z y	\$use automatic size correction for CIE
create-new-phase nbc p	\$new precipitate phase nbc_p0 in austenite
set-precipitation-parameter nbc_p0 c npc	\$use variable npc for number of prec. classes
set-precipitation-parameter nbc_p0 d austenite	\$austenite is precipitation domain
set-precipitation-parameter nbc_p0 n s d	\$nucleation sites are dislocations
set-precipitation-parameter nbc_p0 i y	\$use calculated interfacial energy (CIE)
set-precipitation-parameter nbc_p0 z y	\$use automatic size correction for CIE

\$-----PRECIPITATE PHASES IN FERRITE-----	
remove_phase fcc_A1#01	\$remove the fcc_A1#01 precipitate phase
create-new-phase aln e	\$ new equilibrium phase for aln#01
create-new-phase vn e	\$ new equilibrium phase for vn#01
create-new-phase vc e	\$ new equilibrium phase for vc#01
create-new-phase nbn e	\$ new equilibrium phase for nbn#01
create-new-phase nbc e	\$ new equilibrium phase for nbc#01
create-new-phase aln#01 p	\$ new precipitate phase aln#01_p0 in ferrite
set-precipitation-parameter aln#01_p0 c npc	\$use variable npc for number of prec. classes
set-precipitation-parameter aln#01_p0 d ferrite	\$ferrite is precipitation domain
set-precipitation-parameter aln#01_p0 n s g	\$nucleation sites are grain boundaries
set-precipitation-parameter aln#01_p0 i y	\$use calculated interfacial energy (CIE)
set-precipitation-parameter aln#01_p0 z y	\$use automatic size correction for CIE
set-precipitation-parameter aln#01_p0 s g ( $10^{(12-0,01*T/C)}$ )	\$ratio between diffusion at gb and bulk
set-precipitation-parameter aln#01_p0 s n ( $10^{(12-0,01*T/C)}$ )	\$ratio between diffusion at gb and bulk
create-new-phase aln#01 p1	\$new precipitate phase aln#01_p1 in ferrite
set-precipitation-parameter aln#01_p1 c npc	\$use variable npc for number of prec. classes
set-precipitation-parameter aln#01_p1 d ferrite	\$ferrite is precipitation domain
set-precipitation-parameter aln#01_p1 n s d	\$nucleation sites are dislocations
set-precipitation-parameter aln#01_p1 i y	\$use calculated interfacial energy (CIE)
set-precipitation-parameter aln#01_p1 z y	\$use automatic size correction for CIE
set-precipitation-parameter aln#01_P1 t m n 0,19	\$use volumetric misfit between prec./matrix
set-precipitation-parameter aln#01_p1 n f y	\$use 0.27 instead 0.19 for MC versions $\leq 5.40$
set-precipitation-parameter aln#01_p1 n f y	\$take into account volume misfit in nucleation
create-new-phase vn#01 p	\$ new precipitate phase vn#01_p0 in ferrite
set-precipitation-parameter vn#01_p0 c npc	\$use variable npc for number of prec. classes
set-precipitation-parameter vn#01_p0 d ferrite	\$ferrite is precipitation domain
set-precipitation-parameter vn#01_p0 n s d	\$nucleation sites are dislocations
set-precipitation-parameter vn#01_p0 i y	\$use calculated interfacial energy (CIE)
set-precipitation-parameter vn#01_p0 z y	\$use automatic size correction for CIE
create-new-phase vc#01 p	\$new precipitate phase vc#01_p0 in ferrite
set-precipitation-parameter vc#01_p0 c npc	\$use variable npc for number of prec. classes
set-precipitation-parameter vc#01_p0 d ferrite	\$ferrite is precipitation domain
set-precipitation-parameter vc#01_p0 n s d	\$nucleation sites are dislocations
set-precipitation-parameter vc#01_p0 i y	\$use calculated interfacial energy (CIE)
set-precipitation-parameter vc#01_p0 z y	\$use automatic size correction for CIE
create-new-phase nbn#01 p	\$ new precipitate phase nbn#01_p0 in ferrite
set-precipitation-parameter nbn#01_p0 c npc	\$use variable npc for number of prec. classes
set-precipitation-parameter nbn#01_p0 d ferrite	\$ferrite is precipitation domain
set-precipitation-parameter nbn#01_p0 n s d	\$nucleation sites are dislocations
set-precipitation-parameter nbn#01_p0 i y	\$use calculated interfacial energy (CIE)
set-precipitation-parameter nbn#01_p0 z y	\$use automatic size correction for CIE
create-new-phase nbc#01 p	\$new precipitate phase nbc#01_p0 in ferrite
set-precipitation-parameter nbc#01_p0 c npc	\$use variable npc for number of prec. classes
set-precipitation-parameter nbc#01_p0 d ferrite	\$ferrite is precipitation domain
set-precipitation-parameter nbc#01_p0 n s d	\$nucleation sites are dislocations
set-precipitation-parameter nbc#01_p0 i y	\$use calculated interfacial energy (CIE)
set-precipitation-parameter nbc#01_p0 z y	\$use automatic size correction for CIE
create-new-phase m7c3 p	\$new precipitate phase m7c3_p0 in ferrite
set-precipitation-parameter m7c3_p0 c npc	\$use variable npc for number of prec. classes
set-precipitation-parameter m7c3_p0 d ferrite	\$ferrite is precipitation domain

```

set-precipitation-parameter m7c3_p0 n s d          $nucleation sites are dislocations
set-precipitation-parameter m7c3_p0 i y          $use calculated interfacial energy (CIE)
set-precipitation-parameter m7c3_p0 z y          $use automatic size correction for CIE

create-new-phase m23c6 p
set-precipitation-parameter m23c6_p0 c npc       $use variable npc for number of prec. classes
set-precipitation-parameter m23c6_p0 d ferrite   $ferrite is precipitation domain
set-precipitation-parameter m23c6_p0 n s d       $nucleation sites are dislocations
set-precipitation-parameter m23c6_p0 i y         $use calculated interfacial energy (CIE)
set-precipitation-parameter m23c6_p0 z y         $use automatic size correction for CIE

create-new-phase Cr2N p
set-precipitation-parameter Cr2N_p0 c npc        $use variable npc for number of prec. classes
set-precipitation-parameter Cr2N_p0 d ferrite   $ferrite is precipitation domain
set-precipitation-parameter Cr2N_p0 n s d       $nucleation sites are dislocations
set-precipitation-parameter Cr2N_p0 i y         $use calculated interfacial energy (CIE)
set-precipitation-parameter Cr2N_p0 z y         $use automatic size correction for CIE

create-new-phase cementite p
set-precipitation-parameter cementite_p0 c npc   $use variable npc for number of prec. classes
set-precipitation-parameter cementite_p0 d ferrite $ferrite is precipitation domain
set-precipitation-parameter cementite_p0 n s d   $nucleation sites are dislocations
set-precipitation-parameter cementite_p0 n c o   $nucleus composition: ortho-equilibrium
set-precipitation-parameter cementite_p0 i y     $use calculated interfacial energy (CIE)
set-precipitation-parameter cementite_p0 z y     $use automatic size correction for CIE

$ *****
$ ***** SIMULATION SETUP *****
$ *****

set-temperature-celsius 1000                      $define something
set-automatic-startvalues                         $initiate equilibrium calculation
calc-equilibrium                                 $calculate equilibrium state
set-simulation-parameter e 1e24                  $set simulation end time
set-simulation-parameter t h ht 10                $T-control from defined heat treatment
set-simulation-parameter s r                      $simulation starting conditions: reset

$ *****
$ ***** OUTPUT WINDOWS, PLOTS, ETC. *****
$ *****

new-gui-window p1                                 $generate new plot: temperature
                                                $define values for default x-axis
set-gui-window-property . x stepvalue            $default x-axis variable (time)
set-gui-window-property . s u y                  $use default x-axis for all plots: yes
set-gui-window-property . s t time / s          $default x-axis title
set-gui-window-property . s f 1                  $scaling factor is 1
set-plot-option . s n b t$c                      $add series: temperature
set-plot-option . s m -1 t$c T                   $define series legend
set-plot-option . a y 1 t temperture / °C        $y-axis title
set-plot-option . l n                             $no legend

create-new-plot x .
set-plot-option . s n b f$aIn_p0                 $create new plot: phase fractions
set-plot-option . s m -1 f$aIn_p0 AIN(gamma,gb) $add series: phase fraction aln_p0
set-plot-option . s n b f$aIn_p1                 $rename phase aln_p0 to AIN(gamma,gb)
set-plot-option . s m -1 f$aIn_p1 AIN(gamma,d)   $add series: phase fraction aln_p1
set-plot-option . s n b f$vn_p0                  $rename phase aln_p1 to AIN(gamma,d)
set-plot-option . s m -1 f$vn_p0 VN(gamma)       $add series: phase fraction vn_p0
set-plot-option . s n b f$vc_p0                  $rename phase vn_p0 to VN(gamma)
set-plot-option . s m -1 f$vc_p0                 $add series: phase fraction vc_p0

```

```

set-plot-option . s m -1 f$vc_p0 VC(gamma)
set-plot-option . s n b f$nb_n_p0
set-plot-option . s m -1 f$nb_n_p0 NbN(gamma)
set-plot-option . s n b f$nb_c_p0
set-plot-option . s m -1 f$nb_c_p0 NbC(gamma)
set-plot-option . s n b f$aln#01_p0
set-plot-option . s m -1 f$aln#01_p0 AlN(alpha,gb)
set-plot-option . s n b f$aln#01_p1
set-plot-option . s m -1 f$aln#01_p1 AlN(alpha,d)
set-plot-option . s n b f$vn#01_p0
set-plot-option . s m -1 f$vn#01_p0 VN(alpha)
set-plot-option . s n b f$vc#01_p0
set-plot-option . s m -1 f$vc#01_p0 VC(alpha)
set-plot-option . s n b f$nb_n#01_p0
set-plot-option . s m -1 f$nb_n#01_p0 NbN(alpha)
set-plot-option . s n b f$nb_c#01_p0
set-plot-option . s m -1 f$nb_c#01_p0 NbC(alpha)
set-plot-option . s n b f$m23c6_p0
set-plot-option . s m -1 f$m23c6_p0 M<sub>23</sub>C<sub>6</sub>(alpha)
set-plot-option . s n b f$m7c3_p0
set-plot-option . s m -1 f$m7c3_p0 M<sub>7</sub>C<sub>3</sub>(alpha)
set-plot-option . s n b f$Cr2N_p0
set-plot-option . s m -1 f$Cr2N_p0 Cr<sub>2</sub>N(alpha)
set-plot-option . s n b f$cementite_p0
set-plot-option . s m -1 f$cementite_p0 Cementite(alpha)
set-plot-option . a y 1 t phase fraction
set-plot-option . a y 1 y log
set-plot-option . a y 1 s 1e-8..
set-plot-option . l b

create-new-plot x .
set-plot-option . s n b r_mean$aln_p0
set-plot-option . s m -1 r_mean$aln_p0 AlN(gamma,gb)
set-plot-option . s n b r_mean$aln_p1
set-plot-option . s m -1 r_mean$aln_p1 AlN(gamma,d)
set-plot-option . s n b r_mean$vn_p0
set-plot-option . s m -1 r_mean$vn_p0 VN(gamma)
set-plot-option . s n b r_mean$vc_p0
set-plot-option . s m -1 r_mean$vc_p0 VC(gamma)
set-plot-option . s n b r_mean$nb_n_p0
set-plot-option . s m -1 r_mean$nb_n_p0 NbN(gamma)
set-plot-option . s n b r_mean$nb_c_p0
set-plot-option . s m -1 r_mean$nb_c_p0 NbC(gamma)
set-plot-option . s n b r_mean$aln#01_p0
set-plot-option . s m -1 r_mean$aln#01_p0 AlN(alpha,gb)
set-plot-option . s n b r_mean$aln#01_p1
set-plot-option . s m -1 r_mean$aln#01_p1 AlN(alpha,d)
set-plot-option . s n b r_mean$vn#01_p0
set-plot-option . s m -1 r_mean$vn#01_p0 VN(alpha)
set-plot-option . s n b r_mean$vc#01_p0
set-plot-option . s m -1 r_mean$vc#01_p0 VC(alpha)
set-plot-option . s n b r_mean$nb_n#01_p0
set-plot-option . s m -1 r_mean$nb_n#01_p0 NbN(alpha)
set-plot-option . s n b r_mean$nb_c#01_p0
set-plot-option . s m -1 r_mean$nb_c#01_p0 NbC(alpha)
set-plot-option . s n b r_mean$m23c6_p0
set-plot-option . s m -1 r_mean$m23c6_p0 M<sub>23</sub>C<sub>6</sub>(alpha)
set-plot-option . s n b r_mean$m7c3_p0
set-plot-option . s m -1 r_mean$m7c3_p0 M<sub>7</sub>C<sub>3</sub>(alpha)
set-plot-option . s n b r_mean$Cr2N_p0

$rename phase vc_p0 to VC(gamma)
$add series: phase fraction nb_n_p0
$rename phase nb_n_p0 to NbN(gamma)
$add series: phase fraction nb_c_p0
$rename phase nb_c_p0 to NbC(gamma)
$add series: phase fraction aln#01_p0
$rename phase aln#01_p0 to AlN(alpha,gb)
$add series: phase fraction aln#01_p1
$rename phase aln#01_p1 to AlN(alpha,d)
$add series: phase fraction vn#01_p0
$rename phase vn#01_p0 to VN(alpha)
$add series: phase fraction vc#01_p0
$rename phase vc#01_p0 to VC(alpha)
$add series: phase fraction nb_n#01_p0
$rename phase nb_n#01_p0 to NbN(alpha)
$add series: phase fraction nb_c#01_p0
$rename phase nb_c#01_p0 to NbC(alpha)
$add series: phase fraction m23c6_p0
$rename phase m23c6_p0 to M23C6(alpha)
$add series: phase fraction m7c3_p0
$rename phase m7c3_p0 to M7C3(alpha)
$add series: phase fraction cr2n_p0
$rename phase cr2n_p0 to Cr2N(alpha)
$add series: phase fraction cementite_p0
$rename phase cementite_p0 to Cem. (alpha)
$change y-axis title
$use logarithmic scale for y-axis
$scale the y-axis from 1e-8..
$show legend at the bottom of the figure

$create new plot: mean radii
$add series: mean radius aln_p0
$rename phase aln_p0 to AlN(gamma,gb)
$add series: mean radius aln_p1
$rename phase aln_p1 to AlN(gamma,d)
$add series: mean radius vn_p0
$rename phase vn_p0 to VN(gamma)
$add series: mean radius vc_p0
$rename phase vc_p0 to VC(gamma)
$add series: mean radius nb_n_p0
$rename phase nb_n_p0 to NbN(gamma)
$add series: mean radius nb_c_p0
$rename phase nb_c_p0 to NbC(gamma)
$add series: mean radius aln#01_p0
$rename phase aln#01_p0 to AlN(alpha,gb)
$add series: mean radius aln#01_p1
$rename phase aln#01_p1 to AlN(alpha,d)
$add series: mean radius vn#01_p0
$rename phase vn#01_p0 to VN(alpha)
$add series: mean radius vc#01_p0
$rename phase vc#01_p0 to VC(alpha)
$add series: mean radius nb_n#01_p0
$rename phase nb_n#01_p0 to NbN(alpha)
$add series: mean radius nb_c#01_p0
$rename phase nb_c#01_p0 to NbC(alpha)
$add series: mean radius m23c6_p0
$add series: mean radius m7c3_p0
$add series: mean radius cr2n_p0

```

```

set-plot-option . s m -1 r_mean$Cr2N_p0 Cr<sub>2</sub>N(alpha)      $rename phase cr2n_p0 to Cr2N(alpha)
set-plot-option . s n b r_mean$cementite_p0                        $add series: mean radius cementite_p0
set-plot-option . s m -1 r_mean$cementite_p0 Cementite(alpha)    $rename phase cementite_p0 to Cem.(alpha)
set-plot-option . a y 1 t mean radius / nm                       $change y-axis title
set-plot-option . a y 1 y log                                     $use logarithmic scale for y-axis
set-plot-option . a y 1 f 1e9                                     $scaling factor is 1e9
set-plot-option . l b                                            $show legend at the bottom of the figure

create-new-plot x .
set-plot-option . s n b num_part$aln_p0                          $add series: number density aln_p0
set-plot-option . s m -1 num_part$aln_p0 AIN(gamma,gb)          $rename phase aln_p0 to AIN(gamma,gb)
set-plot-option . s n b num_part$aln_p1                          $add series: number density aln_p1
set-plot-option . s m -1 num_part$aln_p1 AIN(gamma,d)           $rename phase aln_p1 to AIN(gamma,d)
set-plot-option . s n b num_part$vn_p0                          $add series: number density vn_p0
set-plot-option . s m -1 num_part$vn_p0 VN(gamma)               $rename phase vn_p0 to VN(gamma)
set-plot-option . s n b num_part$vc_p0                          $add series: number density vc_p0
set-plot-option . s m -1 num_part$vc_p0 VC(gamma)               $rename phase vc_p0 to VC(gamma)
set-plot-option . s n b num_part$nb_n_p0                        $add series: number density nb_n_p0
set-plot-option . s m -1 num_part$nb_n_p0 NbN(gamma)            $rename phase nb_n_p0 to NbN(gamma)
set-plot-option . s n b num_part$nb_c_p0                        $add series: number density nb_c_p0
set-plot-option . s m -1 num_part$nb_c_p0 NbC(gamma)            $rename phase nb_c_p0 to NbC(gamma)
set-plot-option . s n b num_part$aln#01_p0                      $add series: number density aln#01_p0
set-plot-option . s m -1 num_part$aln#01_p0 AIN(alpha,gb)       $rename phase aln#01_p0 to AIN(alpha,gb)
set-plot-option . s n b num_part$aln#01_p1                      $add series: number density aln#01_p1
set-plot-option . s m -1 num_part$aln#01_p1 AIN(alpha,d)        $rename phase aln#01_p1 to AIN(alpha,d)
set-plot-option . s n b num_part$vn#01_p0                      $add series: number density vn#01_p0
set-plot-option . s m -1 num_part$vn#01_p0 VN(alpha)           $rename phase vn#01_p0 to VN(alpha)
set-plot-option . s n b num_part$vc#01_p0                      $add series: number density vc#01_p0
set-plot-option . s m -1 num_part$vc#01_p0 VC(alpha)           $rename phase vc#01_p0 to VC(alpha)
set-plot-option . s n b num_part$nb_n#01_p0                    $add series: number density nb_n#01_p0
set-plot-option . s m -1 num_part$nb_n#01_p0 NbN(alpha)        $rename phase nb_n#01_p0 to NbN(alpha)
set-plot-option . s n b num_part$nb_c#01_p0                    $add series: number density nb_c#01_p0
set-plot-option . s m -1 num_part$nb_c#01_p0 NbC(alpha)        $rename phase nb_c#01_p0 to NbC(alpha)
set-plot-option . s n b num_part$m23c6_p0                      $add series: number density m23c6_p0
set-plot-option . s m -1 num_part$m23c6_p0 M<sub>23</sub>C<sub>6</sub>(alpha)
set-plot-option . s n b num_part$m7c3_p0                        $add series: number density m7c3_p0
set-plot-option . s m -1 num_part$m7c3_p0 M<sub>7</sub>C<sub>3</sub>(alpha)
set-plot-option . s n b num_part$Cr2N_p0                        $add series: number density cr2n_p0
set-plot-option . s m -1 num_part$Cr2N_p0 Cr<sub>2</sub>N(alpha) $rename phase cr2n_p0 to Cr2N(alpha)
set-plot-option . s n b num_part$cementite_p0                  $add series: number density cementite_p0
set-plot-option . s m -1 num_part$cementite_p0 Cementite(alpha) $rename phase cementite_p0 to Cem.(alpha)
set-plot-option . a y 1 t number density / m<sup>-3</sup>        $change y-axis title
set-plot-option . a y 1 y log                                   $use logarithmic scale for y-axis
set-plot-option . a y 1 s 1..                                  $scale the y-axis from 1..
set-plot-option . l b                                          $show legend at the bottom of the figure

move-gui-window . 10 10 600 1000                                $move window to new position and resize
update-gui-window .                                             $update the GUI window

$ ----- HISTOGRAMM -----

new-gui-window p5                                               $new precipitation distribution plot window

set-plot-option . s n p aln_p0                                  $precipitate distribution AIN
set-plot-option . s m -1 aln_p0 AIN(gamma,gb)                  $add series: precipitate distribution aln_p0
set-plot-option . s n p aln_p1                                  $rename phase aln_p0 to AIN(gamma,gb)
set-plot-option . s m -1 aln_p1 AIN(gamma,d)                   $add series: precipitate distribution aln_p1
set-plot-option . s n p aln#01_p0                              $rename phase aln_p1 to AIN(gamma,d)
set-plot-option . s n p aln#01_p0                              $add series: precipitate distr. aln#01_p0
set-plot-option . s m -1 aln#01_p0 AIN(alpha,gb)               $rename phase aln#01_p0 to AIN(alpha,gb)

```

set-plot-option . s n p aln#01_p1	\$add series: precipitate distr. aln#01_p1
set-plot-option . s m -1 aln#01_p1 AIN(alpha,d)	\$rename phase aln#01_p1 to AIN(alpha,d)
set-plot-option . a x 1 f 1e9	\$use scaling factor 1e9 for x-axis
set-plot-option . a y 1 s 1e10..1e20	\$scale the y-axis from 1e10..1e20
set-plot-option . a y 1 y log	\$use logarithmic scale for y-axis
set-plot-option . a x 1 t radius / nm	\$change x-axis title
set-plot-option . a x 1 s 1..500	\$scale the x-axis from 1..500
set-plot-option . a x 1 y log	\$use logarithmic scale for x-axis
set-plot-option . a y 1 t number density	\$change y-axis title
create-new-plot p .	\$create new plot: precipitate distribution VN
set-plot-option . s n p vn_p0	\$add series: precipitate distribution vn_p0
set-plot-option . s m -1 vn_p0 VN(gamma)	\$rename phase vn_p0 to VN(gamma)
set-plot-option . s n p vn#01_p0	\$add series: precipitate distribution vn#01_p0
set-plot-option . s m -1 vn#01_p0 VN(alpha)	\$rename phase vn#01_p0 to VN(alpha)
set-plot-option . a x 1 f 1e9	\$use scaling factor 1e9 for x-axis
set-plot-option . a y 1 s 1e10..1e20	\$scale the y-axis from 1e10..1e20
set-plot-option . a y 1 y log	\$use logarithmic scale for y-axis
set-plot-option . a x 1 t radius / nm	\$change x-axis title
set-plot-option . a x 1 s 1..500	\$scale the x-axis from 1..500
set-plot-option . a x 1 y log	\$use logarithmic scale for x-axis
set-plot-option . a y 1 t number density	\$change y-axis title
create-new-plot p .	\$create new plot: precipitate distribution VC
set-plot-option . s n p vc_p0	\$add series: precipitate distribution vc_p0
set-plot-option . s m -1 vc_p0 VC(gamma)	\$rename phase vc_p0 to VC(gamma)
set-plot-option . s n p vc#01_p0	\$add series: precipitate distribution vc#01_p0
set-plot-option . s m -1 vc#01_p0 VC(alpha)	\$rename phase vc#01_p0 to VC(alpha)
set-plot-option . a x 1 f 1e9	\$use scaling factor 1e9 for x-axis
set-plot-option . a y 1 s 1e10..1e20	\$scale the y-axis from 1e10..1e20
set-plot-option . a y 1 y log	\$use logarithmic scale for y-axis
set-plot-option . a x 1 t radius / nm	\$change x-axis title
set-plot-option . a x 1 s 1..500	\$scale the x-axis from 1..500
set-plot-option . a x 1 y log	\$use logarithmic scale for x-axis
set-plot-option . a y 1 t number density	\$change y-axis title
create-new-plot p .	\$create new plot: precipitate distribution NbN
set-plot-option . s n p nbn_p0	\$add series: precipitate distribution nbn_p0
set-plot-option . s m -1 nbn_p0 NbN(gamma)	\$rename phase nbn_p0 to NbN(gamma)
set-plot-option . s n p nbn#01_p0	\$add series: precipitate distr. nbn#01_p0
set-plot-option . s m -1 nbn#01_p0 NbN(alpha)	\$rename phase nbn#01_p0 to NbN(alpha)
set-plot-option . a x 1 f 1e9	\$use scaling factor 1e9 for x-axis
set-plot-option . a y 1 s 1e10..1e20	\$scale the y-axis from 1e10..1e20
set-plot-option . a y 1 y log	\$use logarithmic scale for y-axis
set-plot-option . a x 1 t radius / nm	\$change x-axis title
set-plot-option . a x 1 s 1..500	\$scale the x-axis from 1..500
set-plot-option . a x 1 y log	\$use logarithmic scale for x-axis
set-plot-option . a y 1 t number density	\$change y-axis title
create-new-plot p .	\$create new plot: precipitate distribution NbC
set-plot-option . s n p nbc_p0	\$add series: precipitate distribution nbc_p0
set-plot-option . s m -1 nbc_p0 NbC(gamma)	\$rename phase nbn_p0 to NbC(gamma)
set-plot-option . s n p nbc#01_p0	\$add series: precipitate distr. nbc#01_p0
set-plot-option . s m -1 nbc#01_p0 NbC(alpha)	\$rename phase nbc#01_p0 to NbC(alpha)
set-plot-option . a x 1 f 1e9	\$use scaling factor 1e9 for x-axis
set-plot-option . a y 1 s 1e10..1e20	\$scale the y-axis from 1e10..1e20
set-plot-option . a y 1 y log	\$use logarithmic scale for y-axis
set-plot-option . a x 1 t radius / nm	\$change x-axis title
set-plot-option . a x 1 s 1..500	\$scale the x-axis from 1..500
set-plot-option . a x 1 y log	\$use logarithmic scale for x-axis



```
set-plot-option . a y 1 t number density           $change y-axis title

move-gui-window . 615 10 400 1000                 $move window to new position and resize
update-gui-window . $ Update the GUI window       $update the GUI window

$ *****
$ ***** START PRECIPITATE SIMULATION *****
$ *****

start-precipitate-simulation

$ *****
$ ***** PRECIPITATE SIMULATION FINISHED *****
$ *****
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## **B Additional Selected Scientific Investigations carried out from 2007 to 2010**

Beside the numerical analysis of the precipitation kinetics of nitrides and carbides in microalloyed steel, the author of the present thesis carried out additional scientific investigations during his time as a PhD student at the Christian Doppler Laboratory 'Early Stages of Precipitation'.

Particularly, the precipitation kinetics of  $\gamma'$  and  $\delta$  phases in nickel-based superalloys is of great scientific interest. Therefore detailed investigations are performed on the precipitation process of these phases in the commercial nickel-based superalloys UDIMET 720 Li and ATI Allvac<sup>®</sup>718Plus<sup>™</sup>, see appendix B1 to B4.

Furthermore, the precipitation kinetics of NiAl particles is studied along with the formation of reverted austenite in a PH 13-8 Mo maraging steel, see appendix B5.

In the following, these additional scientific investigations are presented, which are also published in international accepted journals.